

GEOMETRIC QUANTIZATION VIA SYZ TRANSFORMS

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ABSTRACT. The so-called quantization problem in geometric quantization is asking whether the space of wave functions is independent of the choice of polarization. In this paper, we apply SYZ transforms to solve the quantization problem in two cases:

- (1) semi-flat Lagrangian torus fibrations over complete compact integral affine manifolds, and
- (2) projective toric manifolds.

More precisely, we prove that the space of wave functions associated to the real polarization is canonically isomorphic to that associated to a complex polarization via SYZ transforms in both cases.

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1. INTRODUCTION

1.1. The quantization problem. Physicists are keen on quantizing classical theories. Given any classical theory, they aim at finding a quantum counterpart whose classical limit recovers the original classical theory. This process is called quantization. Geometric quantization is a mathematical approach to constructing such quantum theories.

A system of classical mechanics is a phase space formalized as a symplectic manifold, so we let $(\check{X}, \check{\omega})$ be a symplectic manifold such that $\check{\omega}$ represents an

integral class $[\tilde{\omega}] \in H^2(\check{X}; \mathbb{Z})$. Then there exists a Hermitian line bundle $\mathcal{L} \rightarrow \check{X}$ which admits a unitary connection ∇ so that

$$\frac{i}{2\pi} F_{\nabla} = \tilde{\omega}.$$

Such a pair (\mathcal{L}, ∇) is called a *prequantum line bundle* (Definition 3.1). Given such a prequantum system, physicists would like to associate a Hilbert space \mathcal{H} to represent the wave functions or quantum states. A natural choice is to take \mathcal{H} as the (L^2 -completion of the) space of all sections of a line bundle \mathcal{L} . However, this space is in general too big to capture the actual physics. To obtain a Hilbert space of reasonable size, physicists introduced the notion of polarization.

A *polarization* is an involutive Lagrangian subbundle P of the complexified tangent bundle $T_{\mathbb{C}}\check{X}$ of \check{X} . Given a polarization, one can then define \mathcal{H} as the space $\Gamma_P(\check{X}, \mathcal{L})$ of *polarized sections* – sections that are covariantly constant along P (Definition 3.2). When $(\check{X}, \check{J}, \tilde{\omega})$ is a Kähler manifold admitting a Lagrangian fibration $\check{p} : \check{X} \rightarrow B$, there are two natural choices for the polarization, namely, the *real polarization* induced by the Lagrangian fibration \check{p} and the *complex polarization* induced by the complex structure \check{J} . But physicists believe that the quantum theory should be independent of the polarization we chose, so it is desirable that the space of polarized sections for the real polarization is canonically isomorphic to that for the complex polarization. This is known as the *quantization problem*.

In this paper, we show how *SYZ transforms*, which are Fourier-type transforms responsible for the interchange between complex-geometric data on \check{X} and symplectic-geometric data on the mirror X [24], can be applied to solve the quantization problem. We demonstrate this idea in two cases (Section 4.1):

- (1) semi-flat Lagrangian torus fibrations over complete compact integral affine manifolds (see Definition 4.4), and
- (2) projective toric manifolds.

1.2. Applying SYZ. Let us explain why SYZ transforms can help to solve the quantization problem. We begin with a semi-flat Lagrangian torus fibration $\check{p} : \check{X} \rightarrow B$ which admits a Lagrangian section.¹ The *SYZ mirror* of \check{X} is simply given by the total space of the fiberwise dual torus fibration $p : X \rightarrow B$. Exploiting the integral affine structure on the base B , we have the identifications

$$\check{X} \cong TB/\Lambda, \quad X \cong T^*B/\Lambda^*,$$

together with the fiberwise dual fibrations

$$\begin{array}{ccc} \check{X} & & X \\ & \searrow \check{p} & \swarrow p \\ & B & \end{array}$$

This is the toy model for SYZ mirror symmetry [24, 18].

Since a torus fiber $F_x := p^{-1}(x)$ is dual to the corresponding fiber $\check{F}_x := \check{p}^{-1}(x)$ (for some point $x \in B$), geometrically F_x is parametrizing the flat $U(1)$ -connections over \check{F}_x . So a section of p gives a $U(1)$ -connection over each fiber of \check{p} and they

¹In general there are singular or degenerate fibers in a Lagrangian torus fibration and the semi-flat one is obtained by restricting to the smooth locus.

combine to produce a $U(1)$ -connection over the total space \check{X} . A simple yet fundamental observation of Arinkin-Polishchuk [2] and Leung-Yau-Zaslow [19] is that the section L of p is *Lagrangian* if and only if the corresponding $U(1)$ -connection defines a *holomorphic* line bundle \check{L} over \check{X} . We call the map $L \mapsto \check{L}$ the *SYZ transform*.

Every Hermitian holomorphic line bundle over \check{X} comes from this construction. In particular, we can consider a prequantum line bundle

$$(\mathcal{L}, \nabla) = (\check{L}, \nabla_{\check{L}})$$

given by the SYZ transform of some Lagrangian section L of $p : X \rightarrow B$.

Śniatycki [23] proved that the space $\Gamma_{P_{\mathbb{R}}}(\check{X}, \check{L})$ of polarized sections for the real polarization $P_{\mathbb{R}}$ (or simply the space of real polarized sections) can be identified with the \mathbb{C} -vector space spanned by the so-called *Bohr-Sommerfeld fibers* of $\check{p} : \check{X} \rightarrow B$ which are, by definition, fibers of \check{p} over which $\nabla_{\check{L}}$ restricts to the trivial connection (see Definition 3.6):

$$\Gamma_{P_{\mathbb{R}}}(\check{X}, \check{L}) \cong \bigoplus_{\check{F}_x: \text{BS fiber}} \mathbb{C} \cdot \check{F}_x.$$

Under the SYZ transform, these fibers correspond precisely to the intersection points between L and the zero section L_0 of the dual fibration $p : X \rightarrow B$, so we have

$$(1) \quad \bigoplus_{\check{F}_x: \text{BS fiber}} \mathbb{C} \cdot \check{F}_x \cong \bigoplus_{p \in L_0 \cap L} \mathbb{C} \cdot p$$

via the SYZ transform.

On the other hand, the space $\Gamma_{P_{\mathbb{C}}}(\check{X}, \check{L})$ of polarized sections for the complex polarization $P_{\mathbb{C}}$ (or simply the space of complex polarized sections) is given by the space $H^0(\check{X}, \check{L})$ of holomorphic sections of \check{L} . Since \check{L} is ample, we have $\text{Ext}^{\bullet}(\mathcal{O}_{\check{X}}, \check{L}) = H^0(\check{X}, \check{L})$. Then the homological mirror symmetry (HMS) conjecture of Kontsevich [16] together with its compatibility with the SYZ conjecture [24] implies the isomorphism

$$H^0(\check{X}, \check{L}) \cong HF^{\bullet}(L_0, L),$$

again via the SYZ transform. If the intersection of L_0 and L is nice enough (e.g. if all the intersection points of L_0 and L are of Maslov index 0 so that the Floer differential \mathfrak{m}_1 is 0), then we further have

$$HF^{\bullet}(L_0, L) \cong \bigoplus_{p \in L_0 \cap L} \mathbb{C} \cdot p.$$

Combining, we arrive at the canonical isomorphism between the spaces of polarized sections

$$\Gamma_{P_{\mathbb{R}}}(\check{X}, \check{L}) \cong \Gamma_{P_{\mathbb{C}}}(\check{X}, \check{L}),$$

obtained by applying SYZ transforms *twice*.

Note that the space $HF^{\bullet}(L_0, L)$ in the above heuristic argument is only playing an auxiliary role. In fact, instead of using $HF^{\bullet}(L_0, L)$ and alluding to the HMS conjecture, we will consider the space

$$\ker(d_W) \cap A_{r,d}^0(P(L_0, L))$$

of d_W -closed functions which rapidly decay (hence the subscript “r.d.”) along the lattice directions of $P(L_0, L)$, where $P(L_0, L)$ is the *fiberwise geodesic path space* of the pair (L_0, L) equipped with the S^1 -valued *area function*

$$\mathcal{A} : P(L_0, L) \rightarrow S^1$$

and d_W is the *Witten differential*.

Applying Witten-Morse theory [26, 6, 20] to $P(L_0, L)$, we obtain the correspondence

$$\bigoplus_{p \in L_0 \cap L} \mathbb{C} \cdot p \cong \ker(d_W) \cap A_{r.d.}^0(P(L_0, L))$$

Then the SYZ transform (see Section 2.3)

$$\mathcal{F} : (A^\bullet(P(L_0, L)), d_W) \rightarrow (A^{0,\bullet}(\check{X}, \check{L}), \bar{\partial})$$

gives the identification

$$\ker(d_W) \cap A_{r.d.}^0(P(L_0, L)) \cong H^0(\check{X}, \check{L}).$$

Combining with (1), which was obtained via another SYZ transform, we solve the quantization problem in the semi-flat setting:

Theorem 1.1 (=Theorem 4.10). *Let $\check{p} : \check{X} \rightarrow B$ be a semi-flat Lagrangian torus fibration over a compact complete special integral affine manifold B . Let g be an integral (Definition 4.1) Hessian type metric on B . With respect to the prequantum line bundle $(\check{L}_g, \nabla_{\check{L}_g})$ associated to the metric g , the SYZ transform \mathcal{F} induces a canonical isomorphism between $\Gamma_{P_{\mathbb{R}}}(\check{X}, \check{L}_g)$ and $\Gamma_{P_{\mathbb{C}}}(\check{X}, \check{L}_g)$*

The prequantum line bundle $(\check{L}_g, \nabla_{\check{L}_g})$ here is explicitly given as the SYZ transform of the Lagrangian section

$$L_g := \{(x, d\phi(x)) \in X : x \in B\},$$

where ϕ is the local potential function for the Hessian type metric g .

Similar ideas can also be applied to projective toric manifolds:

Theorem 1.2 (=Theorem 4.26). *Let $\check{p} : \check{X} \rightarrow B$ be the moment map of a projective toric manifold \check{X} . With respect to the prequantum line bundle $(\check{L}_\phi, \nabla_{\check{L}_\phi})$, the SYZ transform \mathcal{F} induces a canonical isomorphism between $\Gamma_{P_{\mathbb{R}}}(\check{X}, \check{L}_\phi)$ and $\Gamma_{P_{\mathbb{C}}}(\check{X}, \check{L}_\phi)$.*

The Lagrangian section mirror to $(\check{L}_\phi, \nabla_{\check{L}_\phi})$ here is again given by the graph of the differential of the Kähler potential ϕ . However, in this case, due to degeneracy of the moment map fibers over the boundary ∂B of the moment polytope B , not only Bohr-Sommerfeld Lagrangian fibers, but also *Bohr-Sommerfeld isotropic fibers*, would contribute. Hence we also need to consider the intersection points of L_ϕ with L_0 over the boundary ∂B (or at infinity if one considers the Legendre dual $N_{\mathbb{R}}$ of B).

Detailed proofs of Theorems 1.1 and 1.2 will be given in Section 4.1.

Remark 1.3. *The Lagrangian sections L_0 and L are graded, and as pointed out in [17, Remark 13], the Maslov index of an intersection point of L_0 and L coincides with the Morse index of the corresponding critical point of \mathcal{A} . Since the Lagrangian section L satisfies the condition*

$$\text{Hess}(\mathcal{A})(p) > 0,$$

for every $p \in L_0 \cap L$, so indeed all the intersection points of L_0 and L are of Maslov index 0. It follows that

$$HF^\bullet(L_0, L) = HF^0(L_0, L) = \bigoplus_{p \in L_0 \cap L} \mathbb{C} \cdot p.$$

Hence we indeed have $HF^0(L_0, L) \cong H^0(\check{X}, \check{L})$ (as vector spaces), as predicted by the HMS conjecture.

1.3. Relation to other works. The idea of applying mirror symmetry to geometric quantization (or vice versa) has appeared several times in the literature before our work and must be well-known among experts.

As far as we know, Andrei Tyurin was the first to suggest that the *numerical quantization problem*, namely, the dimension equality

$$\dim \Gamma_{P_{\mathbb{R}}}(\check{X}, \check{L}) = \dim \Gamma_{P_{\mathbb{C}}}(\check{X}, \check{L}),$$

follows from the SYZ [24] and HMS [16] conjectures, at least in the case of elliptic curves and algebraic K3 surfaces. Our SYZ transform \mathcal{F} was called the *geometric Fourier transformation (GFT)* in his paper. Let us also point out that his calculations resembled those carried out in an earlier work [11] of Gross.

The belief that the space of polarized sections for a complex polarization $P_{\mathbb{C}}$ is canonically isomorphic to that for the real polarization suggests that when a Kähler manifold \check{X} admits a Lagrangian torus fibration (possibly with singular fibers), the space $H^0(\check{X}, \check{L})$ of holomorphic sections of a holomorphic line bundle $\check{L} \rightarrow \check{X}$ should have a *canonical basis* or a *theta basis*. This again was first stated by Tyurin in [25]. In his thesis [22], Nohara reviewed this idea and, combined with an earlier work of Andersen [1], reproved the numerical quantization equality for semi-flat Lagrangian torus fibrations over compact base.

In a series of works ([12] and upcoming works; see also the nice survey article [13] for an overview), Gross, Hacking, Keel and Siebert show that a large class of varieties including Calabi-Yau manifolds carry *theta functions*. Instead of Lagrangian fibrations and Floer cohomologies (or their de-Rham versions such as the space $\ker(d_W) \cap A_{r.d.}^0(P(L_0, L))$ that we use in this paper), they use tropical geometry and counts of so-called *broken lines* (tropical analogue of holomorphic disks) to construct theta functions and prove a strong form of Tyurin's conjecture. Their work was inspired by Tyurin's work, and the above heuristic argument for the existence of theta functions applying both the SYZ and HMS conjectures has actually already appeared on p.5 of [13]. See also [15, Section 8] for a nice review of these topics.

In [4], Baier, Mourão and Nunes solved the quantization problem for abelian varieties and, together with Florentino, they solved the same problem for projective toric manifolds in [3]. They considered a degenerating family of complex structures $\{\check{J}_s\}_{s \geq 0}$ approaching a large complex structure limit and which are compatible with a fixed symplectic structure $\tilde{\omega}$. As $s \rightarrow +\infty$, they proved that holomorphic sections of the pre-quantum line bundle would converge to distributional sections supported on Bohr-Sommerfeld fibers of the Lagrangian torus fibration, giving rise to the desired canonical isomorphism.

In our mirror symmetric approach, the same limiting process (here we use $\hbar > 0$ and let $\hbar \rightarrow 0$) occurs when we apply Witten-Morse theory to the fiberwise geodesic path space $P(L_0, L)$ [6, 20]. If s_{\hbar} is a holomorphic section of the prequantum line

$$\begin{array}{ccc}
H^0(\check{X}_{\hbar}, \check{L}_{\hbar}) & \xleftarrow{\text{SYZ transform}} & \ker(d_{W,\hbar}) \cap A_{r.d.}^0(P(L_0, L)) \\
\downarrow \begin{array}{l} \text{Degeneration of} \\ \text{complex structures} \\ \hbar \rightarrow 0 \end{array} & & \downarrow \begin{array}{l} \text{Witten deformation} \\ \hbar \rightarrow 0 \end{array} \\
\bigoplus \mathbb{C} \cdot \check{F}_x & \xleftarrow{\text{SYZ transform}} & \bigoplus_{p \in L_0 \cap L} \mathbb{C} \cdot p \\
\check{F}_x: \text{BS fiber} & &
\end{array}$$

FIGURE 1. Degenerating family of complex structures vs. Witten deformation

bundle \check{L}_{\hbar} , then its SYZ transform in $(A^\bullet(P(L_0, L)), d_W)$ (here d_W depends on \hbar) converges to a sum of δ -functions supported on the intersection points between L_0 and L (or the critical points of the area function \mathcal{A}) as $\hbar \rightarrow 0$. Applying the SYZ transform (or Fourier transform) in the reverse direction transforms these δ -functions to distributional sections supported on the Bohr-Sommerfeld fibers. In summary, under the SYZ transforms, the limiting process in [4, 3] is corresponding precisely to the Witten deformation, as illustrated in Figure 1.

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2. SEMI-FLAT SYZ MIRROR SYMMETRY

In this section, we review the constructions of the semi-flat mirror manifold and the SYZ transform of immersed Lagrangian multi-sections.

2.1. Mirror construction in the semi-flat case. Let B be an n -dimensional integral affine manifold, that is, the transition functions of B belongs to $GL(n, \mathbb{Z}) \times \mathbb{R}^n$, the group of \mathbb{Z} -affine linear map. Let $\Lambda \subset TB$ and $\Lambda^* \subset T^*B$ be the natural lattice bundles defined by the integral affine structure. More precisely, on a local affine chart $U \subset B$,

$$\Lambda(U) := \bigoplus_{j=1}^n \mathbb{Z} \cdot \frac{\partial}{\partial x^j}, \quad \Lambda^*(U) := \bigoplus_{j=1}^n \mathbb{Z} \cdot dx^j,$$

where (x^j) are affine coordinates of U . We set

$$X := T^*B/\Lambda^*, \quad \check{X} := TB/\Lambda.$$

Let $(y_j), (\check{y}^j)$ be fiber coordinates (which are dual to each other) of X and \check{X} respectively. Then (x^j, y_j) and (x^j, \check{y}^j) define a set of local coordinates on $T^*U/\Lambda^* \subset X$ and $TU/\Lambda \subset \check{X}$, respectively.

Equip X with the standard symplectic structure

$$\omega_{\hbar} := \hbar^{-1} \sum_j dy_j \wedge dx^j,$$

where $\hbar > 0$ is a small real parameter. There is a natural almost complex structure \check{J}_\hbar on \check{X} given by

$$\check{J}_\hbar \left(\frac{\partial}{\partial x^j} \right) = -\hbar^{-1} \frac{\partial}{\partial \check{y}^j} \text{ and } \check{J}_\hbar \left(\frac{\partial}{\partial \check{y}^j} \right) = \hbar \frac{\partial}{\partial x^j}.$$

It is easy to see that \check{J}_\hbar is indeed integrable with local complex coordinates given by $z_j = \check{y}^j + i\hbar^{-1}x^j$. Hence $(\check{X}, \check{J}_\hbar)$ defines a complex manifold. To obtain Calabi-Yau structure, we need the following

Definition 2.1. *An n -dimensional manifold B is called a special integral affine manifold if its transition functions sit in $SL(n, \mathbb{Z}) \rtimes \mathbb{R}^n$.*

Hence if B is an integral special affine manifold, then the canonical line bundle of \check{X} is trivial. Indeed,

$$\check{\Omega}_\hbar := dz_1 \wedge \cdots \wedge dz_n$$

defines a global holomorphic volume form on \check{X} .

Definition 2.2. *$(\check{X}, \check{J}_\hbar)$ is called the SYZ mirror of (X, ω_\hbar) .*

The \hbar -parameter gives us a family of symplectic manifolds. As $\hbar \rightarrow 0$, the symplectic volume of (X, ω_\hbar) approaches infinity, which is the so-called *large volume limit* of the family $\{(X, \omega_\hbar)\}_{\hbar>0}$.

Next, we equip Kähler structures on both X and \check{X} .

Definition 2.3. *A Riemannian metric on B is said to be of Hessian type if locally there is a smooth convex function ϕ such that the metric is locally given by*

$$\sum_{j,k=1}^n \frac{\partial^2 \phi}{\partial x^j \partial x^k} dx^j \otimes dx^k,$$

where (x^j) are local affine coordinate of B .

Suppose g is a Hessian Riemannian metric on B . Then we can use it to construct natural complex structure on X and symplectic structure on \check{X} such that they compatible with the natural symplectic and complex structure on X and \check{X} , respectively.

Proposition 2.4. *Given a Hessian Riemannian metric on B , there is a natural complex structure J_g on X such that (X, ω_\hbar, J_g) is a Kähler manifold.*

Proof. Define

$$dz_j = dy_j + i \sum_{k=1}^n g_{kj} dx^k.$$

Then one can easily check that $\{z_j\}$ define complex coordinates on X and the complex structure J_g is given by

$$J_g^*(dx^j) = \sum_{k=1}^n g^{jk} dy_k,$$

$$J_g^*(dy_j) = - \sum_{k=1}^n g_{kj} dx^k.$$

The symplectic form on X is given by

$$\omega_{\hbar} = \hbar^{-1} \sum_{j=1}^n dy_j \wedge dx^j.$$

Hence

$$J_g^* \omega_{\hbar} = -\hbar^{-1} \sum_{j,k,l=1}^n g^{jl} g_{kj} dx^k \wedge dy_l = \hbar^{-1} \sum_{k=1}^n dy_k \wedge dx^k = \omega_{\hbar}.$$

and so

$$\omega_{\hbar}(\cdot, J_g(\cdot)) = \sum_{j,k=1}^n (\hbar^{-2} g_{jk} dx^j \otimes dx^k + g^{jk} dy_j \otimes dy_k),$$

which is positive definite. \square

Using this complex structure, we can also define a holomorphic volume form Ω_g on X by

$$\Omega_g := dz_1 \wedge \cdots \wedge dz_n$$

when B is in fact special. Then (X, ω_{\hbar}, J_g) is also a Calabi-Yau manifold.

Similarly, we have the following

Proposition 2.5. *Given a Hessian Riemannian metric on B , there is a natural symplectic structure $\check{\omega}_g$ on \check{X} such that $(\check{X}, \check{\omega}_g, \check{J}_{\hbar})$ is a Kähler manifold.*

Proof. We define

$$\check{\omega}_g := \sum_{j,k=1}^n g_{jk} d\check{y}^j \wedge dx^k.$$

Then $\check{\omega}_g, \check{J}_{\hbar}$ is compatible. \square

As a conclusion, by fixing a Hessian type metric g on B , we obtain two Calabi-Yau manifolds (X, ω_{\hbar}, J_g) and $(\check{X}, \check{\omega}_g, \check{J}_{\hbar})$.

Remark 2.6. *Both the Kähler metrics $\omega_{\hbar}, \check{\omega}_g$ are Calabi-Yau metrics if and only if the local convex function ϕ satisfies the real Monge-Ampère equation:*

$$\det(\text{Hess}(\phi)) = \text{constant}.$$

2.2. SYZ transform of branes. We follow [2, 19] to define the SYZ transform of a Lagrangian section in a semi-flat Lagrangian torus fibration.

Let $\mathcal{P} \rightarrow X \times_B \check{X}$ be the Poincaré line bundle. The total space is defined as the quotient

$$\mathcal{P} := (T^*B \oplus TB) \times \mathbb{C}/\Lambda^* \oplus \Lambda.$$

The fiberwise action of $\Lambda^* \oplus \Lambda$ on $(T^*B \oplus TB) \times \mathbb{C}$ is given by

$$(\lambda, \check{\lambda}) \cdot (y, \check{y}, t) := (y + \lambda, \check{y} + \check{\lambda}, e^{i\pi((y, \check{\lambda}) - (\lambda, \check{y}))} \cdot t).$$

Define a connection $\nabla_{\mathcal{P}}$ on \mathcal{P} by

$$\nabla_{\mathcal{P}} := d + i\pi((y, d\check{y}) - (\check{y}, dy)).$$

The section $e^{i\pi(y, \check{y})}$ is invariant under the $\{0\} \oplus \Lambda$ action:

$$(0, \check{\lambda}) \cdot (y, \check{y}, t) = (y, \check{y} + \check{\lambda}, e^{i\pi(y, \check{\lambda} + \check{y})}).$$

Hence it descends to a section $\check{\mathbb{I}}$ on $T^*B \times_B \check{X}$. With respect to this frame, the connection $\nabla_{\mathcal{P}}$ can be written as

$$\nabla_{\mathcal{P}} = d + 2\pi i(y, d\check{y}).$$

The remaining action of $\Lambda^* \oplus \{0\}$ then becomes

$$\lambda \cdot [(y, \check{y}, e^{i\pi(y, \check{y})})]_{\Lambda} = e^{-2\pi i(\lambda, \check{y})} [y + \lambda, \check{y}, e^{i\pi(y + \lambda, \check{y})}]_{\Lambda}.$$

Let $\mathbb{L} = (L, \xi)$ be a Lagrangian section and \mathcal{L} be a $U(1)$ -local system on L with holonomy $e^{2\pi i b}$, for some $b \in \mathbb{R}$. Define

$$\check{\mathbb{L}}_b := (\pi_{\check{X}})_*((\xi \times id_{\check{X}})^*(\mathcal{P}) \otimes \mathcal{L}).$$

where $\pi_{\check{X}} : L \times_B \check{X} \rightarrow \check{X}$ is the natural projection. The following proposition is standard.

Proposition 2.7. *The connection $\nabla_{\check{\mathbb{L}}_b}$ satisfies $(\nabla_{\check{\mathbb{L}}_b}^2)^{0,2} = 0$ if and only if $\xi : L \rightarrow X$ is a Lagrangian section.*

Hence $\check{\mathbb{L}}_b$ carries a natural holomorphic structure.

Definition 2.8. *$(\check{\mathbb{L}}_b, \nabla_{\check{\mathbb{L}}_b})$ is called the SYZ mirror bundle of the A-brane $(\mathbb{L}, \mathcal{L})$. For convenience, we just write the SYZ mirror of $(\mathbb{L}, \mathcal{L})$ as $\check{\mathbb{L}}_b$ for short.*

Remark 2.9. *It is known that every Lagrangian sections are graded Lagrangian immersions of the Calabi-Yau manifold $(X, \omega_{\hbar}, \Omega_g)$. For a proof, see [17].*

2.3. SYZ transform of morphisms. The exposition here follows Ma's PhD thesis [20].

Let $\mathbb{L}_1 = (B, \xi_1), \mathbb{L}_2 = (B, \xi_2)$ be two Lagrangian sections of the semi-flat fibration $p : X \rightarrow B$ and \check{L}_1, \check{L}_2 be the SYZ mirror bundles. Denote the images of ξ_1 and ξ_2 by L_1 and L_2 respectively. Let

$$P(L_1, L_2) := \coprod_{x \in B} \pi_1(p^{-1}(x); L_1, L_2)$$

be the fiberwise geodesic path space of (L_1, L_2) . It is a (possibly disconnected) covering space of B via the natural projection map.

Let $p \in L_1 \cap L_2$ and $x \in B$. Let $u : [0, 1] \times [0, 1] \rightarrow X$ be a smooth map such that for any $s, t \in [0, 1]$,

$$u(s, 0) \in L_1, \quad u(s, 1) \in L_2, \quad u(0, t) = p, \quad u(1, t) \in p^{-1}(x).$$

Let $\tilde{P}(L_1, L_2)$ be set of all $(x, [u])$, where $x \in B$ and $[u]$ is the homotopy class of some u that satisfies the above properties. There is a natural boundary map $\partial^+ : \tilde{P}(L_1, L_2) \rightarrow P(L_1, L_2)$ given by

$$\partial^+ : (x, [u]) \mapsto (x, [u|_{\{1\} \times [0, 1]}]).$$

Define the area function $\mathcal{A} : \tilde{P}(L_1, L_2) \rightarrow \mathbb{R}$ by

$$\mathcal{A}(x, [u]) = \int_0^1 \int_0^1 u^* \omega_{\hbar}.$$

Since the boundary of u lies in Lagrangian submanifolds, the integral depends only on the homotopy class of the map u . In general, \mathcal{A} may not be well-defined as a

function on $P(L_1, L_2)$ but its differential

$$d\mathcal{A}(x, \gamma) = \dot{\gamma} \lrcorner \omega_{\hbar} = \hbar^{-1} \sum_{j=1}^n \left(\xi_j^{(2)} - \xi_j^{(1)} + m_j \right) dx^j,$$

where $(m_1, \dots, m_n) \in \mathbb{Z}^n$ corresponds to $\gamma \in \pi_1(p^{-1}(x); L_1, L_2) \cong \mathbb{Z}^n$, is a well-defined 1-form on $P(L_1, L_2)$. That is, there exists a 1-form θ on $P(L_1, L_2)$ such that $(\partial^+)^*\theta = d\mathcal{A}$. By abuse of notation, we use $d\mathcal{A}$ to stand for the 1-form θ . But one should keep in mind that θ is, in general, not exact. Note that the critical points of \mathcal{A} are precisely the intersection points of L_1 and L_2 .

On $P(L_1, L_2)$, one can equip its space of differential forms $A^\bullet(P(L_1, L_2))$ with the Witten differential:

$$d_{W, \hbar} := d + 2\pi(\dot{\gamma} \lrcorner \omega_{\hbar}) \wedge.$$

Following [20] (see also [8, 17]), we define the SYZ transform $\mathcal{F} : A^\bullet(P(L_1, L_2)) \rightarrow A^{0, \bullet}(\check{X}, \check{L}_1^* \otimes \check{L}_2)$ by

$$\mathcal{F} \left(\sum_I \alpha_I(x, m) dx^I \right) := \left(-\frac{\hbar}{2i} \right)^{|\alpha|} \sum_{m \in \mathbb{Z}^n} \sum_I \alpha_I(x, m) e^{2\pi i(m, \check{y})} d\check{z}^I \otimes \check{1}_1^* \otimes \check{1}_2.$$

To avoid any convergence issues, we restrict ourselves to rapidly decay differential forms:

Definition 2.10. *An element $\alpha \in A^\bullet(P(L_1, L_2))$ is said to rapidly decay if for any compact subset $K \subset B$, integer $k, l \geq 0$,*

$$\lim_{|\dot{\gamma}| \rightarrow \infty} \sup_{x \in K} |\dot{\gamma}|^k |\nabla^l \alpha(x, \gamma)| = 0,$$

where ∇^l stands for the l -times covariant derivative with respect to the affine connection ∇ on $P(L_0, L)$ (induced from the one on B) and $|\cdot|$ is the norm induced by the Hessian-type metric g on B . The space of all rapidly decay forms is denoted by $A_{r.d.}^\bullet(P(L_1, L_2))$ where the subscript “r.d.” stands for “rapidly decay”.

It is easy to see that $d_{W, \hbar}$ preserves $A_{r.d.}^\bullet(P(L_1, L_2))$. We can also define the inverse SYZ transform:

$$\mathcal{F}^{-1} \left(\sum_I \check{\alpha}_I(x, \check{y}) d\check{z}^I \otimes \check{1}_1^* \otimes \check{1}_2 \right) := \left(-\frac{\hbar}{2i} \right)^{-|\check{\alpha}|} \sum_I \left(\int_{\check{p}^{-1}(x)} \check{\alpha}(x, \check{y}) e^{-2\pi i(m, \check{y})} d\check{y} \right) \otimes dx^I.$$

It is not hard to see that \mathcal{F} and \mathcal{F}^{-1} are indeed well-defined and inverse to each other. More importantly, \mathcal{F} and \mathcal{F}^{-1} exchange the Witten complex $(A_{r.d.}^\bullet(P(L_1, L_2)), d_{W, \hbar})$ with the Dolbeault complex $(A^{0, \bullet}(\check{X}, \check{L}_1^* \otimes \check{L}_2), \bar{\partial}_{\hbar})$.

Proposition 2.11. $\mathcal{F} : (A_{r.d.}^\bullet(P(L_1, L_2)), d_{W, \hbar}) \rightarrow (A^{0, \bullet}(\check{X}, \check{L}_1^* \otimes \check{L}_2), \bar{\partial}_{\hbar})$ is an isomorphism of differential graded vector spaces.

Proof. For $\alpha \in A_{r,d}^\bullet(P(L_1, L_2))$, we compute that

$$\begin{aligned}
 & \mathcal{F}(d_{W,\hbar}\alpha) \\
 &= \left(-\frac{\hbar}{2i}\right)^{|\alpha|+1} \sum_{m \in \mathbb{Z}^n} \sum_I \sum_{k=1}^n \left(\frac{\partial \alpha_I}{\partial x^k} + 2\pi\hbar^{-1} (\xi_1^k - \xi_2^k + m_k) \alpha_I \right) e^{2\pi i(m,\tilde{y})} d\bar{z}^k \wedge d\bar{z}^I \otimes \check{1}_1^* \otimes \check{1}_2 \\
 &= \left(-\frac{\hbar}{2i}\right)^{|\alpha|+1} \sum_{m \in \mathbb{Z}^n} \sum_I \sum_{k=1}^n \hbar^{-1} \left(\hbar \frac{\partial}{\partial x^k} \alpha_I e^{2\pi i(m,\tilde{y})} - i \frac{\partial}{\partial \tilde{y}^k} \alpha_I e^{2\pi i(m,\tilde{y})} \right) d\bar{z}^k \wedge d\bar{z}^I \otimes \check{1}_1^* \otimes \check{1}_2 \\
 &\quad + \left(-\frac{\hbar}{2i}\right)^{|\alpha|+1} \sum_{m \in \mathbb{Z}^n} \sum_I \sum_{k=1}^n 2\pi\hbar^{-1} (\xi_1^k - \xi_2^k) \alpha_I e^{2\pi i(m,\tilde{y})} d\bar{z}^k \wedge d\bar{z}^I \otimes \check{1}_1^* \otimes \check{1}_2 \\
 &= \left(\bar{\partial} + i\pi \sum_{k=1}^n (\xi_2^k - \xi_1^k) d\bar{z}^k \right) \mathcal{F}(\alpha) \\
 &= \bar{\partial}_\hbar \mathcal{F}(\alpha),
 \end{aligned}$$

as desired. \square

In order to connect (S^1 -valued) Morse theory of $(P(L_1, L_2), \mathcal{A})$ with the de Rham model $(A^\bullet(P(L_1, L_2)), d_W)$, we apply the idea of Witten deformations, or Witten-Morse theory [26].

For a Morse function $f : M \rightarrow \mathbb{R}$ defined on a closed Riemannian manifold (M, g) , Witten considered the twisted differential

$$d_{f,\hbar} := e^{-f/\hbar} de^{f/\hbar} = d + \hbar^{-1} df \wedge .$$

According to Witten-Morse theory, when $\hbar > 0$ is small enough, there is a 1-1 correspondence between index k critical points of f and small eigenforms of degree k , concentrated at p :

$$\text{Span}_{\mathbb{R}}(\text{Crit}^k(f)) \xrightarrow{\sim} \Omega_{sm}^k(M, f), \quad p \mapsto \alpha_{p,\hbar},$$

where

$$\Omega_{sm}^k(M, f) := \left\{ \alpha_\hbar \in \Omega^k(M) : \Delta_\hbar \alpha_\hbar = \lambda_\hbar \alpha_\hbar, \lim_{\hbar \rightarrow 0} \lambda_\hbar = 0 \right\},$$

and Δ_\hbar is the Witten Laplacian with respect to the metric g . The Morse cohomology of the pair (M, f) is defined to be the cohomology of

$$(\text{Span}_{\mathbb{R}}(\text{Crit}^\bullet(f)), \delta),$$

where δ is defined by counting gradient flow lines from index k critical points to index $k+1$ critical points. Witten also argued that the assignment

$$p \mapsto \alpha_{p,\hbar}$$

gives an isomorphism between the Morse cohomology and the cohomology of the Witten complex

$$(\Omega_{sm}^\bullet(M, f), d_{f,\hbar}),$$

which is nothing but the de Rham cohomology.

In Section 4.1, we will apply Witten-Morse theory to $(P(L_1, L_2), \mathcal{A})$ to solve the quantization problem.

Remark 2.12. *On the A-side, morphisms between two Lagrangian submanifolds should be given by their Floer complex. But here we are using a twisted version of the de Rham complex as the morphism space. Let us clarify their (conjectural) relations.*

As we mentioned in Remark 1.3, our Lagrangian sections L_1, L_2 are graded and the Maslov index of an intersection point between L_1 and L_2 coincides with the Morse index of the corresponding critical point of the area function \mathcal{A} by [17, Remark 13]. It was further pointed out in [17, Section 5.2] that the Morse complex of $(P(L_1, L_2), \mathcal{A})$ (together with higher products) is an approximation of the Floer complex $CF^\bullet(L_1, L_2)$ when \hbar is small (based on results of Floer [7] and Fukaya-Oh [9] in the case of cotangent bundles).

On the other hand, Witten's original work [26] together with the recent work of Chan, Leung and Ma [6] (see also [20]) proved that the Morse complex is also an approximation of the twisted de Rham complex of $(P(L_1, L_2), \mathcal{A})$ when \hbar is small.

In view of these results, we expect that the Floer complex $CF^\bullet(L_1, L_2)$ is quasi-isomorphic to the twisted de Rham complex of $(P(L_1, L_2), \mathcal{A})$. This is why we use the term "morphisms" in this section.

3. GEOMETRIC QUANTIZATION AND THE QUANTIZATION PROBLEM

In this section, we apply SYZ transforms to solve the quantization problem for Lagrangian torus fibrations over compact complete integral affine manifolds and projective toric manifolds.

3.1. Prequantum line bundles and polarizations. We first recall the precise statement of the quantization problem.

Definition 3.1. *A prequantum line bundle on a Kähler manifold (M, J, ω) is a pair (\mathcal{L}, ∇) , where \mathcal{L} is complex line bundle on M and ∇ is a complex integrable connection satisfying*

$$\frac{i}{2\pi} F_\nabla = \omega.$$

In particular, \mathcal{L} is an ample holomorphic line bundle.

Definition 3.2. *A polarization is an integrable Lagrangian subbundle $P \subset T_{\mathbb{C}}M$ meaning that $\text{rk}(P) = n$ and $[P, P] \subset P$ and $\omega|_P = 0$. The space of all polarized sections of \mathcal{L} is defined to be*

$$\Gamma_P(M, \mathcal{L}) := \{s \in \bar{\Gamma}_{sm}(M, \mathcal{L}) : \nabla_v s = 0 \text{ for all } v \in P\},$$

where $\bar{\Gamma}_{sm}(M, \mathcal{L})$ denotes the L^2 -completion of the space of smooth sections of \mathcal{L} .

There are two typical choices of polarization, namely, the real and complex polarizations.

Given a Lagrangian fibration $f : M \rightarrow B$ of M over an integral affine manifold B , we define

$$P_{\mathbb{R}} := \ker(f_* : TM \rightarrow f^*TB) \otimes \mathbb{C}.$$

Lemma 3.3. *$P_{\mathbb{R}}$ is a polarization.*

Proof. Since $f : M \rightarrow B$ is a Lagrangian fibration, $P_{\mathbb{R}}$ is of rank n and $\omega|_{P_{\mathbb{R}}} = 0$. Let (x^j) be coordinate of B and (y_j) be fiber coordinates. Sections of $P_{\mathbb{R}}$ is of form

$$\sum_{j=1}^n g^j(x, y) \frac{\partial}{\partial y_j},$$

and it follows that $[P_{\mathbb{R}}, P_{\mathbb{R}}] \subset P_{\mathbb{R}}$. \square

Definition 3.4. $P_{\mathbb{R}}$ is called the real polarization.

Remark 3.5. The terminology “real polarization” is justified by the property that $\bar{P}_{\mathbb{R}} = P_{\mathbb{R}}$.

Definition 3.6. A submanifold S of M is called Bohr-Sommerfeld if $\nabla|_S = d$.

Proposition 3.7 (Śniatycki [23]). There is a canonical identification

$$\Gamma_{P_{\mathbb{R}}}(M, \mathcal{L}) \cong \bigoplus_{F_x: BS \text{ fiber}} \mathbb{C} \cdot F_x.$$

Next, we define

$$P_{\mathbb{C}} := T^{0,1}M = \{v \in T_{\mathbb{C}}M : \check{J}(v) = -iv\}.$$

Lemma 3.8. $P_{\mathbb{C}}$ is a polarization.

Proof. Clearly $rk(P_{\mathbb{C}}) = n$. Since the complex structure of M is integrable, we have $[P_{\mathbb{C}}, P_{\mathbb{C}}] \subset P_{\mathbb{C}}$. Finally, ω is a Kähler form which is of $(1, 1)$ -type, so we have $\omega|_{P_{\mathbb{C}}} = 0$. \square

Definition 3.9. $P_{\mathbb{C}}$ is called the complex polarization.

The space of complex polarized sections is nothing but the space of holomorphic sections of \mathcal{L} :

$$\Gamma_{P_{\mathbb{C}}}(M, \mathcal{L}) = H^0(M, \mathcal{L}).$$

In geometric quantization, physicists ask whether a quantum theory is independent of the choice of polarizations. In other words, for two given polarizations P, P' , they are asking if the spaces of polarized sections are canonically isomorphic to each other. This is commonly known as the *quantization problem*. In particular, for a Kähler manifold admitting a Lagrangian fibration, we expect that the real and complex polarizations should give canonically isomorphic spaces of polarized sections.

4. MAIN RESULTS

We prove that, in the case of semi-flat Lagrangian torus fibrations over compact complete bases and projective toric manifolds, the spaces of real and complex polarized sections are canonically isomorphic via the SYZ transforms defined in Section 2. The proofs also use Witten-Morse theory on the fiberwise geodesic path spaces.

4.1. Semi-flat Lagrangian torus fibrations. Throughout this section, B is assumed to be an n -dimensional compact special integral affine manifold without boundary. We also assume $\hbar = 1$.

Choose a Hessian type metric g on B and let $\check{X} := TB/\Lambda$, which is equipped with the natural complex structure \check{J} . By Proposition 2.5, $(\check{X}, \check{J}, \check{\omega}_g)$ is a Kähler manifold (in fact a Calabi-Yau manifold).

Definition 4.1. A Hessian type metric g is said to be integral if there exists an open cover $\{U\}$ of B such that on each non-empty overlap $U \cap V$, the potential functions $\phi_U : U \rightarrow \mathbb{R}$, $\phi_V : V \rightarrow \mathbb{R}$ of g satisfy

$$\phi_U(x_U(x_V)) - \phi_V(x_V) = \langle m_{UV}, x_V \rangle + a_{UV},$$

for some $m_{UV} \in \mathbb{Z}^n$ and $a_{UV} \in \mathbb{R}$.

Let g be an integral Hessian type metric on B . Then we can define a Lagrangian section L_g of the dual fibration $p : X \rightarrow B$ by

$$L_g := \{(x, d\phi(x)) \in X : x \in B\}.$$

The integrality condition on g ensures that L_g is independent of the potential function ϕ . Hence it is a well-defined Lagrangian section of $p : X \rightarrow B$.

The SYZ transform \check{L}_g of L_g has connection

$$\nabla_{\check{L}_g} = d + 2\pi i \sum_{j=1}^n \frac{\partial \phi}{\partial x^j} d\check{y}_j.$$

Moreover, the curvature of $\nabla_{\check{L}_g}$ satisfies

$$\frac{i}{2\pi} F_{\nabla_{\check{L}_g}} = - \sum_{j,k} \frac{\partial^2 \phi}{\partial x^j \partial x^k} dx^j \wedge d\check{y}_k = \check{\omega}_g.$$

Hence we obtain the following

Proposition 4.2. *The SYZ transform $(\check{L}_g, \nabla_{\check{L}_g})$ of the Lagrangian section L_g is a prequantum line bundle on \check{X} .*

A key observation is given by the following

Proposition 4.3. *With respect to the prequantum line bundle $(\check{L}_g, \nabla_{\check{L}_g})$, there is a 1-1 correspondence between Bohr-Sommerfeld fibers of $\check{p} : \check{X} \rightarrow B$ and intersection points of L_0 and L_g .*

Proof. The connection $\nabla_{\check{L}_g}$ is trivial on a fiber \check{F}_x of $\check{p} : \check{X} \rightarrow B$ if and only if $d\phi(x) \in \Lambda_x^*$, that is, a point in $L_0 \cap L_g$. \square

In order to prove the next lemma, we need to introduce the following

Definition 4.4. *An affine manifold is said to be complete if its universal cover is diffeomorphic to \mathbb{R}^n as affine manifold.*

Remark 4.5. *There is a famous conjecture by Markus [21], stating that any compact special affine manifold is complete. Hence if Markus' conjecture is true, then the completeness condition is automatic. since we have assumed B to be a compact special integral affine manifold.*

Lemma 4.6. *Suppose the base B is complete. Then we have a decomposition*

$$P(L_0, L_g) = P \amalg \coprod_{p \in L_0 \cap L_g} P_p(L_0, L_g),$$

where $P_p(L_0, L_g)$ is the connected component of p . Moreover, each $P_p(L_0, L_g)$ is contractible.

Proof. Suppose $P_p(L_0, L_g) \cap P_q(L_0, L_g) \neq \emptyset$ with $p \neq q$. Since $P_p(L_0, L_g)$ and $P_q(L_0, L_g)$ are connected components, we have $P_p(L_0, L_g) = P_q(L_0, L_g)$. Since B is complete and $P_p(L_0, L_g)$ is a covering of B , the universal cover $\pi : \tilde{P}_p \rightarrow P_p(L_0, L_g)$ of $P_p(L_0, L_g)$ is affinely diffeomorphic to \mathbb{R}^n . Hence the connection ∇ pulls back to a connection on \tilde{P}_p which is gauge equivalent to the trivial connection d . Since any

two points in \mathbb{R}^n can be join by a geodesic with respect to the trivial connection, so is $P_p(L_0, L_g)$. Let $(l(s), \gamma_{l(s)})$ be a geodesic connecting p and q . Then

$$u(s, t) := (l(s), \gamma_{l(s)}(t))$$

is a disk in X with the properties

$$u(0, t) = p, \quad u(1, t) = q, \quad u(s, 0) \in L_0, \quad u(s, 1) \in L_g.$$

Consider the function $f : [0, 1] \rightarrow \mathbb{R}$ defined by

$$f(s) := g_X \left(\frac{\partial u}{\partial s}(s, 1), \frac{\partial u}{\partial t}(s, 1) \right).$$

In an affine chart, it can be written as

$$f(s) = \sum_{j=1}^n \dot{l}^j(s) \left(\frac{\partial \phi}{\partial x^j}(l(s)) + m_j \right).$$

Since $f(0) = f(1) = 0$, there exists $s_0 \in (0, 1)$ such that $f'(s_0) = 0$. In terms of the local expression of f , we have

$$\sum_{j=1}^n \ddot{l}^j(s_0) \left(\frac{\partial u}{\partial x^j}(l(s_0)) + m_j \right) + \sum_{j,k=1}^n \frac{\partial^2 \phi}{\partial x^j \partial x^k}(l(s_0)) \dot{l}^j(s_0) \dot{l}^k(s_0) = 0.$$

The path $l : [0, 1] \rightarrow B$ is also a geodesic with respect to the flat connection ∇ , so $\ddot{l}^j(s_0) = 0$ for all j . Hence

$$\sum_{j,k=1}^n \frac{\partial^2 \phi}{\partial x^j \partial x^k}(l(s_0)) \dot{l}^j(s_0) \dot{l}^k(s_0) = 0.$$

But $\dot{l}(s_0)$ is a nonzero vector, this contradicts positivity of $Hess(\phi)$. Therefore, $P_p(L_0, L_g) \cap P_q(L_0, L_g) = \emptyset$ whenever $p \neq q$.

Let $[\alpha]$ be a homotopy class of loops of $P_p(L_0, L_g)$ based at p . Let $\tilde{p} \in \tilde{P}_p$ be a lift of p . Suppose $[\alpha] \neq 0$. Choose any loop $\alpha \in [\alpha]$. By the path lifting property, there is a lift $\tilde{\alpha} : [0, 1] \rightarrow \tilde{P}_p$ of α such that $\tilde{\alpha}(0) = \tilde{p}$. Since $[\alpha] \neq 0$, $\tilde{p}' := \tilde{\alpha}(1) \neq \tilde{p}$. Choose a geodesic $c \subset \tilde{P}_p$ with starting point \tilde{p} and end point $\tilde{p}' \neq \tilde{p}$. Then $[\pi \circ c] = [\alpha]$. In particular, $\pi \circ c$ is a non-constant geodesic loop in $P_p(L_0, L_g)$, based at p . Let's write $\pi \circ c$ as

$$(\pi \circ c)(s) = (l(s), \gamma_{l(s)}).$$

Then

$$u(s, t) := (l(s), \gamma_{l(s)}(t))$$

is a disk in X satisfying

$$u(s, 0) \in L_0, \quad u(s, 1) \in L_g, \quad u(0, t) = p, \quad u(1, t) = p.$$

Again, we can consider the function

$$f(s) := g_X \left(\frac{\partial u}{\partial s}(s, 1), \frac{\partial u}{\partial t}(s, 1) \right).$$

Since $l(s)$ is an non-constant geodesic, the same proof above applies to conclude there exists some s_0 such that

$$Hess(\phi)(l(s_0))(\dot{l}(s_0), \dot{l}(s_0)) = 0.$$

Hence we must have $\alpha = 0$, that is, $P_p(L_0, L_g)$ is simply connected. In particular, $P_p(L_0, L_g) \cong \mathbb{R}^n$ as affine manifold. \square

In order to rule out the solution on P , we need the following

Lemma 4.7. *On a component of $P(L_0, L_g)$, any rapidly decay positive (resp. negative) functions have a maximum (resp. minimum).*

Proof. Let P_0 be a component of $P(L_0, L_g)$. Then P_0 is a covering of B . If P_0 is a finite covering of B , then we are done since B is compact. Suppose P_0 is an infinite covering of B . Then there is an infinite set $S \subset \mathbb{Z}^n$ such that P_0 is a S -covering of B . Now, since B is compact, we can choose a finite open covering $\mathcal{U} = \{U\}$ for B . We assume each U is evenly covered by open sets in P_0 . Let $\mathcal{V} = \{V\}$ be a pre-compact refinement of \mathcal{U} , that is, $\bar{V} \subset U$ and \bar{V} is compact. Then by identifying \bar{V} with a compact subset of P_0 , we have

$$P_0 = \bigcup_{m \in S} \bigcup_{V \in \mathcal{V}} \bar{V} \times \{m\}.$$

Let $f : P_0 \rightarrow \mathbb{R}$ be a rapidly decay positive function. Then for any $\epsilon > 0$ and $V \in \mathcal{V}$, there exists N_V such that for $|m| \geq N_V$,

$$\sup_{x \in \bar{V}} f(x, m) < \epsilon.$$

Let $N := \max_{V \in \mathcal{V}} N_V$. Then for $|m| \geq N$, the above inequality holds. Restricting f to the compact set

$$\bigcup_{|m| < N} \bigcup_{V \in \mathcal{V}} \bar{V} \times \{m\},$$

f achieves a maximum in it. Since f is bounded on P_0 , $\sup_{P_0} f$ exists and is positive. By taking ϵ smaller than $\sup_{P_0} f$, we conclude that f has a global maximum. \square

Let \mathcal{A}_p be the primitive of $d\mathcal{A}$ defined on the component $P_p(L_0, L_g)$ with $\mathcal{A}_p(p) = 0$, where $d\mathcal{A}(p) = 0$. Then we obtain

Proposition 4.8. *Suppose B is complete. Then the assignment*

$$A : p \mapsto e^{-2\pi\mathcal{A}_p}$$

gives an identification between $\bigoplus_{p \in L_0 \cap L_g} \mathbb{C} \cdot p$ and $\ker(d_W) \cap A_{r.d.}^0(P(L_0, L_g))$. In particular, by composing with the SYZ transform \mathcal{F} , we obtain the identification of morphism spaces:

$$\mathcal{F} \circ A : \bigoplus_{p \in L_0 \cap L_g} \mathbb{C} \cdot p \xrightarrow{\sim} H^0(\check{X}, \check{L}_g).$$

Proof. We claim that non-trivial rapidly decay solution to $d_W f = 0$ has a critical point and this implies $d_W f = 0$ has no non-trivial solution on P . To prove our claim, suppose f is a non-trivial rapidly decay solution defined on a component P_0 of P . By Lemma 4.7, f achieves its global maximum or minimum at some point $p_0 \in P_0$. Hence $df(p_0) = 0$ and so

$$f(p_0)d\mathcal{A}(p_0) = 0.$$

By replacing f by $-f$, we simply assume $f(p_0)$ is the maximum of f . If $f(p_0) \neq 0$, then $d\mathcal{A}(p_0) = 0$. If $f(p_0) = 0$, then $f \leq 0$. We can then find p_1 such that $f(p_1)$ is the global minimum of f . Hence $df(p_1) = 0$ and so

$$f(p_1)d\mathcal{A}(p_1) = 0.$$

Since f is non-trivial, $f(p_1) \neq 0$. Again, we have $d\mathcal{A}(p_1) = 0$. But P contains no critical point of \mathcal{A} . Therefore, f must be trivial.

Next, we prove that $e^{-2\pi\mathcal{A}_p}$ is rapidly decay. Note that the equation $d_W f = 0$ has a unique solution (up to a constant multiple) because $d\mathcal{A}$ is exact on $P_p(L_0, L_g)$. Hence for each critical point p of \mathcal{A} , the vector space $\ker(d_W) \cap A^0(P_p(L_0, L_g))$ is 1-dimensional. Now, for two critical point p, q , both $P_p(L_0, L_g)$ and $P_q(L_0, L_g)$ are universal covering spaces of B . Hence there exists a (non-unique) diffeomorphism $\psi_{qp} : P_p(L_0, L_g) \rightarrow P_q(L_0, L_g)$ covering the identity id_B . In particular, it preserves the 1-form $d\mathcal{A}$, that is,

$$\psi_{qp}^* d\mathcal{A}_q = d\mathcal{A}_p.$$

Hence $e^{-2\pi\mathcal{A}_q \circ \psi_{qp}}$ is a constant multiple of $e^{-2\pi\mathcal{A}_p}$. Now, the inverse SYZ transform and Lemma 4.6 give the identification:

$$\mathcal{F}^{-1} : H^0(\check{X}, \check{L}_g) \cong \ker(d_W) \cap A_{r,d}^0(P(L_0, L_g)) = \bigoplus_{p \in L_0 \cap L_g} \ker(d_W) \cap A_{r,d}^0(P_p(L_0, L_g)).$$

This implies $\dim(\ker(d_W) \cap A_{r,d}^0(P_p(L_0, L_g))) = 1$ for some p . As rapidly decay condition is preserved by any fiber preserving diffeomorphism, we have, for any q ,

$$e^{-2\pi\mathcal{A}_q} \propto e^{-2\pi\mathcal{A}_p \circ \psi_{qp}^{-1}}$$

is also rapidly decay.

Therefore, we conclude that A gives the isomorphism

$$\bigoplus_{p \in L_0 \cap L_g} \mathbb{C} \cdot p \cong \ker(d_W) \cap A_{r,d}^0(P(L_0, L_g)).$$

□

Therefore, the function $e^{-2\pi\mathcal{A}_p}$ in Proposition 4.8 is the Witten-deformation of the intersection point $p \in L_0 \cap L_g$.

Proposition 4.8 can be applied to abelian varieties:

Example 4.9. Let Ω be a positive definite symmetric $n \times n$ matrix with real entries. Let $\check{X} := \mathbb{C}^n / \mathbb{Z}^n \oplus i\Omega\mathbb{Z}^n$ be the abelian variety with period Ω . Complex coordinates are given by $z^j := \check{y}^j + ix^j$. The mirror of \check{X} is $X := \mathbb{R}^{2n} / \mathbb{Z}^{2n}$ equipped with the symplectic structure

$$\omega = \sum_{j,k=1}^n \Omega_{jk} dy_j \wedge dx^k.$$

Let Q be any positive definite symmetric integral matrix with the property that

$$Q\Omega = \Omega Q.$$

Then

$$L_Q := \{(x, Qx) \in X : x \in \mathbb{R}^n / \mathbb{Z}^n\}.$$

is a Lagrangian section of $p : X \rightarrow \mathbb{R}^n / \mathbb{Z}^n$. Let $x_1, \dots, x_N \in B$ such that $Qx_k \in \mathbb{Z}^n$ for all $k = 1, \dots, N$. In this case, the geodesic path space is a disjoint union of N

copies of \mathbb{R}^n and the Witten differential on a component $P_{(x_k, 0)}(L_0, L_Q)$ is given by

$$d + 2\pi \sum_{j,l,r=1}^n Q_{jl} (x^l - x_k^l + m_l) \Omega_{jr} dx^r.$$

Hence the function $A_k := A(x_k, 0)$ is given by

$$\begin{aligned} A_k(x, m) &= \exp \left(-2\pi \int_{x_k}^{x+m} \int_0^{Q(u-x_k)} \omega \right) \\ &= \exp \left(-\pi(x - x_k + m)^t Q \Omega (x - x_k + m) \right) \\ &= \exp \left(-\pi(x - x_k)^t Q \Omega (x - x_k) \right) \exp \left(-2\pi \left(m^t Q \Omega (x - x_k) + \frac{1}{2} m^t Q \Omega m \right) \right) \end{aligned}$$

Note that $Q\Omega$ is still positive definite since $[Q, \Omega] = 0$. The local holomorphic frame \check{e}_1 for \check{L}_Q is given by

$$\check{e}_1 = e^{-\pi(x-x_k)^t Q \Omega (x-x_k)} \check{1}_1.$$

The SYZ transform of A_k is given by

$$\begin{aligned} \mathcal{F}(A_k)(x, \check{y}) &= \sum_{m \in \mathbb{Z}^n} e^{-2\pi m^t Q \Omega m} e^{-2\pi(Q \Omega m, x - x_k)} e^{2\pi i(Q m, \check{y})} \otimes \check{e}_1 \\ &= \sum_{m \in \mathbb{Z}^n} e^{-2\pi m^t Q \Omega m} e^{2\pi i(m, Q(\check{y} + i\Omega(x - x_k)))} \otimes \check{e}_1 \end{aligned}$$

and its pullback via the covering map $\mathbb{C}^n \rightarrow \check{X}$ is the Riemann theta function

$$\theta \begin{bmatrix} 0 \\ -iQ\Omega x_k \end{bmatrix} (iQ\Omega, Qz_\Omega) = \sum_{m \in \mathbb{Z}^n} e^{-2\pi m^t Q \Omega m} e^{2\pi i(m, Qz_\Omega - iQ\Omega x_k)}$$

on \mathbb{C}^n , where $z_\Omega := \check{y} + i\Omega x$. In fact, the conditions $[Q, \Omega] = 0$ and $Q > 0$ are equivalent to the Riemann bilinear relations:

$$\begin{pmatrix} I_n & i\Omega \end{pmatrix} \begin{pmatrix} O_n & Q \\ -Q & O_n \end{pmatrix}^{-1} \begin{pmatrix} I_n \\ i\Omega \end{pmatrix} = 0, \quad -i \begin{pmatrix} I_n & i\Omega \end{pmatrix} \begin{pmatrix} O_n & Q \\ -Q & O_n \end{pmatrix}^{-1} \begin{pmatrix} I_n \\ -i\Omega \end{pmatrix} > 0,$$

which gives the ample line bundle \check{L}_Q on \check{X} and the space of global holomorphic sections of \check{L}_Q is generated by the theta functions $\mathcal{F}(A_k)$, $k = 1, \dots, N$.

Combining Propositions 3.7 and 4.8, we obtain our first main result:

Theorem 4.10. *Let $\check{p} : \check{X} \rightarrow B$ be a semi-flat Lagrangian torus fibration over a compact complete special integral affine manifold B . Let g be an integral Hessian type metric on B . With respect to the prequantum line bundle $(\check{L}_g, \nabla_{\check{L}_g})$, the SYZ transform \mathcal{F} induces a canonical isomorphism between the spaces of real and complex polarized sections.*

We expect that Proposition 4.8 and Theorem 4.10 also hold for non-compact but complete base B if we impose suitable growth conditions on the space of holomorphic sections $H^0(\check{X}, \check{L}_g)$, as in the following example.

Example 4.11. Let $B = \mathbb{R}^n$ and $\check{X} = (\mathbb{C}^\times)^n$. We have the natural torus fibration $\check{p} : \check{X} \rightarrow B$. The mirror of \check{X} is given by $X = (\mathbb{C}^\times)^n$. Consider the Lagrangian section

$$L_1 := \{(x, [x]) \in X : x \in B\}$$

of the dual fibration $p : X \rightarrow B$. The mirror line bundle is isomorphic to the trivial line bundle $\mathcal{O}_{\check{X}}$ since \check{X} is affine. The set of intersection points between L_1 and the zero section L_0 is given by

$$\{(k, 0) \in X : k \in \mathbb{Z}\}.$$

On the component $P_{(k,0)}(L_0, L_1)$, the function $A_k := A((k, 0))$ is given by

$$A_k(x) = \exp(-\pi(x - k)^2).$$

If we let $1 := \mathcal{F}(A_0)$, then the SYZ transform $\mathcal{F}(A_k)$ of A_k is proportional to monomial z^k on $\check{X} = (\mathbb{C}^\times)^n$. Hence if we restrict our attention to $H_{poly}^0(\check{X}, \mathcal{O}_{\check{X}})$, the space of all holomorphic functions on \check{X} that have polynomial growth (which are just Laurent polynomials), then we obtain the isomorphisms

$$\Gamma_{P_{\mathbb{R}}}(\check{X}, \check{L}_1) \cong \bigoplus_{p \in L_0 \cap L_1} \mathbb{C} \cdot p \cong H_{poly}^0(\check{X}, \mathcal{O}_{\check{X}}),$$

via Witten-Morse theory and the SYZ transform.

4.2. Projective toric manifolds. We now turn to the quantization problem for projective toric manifolds. Let us first recall some basic facts in toric geometry.

Let $N \cong \mathbb{Z}^n$ be a lattice of rank n and set

$$N_{\mathbb{R}} := N \otimes_{\mathbb{Z}} \mathbb{R}, \quad M := \text{Hom}_{\mathbb{Z}}(N, \mathbb{Z}), \quad M_{\mathbb{R}} := M \otimes_{\mathbb{Z}} \mathbb{R}.$$

Let Σ be a fan with primitive generators $v_1, \dots, v_d \in N$. Let $\check{X} := X_{\Sigma}$ be the toric variety associated to Σ . We assume \check{X} is smooth and projective. The Picard group of \check{X} has an explicit description as follows. Let $\iota : M \rightarrow \mathbb{Z}^d$ be given by

$$\iota : u \mapsto (\langle u, v_1 \rangle, \dots, \langle u, v_d \rangle).$$

The assignment

$$[a] \mapsto \mathcal{L}_{[a]} := \mathcal{O} \left(\sum_{j=1}^d a_j D_j \right),$$

where $D_j \subset \check{X}$ is the toric divisor corresponds to the ray v_j , gives the identification

$$\text{Pic}(\check{X}) \cong \mathbb{Z}^d / \iota(M).$$

Since \check{X} is assumed to be projective, we can take $\mathcal{L}_{[\lambda]}$ to be an ample line bundle on \check{X} . It is a well-known fact in toric geometry that such a line bundle is in fact very ample (see [10]). Hence it determines an embedding $i : \check{X} \hookrightarrow \mathbb{P}^N$ into some projective space \mathbb{P}^N . Let $\tilde{\omega} := i^* \omega_{FS}$, where ω_{FS} is the Fubini-Study metric on \mathbb{P}^N . The dense torus in \check{X} can be identified with $TN_{\mathbb{R}}/N$. Denote the coordinates on $N_{\mathbb{R}}$ by ξ^j and the induced complex coordinates $z^j := \check{y}^j + i\xi^j$ on $TN_{\mathbb{R}}/N$. It is well-known (see [14]) that on $TN_{\mathbb{R}}/N$, the Kähler form $\tilde{\omega}$ can be written as

$$\tilde{\omega} := 2i\partial\bar{\partial}\phi, \quad \phi(\xi) := \frac{1}{4\pi} \log \left(\sum_{u \in B \cap M} c_u e^{4\pi \langle u, \xi \rangle} \right).$$

Here $c_u \geq 0$ are constants that depend on the embedding $i : \check{X} \hookrightarrow \mathbb{P}^N$.

The $(\mathbb{C}^\times)^n$ -action on \check{X} restricted to a Hamiltonian T^n -action on $(\check{X}, \check{\omega})$ such that the moment map $\check{p} : \check{X} \rightarrow M_{\mathbb{R}}$ has image

$$B := \{x \in M_{\mathbb{R}} : \langle x, v_k \rangle + \lambda_j \geq 0, k = 1, \dots, d\},$$

which is a convex polytope in the vector space $M_{\mathbb{R}}$. The interior fibers of $\check{p} : \check{X} \rightarrow B$ are special Lagrangian tori with respect to the following holomorphic volume form

$$\check{\Omega} := \frac{dz^1}{z^1} \wedge \dots \wedge \frac{dz^n}{z^n},$$

which has a simple pole along the toric divisors. The fibers of \check{p} get degenerated to isotropic tori on the boundary ∂B . Moreover, the space of holomorphic sections of $\mathcal{L}_{[\lambda]}$ can be identified with the vector space spanned by the lattice points in B , that is $M \cap B$.

Remark 4.12. *One can identify $N_{\mathbb{R}}$ with \mathring{B} , the interior of B , via the differential $d\phi : N_{\mathbb{R}} \rightarrow \mathring{B}$. Moreover, if we let $\pi : TN_{\mathbb{R}}/N \rightarrow N_{\mathbb{R}}$ be the natural projection, the moment map \check{p} factors through $N_{\mathbb{R}}$ as $\check{p} = d\phi \circ \pi$.*

Before going into geometric quantization, let us recall mirror symmetry for a projective toric manifold. The mirror of \check{X} is given by the Landau-Ginzburg model (X, W) , where $X := T^*\mathring{B}/\Lambda^*$ and W is a holomorphic function on X , called the superpotential. Explicitly, we have

$$X = \{(z_1, \dots, z_n) \in (\mathbb{C}^\times)^n : |e^{-\lambda_j} z^{v_j}| < 1, \forall j = 1, \dots, d\}$$

and

$$W(z_1, \dots, z_n) := \sum_{j=1}^d e^{-\lambda_j} z^{v_j}.$$

Here, the complex coordinate z_j is given by

$$z_j := e^{2\pi i(y_j + i\phi_j)},$$

where $\phi_j = \frac{\partial \phi}{\partial \xi^j}$. Hence we can identify X with $T^*N_{\mathbb{R}}/M$ via

$$(y_1 + i\phi_1, y_2 + i\phi_2) \mapsto (y_1 + i\xi_1, y_2 + i\xi_2).$$

Hence X admits a special Lagrangian torus fibration over \mathring{B} with respect to the symplectic form

$$\sum_{j=1}^n dy_j \wedge d\xi^j$$

and the holomorphic volume form

$$\Omega = \frac{dz_1}{z_1} \wedge \dots \wedge \frac{dz_n}{z_n}.$$

In [5], the first named author of this paper defined a certain class of Lagrangian sections of the dual fibration $p : X \rightarrow \mathring{B}$ whose SYZ transform are in $\text{Pic}(\check{X})$. Moreover, every such line bundle carries a natural T^n -invariant Hermitian metric. Let's recall this class of objects. We identify \mathring{B} with $N_{\mathbb{R}}$ by $d\phi$.

Definition 4.13. *Let $a := (a_1, \dots, a_d) \in \mathbb{Z}^d$. A Lagrangian section \tilde{L} of $p : T^*N_{\mathbb{R}} \rightarrow N_{\mathbb{R}}$ is said to satisfy Condition $(*_a)$ if there is a C^2 -potential function $g : N_{\mathbb{R}} \rightarrow \mathbb{R}$ of \tilde{L} satisfies the following conditions: For any top dimensional cone $\sigma \in \Sigma(n)$, without loss of generality, assume σ is generated by v_1, \dots, v_n and let $\xi(t) = t_1 v_1 + \dots + t_n v_n$, for $t_j \in \mathbb{R}$, we have*

1. For any $j = 1, \dots, n$, the limits

$$\lim_{t_j \rightarrow -\infty} 2e^{-4\pi t_j} (\langle dg(\xi(t)), v_j \rangle + a_j),$$

$$\lim_{t_j \rightarrow -\infty} e^{-4\pi t_j} \text{Hess}(g(\xi(t)))(v_j, v_j)$$

exist and equal.

2. For any $j, k, l = 1, \dots, n$, the limit

$$\lim_{t_l \rightarrow -\infty} \text{Hess}(g(\xi(t)))(v_j, v_k)$$

exists.

3. For any $j, k = 1, \dots, n$ with $j \neq k$, we have

$$\lim_{t_j \rightarrow -\infty} e^{-2\pi(t_j+t_k)} \text{Hess}(g(\xi(t)))(v_j, v_k) = 0$$

or

$$\lim_{t_k \rightarrow -\infty} e^{-2\pi(t_j+t_k)} \text{Hess}(g(\xi(t)))(v_j, v_k) = 0.$$

Let $[a] \in \mathbb{Z}^d / \iota(M)$. A Lagrangian section of $p : X \rightarrow N_{\mathbb{R}}$ is said to satisfy Condition $(*_{[a]})$ if for some lift $\tilde{L} \subset T^*N_{\mathbb{R}}$ of L , \tilde{L} satisfies Condition $(*_a)$ for some representative a of $[a]$.

Remark 4.14. There is a slight difference between our definition and that given in [5]. We require all the exponential terms in the limits consist of a fact of 2π . This difference is due to our choice of complex coordinates on X being $e^{2\pi i(y_j + i\phi_j)}$, while in [5], the author used $e^{\phi_j + iy_j}$ instead.

The main result in [5] is the following

Theorem 4.15 (Chan [5]). Fix $a \in \mathbb{Z}^d$. The SYZ transform gives a 1-1 correspondence between Lagrangian sections satisfying Condition $(*_{[a]})$ and T^n -invariant Hermitian metrics on $\mathcal{L}_{[a]}$ of C^2 -class.

In particular, we know that whenever a Lagrangian section satisfies Condition $(*_{[a]})$, then its SYZ transform is isomorphic to the holomorphic line bundle $\mathcal{L}_{[a]}$.

Now we are ready to work on the quantization problem for the projective toric manifold \check{X} . As the fibers of the moment map \check{p} degenerates to isotropic tori on the boundary ∂B , it is natural consider Bohr-Sommerfeld isotropic submanifolds. It is already known that Bohr-Sommerfeld isotropic fibers are precisely those fibers above the lattice points of B (see [3]), and these lattice points also correspond to holomorphic sections of the associated line bundle of the polytope B , the equivalence of real and complex polarization is almost trivial. Here we give a mirror symmetric interpretation of this equivalence.

We consider the Lagrangian section

$$L_{\phi} := \{(\xi, d\phi(\xi)) \in X : \xi \in N_{\mathbb{R}}\}.$$

Applying the SYZ transform, we obtain a holomorphic line bundle \check{L}_{ϕ} on the dense torus $TN_{\mathbb{R}}/N \subset \check{X}$ together with a connection $\nabla_{\check{L}_{\phi}}$. The following proposition shows that the pair $(\check{L}_{\phi}, \nabla_{\check{L}_{\phi}})$ in fact extends to a holomorphic prequantum line bundle on \check{X} .

Proposition 4.16. *The line bundle \check{L}_ϕ extends to a holomorphic line bundle on \check{X} and is isomorphic to the line bundle $\mathcal{L}_{[\lambda]}$. Moreover, $(\check{L}_\phi, \nabla_{\check{L}_\phi})$ defines a pre-quantum line bundle on \check{X} .*

Proof. Fix a top dimensional cone $\sigma \in \Sigma$. By renaming, we assume σ is generated by v_1, \dots, v_n . Let $\xi(t) = \sum_{j=1}^n t_j v_j$. We show that

$$\lim_{t_j \rightarrow -\infty} \langle d\phi(\xi(t)), v_j \rangle = -\lambda_j.$$

We compute that

$$\langle d\phi(\xi(t)), v_j \rangle = \frac{\sum_{u \in B \cap M} \langle u, v_j \rangle c_u e^{4\pi \langle u, \xi \rangle}}{\sum_{u \in B \cap M} c_u e^{4\pi \langle u, \xi \rangle}}.$$

Then

$$\lim_{t_j \rightarrow -\infty} \langle d\phi(\xi(t)), v_j \rangle = \min_{u \in B \cap M} \langle u, v_j \rangle = -\lambda_j.$$

It is not hard to see that the Lagrangian L_ϕ satisfies Condition $(*_{[\lambda]})$. Hence by Theorem 4.15, $(\check{L}_\phi, \nabla_{\check{L}_\phi})$ extends to \check{X} and $\check{L}_\phi \cong \mathcal{L}_{[\lambda]}$ as holomorphic line bundle.

For the prequantum condition, recall that the connection $\nabla_{\check{L}_\phi}$ is given by

$$\nabla_{\check{L}_\phi} = d + 2\pi i \sum_{j=1}^n \frac{\partial \phi}{\partial \xi^j} d\check{y}^j.$$

Hence

$$\frac{i}{2\pi} F_{\nabla_{\check{L}_\phi}} = - \sum_{j=1}^n \frac{\partial^2 \phi}{\partial \xi^k \partial \xi^j} d\xi^k \wedge d\check{y}^j = \check{\omega}.$$

This completes the proof of the proposition. \square

As we have mentioned, the spaces of real and complex polarized sections are respectively given by

$$\begin{aligned} \Gamma_{P_{\mathbb{R}}}(\check{X}, \check{L}_\phi) &\cong \bigoplus_{\check{F}_x: \text{BS fiber}} \mathbb{C} \cdot \check{F}_x, \\ \Gamma_{P_{\mathbb{C}}}(\check{X}, \check{L}_\phi) &= H^0(\check{X}, \check{L}_\phi). \end{aligned}$$

Note that the direct sum is taken over all Bohr-Sommerfeld isotropic fibers.

As the semi-flat case, we would like to establish a correspondence between $L_0 \cap L_\phi$ and Bohr-Sommerfeld fibers. However, this is not true if we only consider interior intersections. What we need is to include the intersections at infinity:

Definition 4.17. *A Lagrangian section L of $p : X \rightarrow B$ is said to be intersecting the zero section L_0 at infinity if there exists a lift $\tilde{L} \subset T^*N_{\mathbb{R}}$ of L , a potential function $g : N_{\mathbb{R}} \rightarrow \mathbb{R}$ of \tilde{L} , a ray $\xi : [0, \infty) \rightarrow N_{\mathbb{R}}$ such that the limit*

$$\lim_{t \rightarrow \infty} dg(\xi(t))$$

exists in M . Fix a potential g . Let

$$\mathcal{R}_{\mathbb{Z}}(L) := \left\{ \xi : [0, \infty) \rightarrow N_{\mathbb{R}} \mid \lim_{t \rightarrow \infty} dg(\xi(t)) \text{ exists in } M \right\}.$$

Two rays $\xi_1, \xi_2 \in \mathcal{R}_{\mathbb{Z}}(L)$ are said to be integrally equivalent if

$$\lim_{t \rightarrow \infty} dg(\xi_1(t)) = \lim_{t \rightarrow \infty} dg(\xi_2(t)).$$

Denote the set of all equivalence classes of such rays by $\overline{L_0 \cap L}$.

Remark 4.18. *Interior intersection points are precisely those equivalence classes that can be represented (uniquely) by constant path.*

In view of Section 4.1, one would like to consider the fiberwise geodesic path space $P(L_0, L_\phi)$. However, there are not enough critical points if we only consider interior intersections. To recall the information coming from infinity, we consider

$$\mathcal{R}(L_\phi) := \left\{ \xi : [0, \infty) \rightarrow N_{\mathbb{R}} \mid \lim_{t \rightarrow \infty} d\phi(\xi(t)) \text{ exists in } M_{\mathbb{R}} \right\}.$$

Similar to Definition 4.17, we say two rays ξ_1, ξ_2 are equivalent if

$$\lim_{t \rightarrow \infty} dg(\xi_1(t)) = \lim_{t \rightarrow \infty} dg(\xi_2(t)).$$

Denote $\overline{N_{\mathbb{R}}} := \mathcal{R}(L_\phi) / \sim$ and $\overline{P(L_0, L_\phi)} := \overline{N_{\mathbb{R}}} \times M$.

Remark 4.19. *There is a natural identification between $\overline{N_{\mathbb{R}}}$ and the polytope B , given by*

$$[\xi] \mapsto \lim_{t \rightarrow \infty} d\phi(\xi(t)).$$

Moreover, $\overline{L_0 \cap L_\phi}$ is mapped to $M \cap B$, the set of lattice points of B . See also [10].

The original path space $P(L_0, L_\phi)$ can be identified with $N_{\mathbb{R}} \times M$ via

$$(\xi, s(d\phi(\xi) + m)) \mapsto (\xi, m).$$

Hence $P(L_0, L_\phi)$ sits inside $\overline{P(L_0, L_\phi)}$ naturally by

$$(\xi, m) \mapsto ([\xi], m),$$

where $[\xi]$ is the equivalence class of the constant path ξ . Note that the interior intersection points are of form

$$([d\phi^{-1}(u)], -u),$$

with $u \in \mathring{B} \cap M$. For the intersections at infinity, they are of the form

$$([\xi], -u),$$

where $\xi \in \mathcal{R}_{\mathbb{Z}}(L_\phi)$ and $u \in \partial B \cap M$ is nothing but the limit

$$\lim_{t \rightarrow \infty} d\phi(\xi(t)).$$

Next, we study the Morse theory on the fiberwise geodesic path space $P(L_0, L_\phi)$. However, not all functions on $P(L_0, L_\phi)$ can be transformed to global sections of \check{L}_ϕ .

Let us recall the the coordinate charts associated to a vertex of B . Let $V(B)$ be the vertex set of B . Consider the coordinate chart $\check{p}^{-1}(B_v)$, where

$$B_v := \{v\} \cup \mathring{B} \cup \bigcup_{F: \text{face contains } v} \mathring{F}$$

and $v \in V(B)$. Let

$$l_k(x) := \langle x, v_k \rangle + \lambda_k, \quad k = 1, \dots, d.$$

By smoothness, without loss of generality, we assume the vertex v is given by

$$l_1(x) = \dots = l_n(x) = 0.$$

Let $A = (A_{jk}) \in GL(n, \mathbb{Z})$ be the differential of the affine map

$$(x_1, \dots, x_n) \mapsto (l_1(x), \dots, l_n(x)).$$

The gluing map $F_v : TN_{\mathbb{R}}/N \rightarrow \check{p}^{-1}(B_v)$ of $\check{p}^{-1}(B_v) \cong \mathbb{C}^n$ to $TN_{\mathbb{R}}/N \cong (\mathbb{C}^\times)^n$ is given by

$$F_v : \left(e^{2\pi i(\check{y}^j + i\xi^j)} \right)_{j=1, \dots, n} \mapsto \left(e^{2\pi i \sum_{k=1}^n {}^t A^{jk}(\check{y}^k + i\xi^k)} \right)_{j=1, \dots, n},$$

where $({}^t A^{jk})$ is the transpose of the inverse of A . We write

$$\xi_v^j := \sum_{k=1}^n {}^t A^{jk} \xi^k, \quad \check{y}_v^j := \sum_{k=1}^n {}^t A^{jk} \check{y}^k.$$

The complex coordinates

$$Z_v^j := X_v^j + i\check{Y}_v^j := e^{2\pi i z_v^j} := e^{2\pi i(\check{y}_v^j + i\xi_v^j)}$$

extend to $\check{p}^{-1}(B_v) \cong \mathbb{C}^n$.

Recall the SYZ transform of a function $f : P(L_0, L) \rightarrow \mathbb{C}$ is given by

$$\mathcal{F}(f) = \sum_{m \in \mathbb{Z}^n} \alpha_I(\xi, m) e^{2\pi i(m, \check{y})} \otimes \check{1}_L.$$

On $\check{p}^{-1}(B_v)$, there is an unitary frame $\check{1}_v$ such that

$$\check{1}_v = e^{2\pi i(v, \check{y})} e^{2\pi i(\lambda, \check{y})} \check{1}_L,$$

where $v \in V(B)$ and $\lambda \in \mathbb{Z}^n$ corresponds to the difference between the choices of lift of L to $TN_{\mathbb{R}}$ on the two charts $\check{p}^{-1}(\check{B})$, $\check{p}^{-1}(\check{B}_v)$. In terms of the coordinates ξ_v^j, \check{y}_v^j , we have

$$\mathcal{F}(f) = \sum_{m \in \mathbb{Z}^n} f({}^t A \xi_v, A^{-1}m + \lambda) e^{2\pi i(m, \check{y}_v)} \otimes \check{1}_v.$$

Lemma 4.20. *Let $F : (\mathbb{C}^\times)^n \rightarrow \mathbb{C}$ be a smooth function. Equip \mathbb{C}^n with the standard flat metric \bar{g} . Then F extends to a smooth function on \mathbb{C}^n if and only if for any pre-compact open subset $U \subset \mathbb{C}^n$ and $j \in \mathbb{Z}_{\geq 0}$, the covariant derivatives $\nabla^j F$ are bounded on $U \cap (\mathbb{C}^\times)^n$.*

Proof. We cover \mathbb{C}^n by pre-compact open ball $\mathcal{U} := \{U\}$. Fix $U \in \mathcal{U}$. Since U is connected and the divisor $\bigcup_{j=1}^n \{z_j = 0\}$ is of real co-dimension 2, for each two points $z_0, z_1 \in U \cap (\mathbb{C}^\times)^n$, we can choose a regular path $\gamma : [0, 1] \rightarrow U \cap (\mathbb{C}^\times)^n$ joining z_0 to z_1 . By reparametrizing γ , we can assume γ has constant speed $|\dot{\gamma}|_{\bar{g}} = |z_1 - z_0|$, the distance between z_0 and z_1 . It follows from the mean value theorem that

$$|\nabla^j F(z_1) - \nabla^j F(z_0)|_{\bar{g}} \leq M_U |z_1 - z_0|,$$

for some constant $M_U > 0$, depending on the open ball U . This implies $\nabla^j F$ is uniformly continuous on $U \cap (\mathbb{C}^\times)^n$. Hence we obtain an unique continuous extension of $\nabla^j F$ to \bar{U} . If U, V are two overlapping pre-compact open sets in \mathbb{C}^n , by uniqueness, the extensions coincide on $\bar{U} \cap \bar{V}$. Hence $\nabla^j F$ extends to a global continuous section on \mathbb{C}^n . The converse is trivial. \square

Remark 4.21. *The result of Lemma 4.20 dose not depend on the chosen metric on \mathbb{C}^n . We use flat metrics only for the sake of convenience.*

For a function $F : p^{-1}(\check{B}_v) \cong (\mathbb{C}^\times)^n \rightarrow \mathbb{C}$, write

$$\nabla^k F = \sum_{|I|+|J|=k} a_{IJ}(\xi_v, \check{y}_v) d\xi_v^I \otimes d\check{y}_v^J + \sum_{|I|+|J|=k} b_{IJ}(\xi_v, \check{y}_v) d\check{y}_v^I \otimes d\xi_v^J,$$

where we put $b_{IJ} \equiv 0$ if $I = \emptyset$ or $J = \emptyset$. For each multi-sets I, J , let P_{IJ}, Q_{IJ} be the linear differential operator so that

$$\begin{aligned} P_{IJ}F &= a_{IJ}, \\ Q_{IJ}F &= b_{IJ}. \end{aligned}$$

In terms of the coordinates ξ_v^j, \check{y}_v^j , the flat metric \bar{g} reads

$$\bar{g} = 4\pi^2 \sum_{j=1}^n e^{-4\pi\xi_v^j} (d\xi_v^j \otimes d\xi_v^j + d\check{y}_v^j \otimes d\check{y}_v^j)$$

and the only non-zero Christoffel symbols are $\Gamma_{jj}^j = -2\pi$, which are constants. We see that the coefficients of P_{IJ}, Q_{IJ} are in fact constants. Let's denote $\hat{P}_{IJ}, \hat{Q}_{IJ}$ the differential operator on $N_{\mathbb{R}} \times \{m\}$ by replacing $\frac{\partial}{\partial \check{y}_v^j}$ by multiplication by $2\pi i m_j$.

Definition 4.22. *Let $L \subset X$ be a Lagrangian section that satisfies Condition $(*_{[a]})$ for some $a \in \mathbb{Z}^d$. A function $f \in A_{r,d}^0(P(L_0, L))$ is said to be weakly extendable if for any vertex $v \in V(B)$, multi-sets I, J and pre-compact open subset $U \subset B_v$, there exists constant $M_{I,J,U} > 0$ such that*

$$\begin{aligned} \left| \hat{P}_{IJ}f(tA\xi_v, A^{-1}m + \lambda) \right| &\leq M_{I,J,U} \prod_{i \in I} e^{-2\pi\xi_v^i} \prod_{j \in J} e^{-2\pi\xi_v^j}, \\ \left| \hat{Q}_{IJ}f(tA\xi_v, A^{-1}m + \lambda) \right| &\leq M_{I,J,U} \prod_{i \in I} e^{-2\pi\xi_v^i} \prod_{j \in J} e^{-2\pi\xi_v^j}, \end{aligned}$$

for all $m \in M$, $\xi_v \in d\phi^{-1}(U)$ and $\check{y}_v \in \check{T}^m$.

In particular, if f is weakly extendable, we see that it is bounded on $N_{\mathbb{R}} = d\phi^{-1}(\check{B})$. The terminology ‘‘weakly extendable’’ is justified by the following

Lemma 4.23. *If $\mathcal{F}(f)$ extends to a smooth section of \check{L} on \check{X} , then f is weakly extendable.*

Proof. Since $\check{\Gamma}_v$ is a section defined on $\check{p}^{-1}(B_v)$, in order to obtain a global section of \check{L} , the functions

$$F(\xi_v, \check{y}_v) := \sum_{m \in \mathbb{Z}^n} f_m(\xi_v) e^{2\pi i(m, \check{y}_v)} := \sum_{m \in \mathbb{Z}^n} f(tA\xi_v, A^{-1}m + \lambda) e^{2\pi i(m, \check{y}_v)}$$

need to be extended to smooth functions defined on $\check{p}^{-1}(B_v)$. Equivalently, by Lemma 4.20, for every pre-compact subset $U \subset \mathbb{C}^n$, they have bounded covariant derivatives on $U \cap (\mathbb{C}^\times)^n$.

We have

$$\begin{aligned} |\nabla^k F|_{\bar{g}}^2 &= \sum_{|I|+|J|=k} (|P_{IJ}F|^2 + |Q_{IJ}F|^2) \prod_{i \in I} e^{4\pi\xi_v^i} \prod_{j \in J} e^{4\pi\xi_v^j}, \\ P_{IJ}F &= \sum_{m \in \mathbb{Z}^n} \hat{P}_{IJ}f_m(\xi_v) e^{2\pi i(m, \check{y}_v)}, \\ Q_{IJ}F &= \sum_{m \in \mathbb{Z}^n} \hat{Q}_{IJ}f_m(\xi_v) e^{2\pi i(m, \check{y}_v)}. \end{aligned}$$

Now, suppose $|\nabla^k F|_{\bar{g}} \leq M_U$ for some constant M_U . Then

$$\begin{aligned} |P_{IJ}F| &\leq M_{IJ,U} \prod_{i \in I} e^{-2\pi\xi_v^i} \prod_{j \in J} e^{-2\pi\xi_v^j}, \\ |Q_{IJ}F| &\leq M_{IJ,U} \prod_{i \in I} e^{-2\pi\xi_v^i} \prod_{j \in J} e^{-2\pi\xi_v^j}, \end{aligned}$$

for some constants $M_{IJ,U} > 0$. Hence by absorbing the constants, we have

$$\begin{aligned} |\hat{P}_{IJ}f_m| &\leq \int_{\check{T}^n} |P_{IJ}F| d\check{y} \leq M_{IJ,U} \prod_{i \in I} e^{-2\pi\xi_v^i} \prod_{j \in J} e^{-2\pi\xi_v^j}, \\ |\hat{Q}_{IJ}f_m| &\leq \int_{\check{T}^n} |Q_{IJ}F| d\check{y} \leq M_{IJ,U} \prod_{i \in I} e^{-2\pi\xi_v^i} \prod_{j \in J} e^{-2\pi\xi_v^j}. \end{aligned}$$

Thus, f is weakly extendable. \square

Remark 4.24. *We expect there should be a stronger condition on f so that $\mathcal{F}(f)$ extends to a smooth section of \check{L} .*

Now, we return to the Lagrangian section L_ϕ . Lemma 4.23 shows that not all functions on $P(L_0, L_\phi)$ can be transformed into global sections of \check{L}_ϕ . Next, we deal with holomorphic sections. Via the inverse SYZ transform, we need to consider the solution to $d_W f = 0$. We have the following

Lemma 4.25. *$f : N_{\mathbb{R}} \times \{m\} \rightarrow \mathbb{R}$ is an non-trivial weakly extendable solution to $d_W f = 0$ on the component $N_{\mathbb{R}} \times \{m\}$ of $P(L_0, L_\phi) \cong N_{\mathbb{R}} \times M$ if and only if f is proportional to*

$$f_u(\xi) = e^{-2\pi\phi(\xi)} e^{2\pi\langle u, \xi \rangle}$$

where $u \in M \cap B$.

Proof. On each component $P_m := N_{\mathbb{R}} \times \{m\}$ of $P(L_0, L_\phi)$, solution to $d_W f = 0$ is proportional to

$$f_m(\xi) = e^{-2\pi\phi(\xi)} e^{2\pi\langle m, \xi \rangle}.$$

We claim that $m \in M \cap B$ if and only if f_m weakly extendable.

For $u \in B \cap M$, we apply the SYZ transform to f_u to obtain

$$\mathcal{F}(f_u)(z) = e^{-2\pi i\langle u, z \rangle} \otimes \check{e}_\phi,$$

which is nothing but the character corresponds to u . It is well-known in toric geometry that $\mathcal{F}(f_u)$ extends to a holomorphic section of $\mathcal{L}_{[\lambda]} \cong \check{L}_\phi$. Hence f_u is weakly extendable by Lemma 4.23.

Conversely, note that $\phi(\xi) - \langle m, \xi \rangle$ is convex, we have

$$\phi(\xi) - \langle m, \xi \rangle \geq \phi(\xi') - \langle m, \xi' \rangle + \langle d\phi(\xi') - m, \xi - \xi' \rangle,$$

for any $\xi, \xi' \in N_{\mathbb{R}}$. As $m \notin M \cap B$, there exists $k \in \{1, \dots, d\}$ such that

$$\langle m, v_k \rangle + \lambda_k < 0.$$

Hence

$$\langle d\phi(\xi') - m, v_k \rangle > 0.$$

for all $\xi' \in N_{\mathbb{R}}$. In terms of the polytope coordinates, this means

$$\langle x - m, v_k \rangle > 0,$$

for all $x \in B$. Since B is compact, there exists $\delta > 0$ such that

$$\langle x - m, v_k \rangle > \delta,$$

for all $x \in B$. Equivalently,

$$\langle d\phi(\xi') - m, v_k \rangle > \delta,$$

for all $\xi' \in N_{\mathbb{R}}$. Putting $\xi = 0$ and $\xi' = tv_k$, for $t > 0$, we get

$$\phi(tv_k) - \langle m, tv_k \rangle < \phi(0) - t\delta.$$

Hence

$$\lim_{t \rightarrow \infty} (\phi(tv_k) - \langle m, tv_k \rangle) = -\infty$$

and hence f_m is unbounded in the v_k -direction. Therefore, f_m cannot be weakly extendable for $m \notin B \cap M$. \square

As before, the function f_u in Lemma 4.25 is regarded as the Witten-deformation of the intersection point $[\xi(t)] \in \overline{L_0} \cap \overline{L_\phi}$, with

$$\lim_{t \rightarrow \infty} \xi(t) = u.$$

We are now ready to give another proof of the quantization problem via SYZ transforms; the original proof can be found in [3].

Theorem 4.26. *Let $\check{p}: \check{X} \rightarrow B$ be the moment map of the projective toric manifold \check{X} . With respect to the prequantum line bundle $(\check{L}_\phi, \nabla_{\check{L}_\phi})$, the SYZ transform \mathcal{F} induces a canonical isomorphism between the space of real and complex polarized sections.*

Proof. In terms of the coordinates $(\xi, \check{y}) \in TN_{\mathbb{R}}/N$, the connection $\nabla_{\check{L}_\phi}$ is given by

$$\nabla_{\check{L}_\phi} = d + 2\pi i \sum_{j=1}^n \frac{\partial \phi}{\partial \xi^j} d\check{y}^j.$$

Suppose $\xi(t)$ is a ray so that

$$\lim_{t \rightarrow \infty} d\phi(\xi(t)) = u,$$

for some $u \in M$ ($u \in M \cap B$ in fact). If $u \in \mathring{B}$, then

$$\lim_{t \rightarrow \infty} \nabla_{\check{L}_\phi}|_{\check{F}_{\xi(t)}} = d + 2\pi i \sum_{j=1}^n u_j d\check{y}^j,$$

which is equivalent to the trivial connection on the Lagrangian fiber \check{F}_u . If $u \in \partial B$, we choose $v \in V(B)$ such that $u \in B_v$. Since $u \in \partial B_v$, there exists subset $S_u \subset \{1, \dots, n\}$ such that the fiber \check{F}_u sits inside the divisor

$$\left\{ \prod_{j \in S_u} Z_j = 0 \right\} \subset \mathbb{C}^n \cong \check{p}^{-1}(B_v).$$

Recall we have a frame $\check{\mathbb{I}}_v$ on $\check{p}^{-1}(B_v)$ and coordinates ξ_v^j, \check{y}_v^j on $\check{p}^{-1}(B_v)$ such that

$$\begin{aligned}\check{\mathbb{I}}_v &= e^{2\pi i \langle v, \check{y} \rangle} e^{2\pi i \langle \lambda, \check{y} \rangle} \check{\mathbb{I}}_\phi, \\ \xi_v^j &= \sum_{k=1}^n {}^t A^{jk} \xi^k, \\ \check{y}_v^j &= \sum_{k=1}^n {}^t A^{jk} \check{y}^k.\end{aligned}$$

Then with respect to $\check{\mathbb{I}}_v$ and the coordinates ξ_v^j, \check{y}_v^j , we have

$$\nabla_{\check{L}_\phi} = d + 2\pi i \sum_{j,k=1}^n A_{jk} \left(v_k + \lambda_k + \frac{\partial \phi}{\partial \xi^k} \right) d\check{y}_v^j.$$

Hence

$$\lim_{t \rightarrow \infty} \nabla_{\check{L}_\phi} |_{\check{F}_{\xi(t)}} = d + 2\pi i \sum_{j \notin S_u} (A(v + \lambda + u))_j d\check{y}_v^j.$$

Since $A(v + \lambda + u) \in \mathbb{Z}^n$, it is equivalent to the trivial connection on the isotropic fiber \check{F}_u . Hence we obtain the identification

$$\bigoplus_{\check{F}_x: \text{BS fiber}} \mathbb{C} \cdot \check{F}_x \cong \overline{L_0 \cap L_\phi}.$$

On a component of the fiberwise geodesic path space $P(L_0, L_\phi) \cong N_{\mathbb{R}} \times M$, by Lemma 4.25, bounded non-trivial solution to $d_W f = 0$ is proportional to

$$f_u(\xi) = e^{-2\pi \phi(\xi)} e^{2\pi \langle u, \xi \rangle},$$

which we can regard it as a function on $P(L_0, L_\phi)$ by setting

$$f_u(\xi, m) := \begin{cases} e^{-2\pi \phi(\xi)} e^{2\pi \langle u, \xi \rangle} & \text{if } m = -u \\ 0 & \text{if } m \neq -u \end{cases}$$

As we have seen in the proof of Lemma 4.25, the SYZ transform of f_u is given by the character

$$\mathcal{F}(f_u)(z) = e^{-2\pi i \langle u, z \rangle} \otimes \check{e}_\phi$$

corresponds to u , which extends to a global holomorphic section of $\mathcal{L}_{[\lambda]}$. Since $\mathcal{L}_{[\lambda]} \cong \check{L}_g$ as holomorphic line bundles by Lemma 4.16 and any holomorphic section of $\mathcal{L}_{[\lambda]}$ is a linear combination of these holomorphic sections, the result follows. \square

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