Cosine Transform Preconditioners for High Resolution Image Reconstruction

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Abstract

This paper studies the application of preconditioned conjugate gradient methods in highresolution image reconstruction problems- We consider reconstructing high resolution images from multiple undersampled shifted degraded frames with subpixel displacement errors-The resulting blurring matrices are spatially variant- The classical Tikhonov regularization and the Neumann boundary condition are used in the reconstruction process- The precon ditioners are derived by taking the cosine transform approximation of the blurring matrices. We prove that when the L_2 or H_1 horm regularization functional is used, the spectra or the $\overline{}$ preconditioned normal systems are clustered around 1 for sufficiently small subpixel displacement errors- Conjugate gradient methods will hence converge very quickly when applied to solving these preconditioned normal equations- Numerical examples are given to illustrate the fast convergence.

$\mathbf 1$ Introduction

of the distributions of the limitations of the experimental systems provided users with only multiple low resolutions images However-Committee in many resolutions-committee in many and the committee in the committee in the committee of resolution of the pictures of the ground taken from a satellite is relatively low and retrieving details on the ground becomes impossible Increasing the image resolution by using digital signal processing techniques - - - - - is therefore of great interest

The authors would like to dedicate this work to Prof- Robert Plemmons in celebration of his th birthday-

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We consider the reconstruction of a high resolution image from multiple undersampled, shifted- degraded and noisy images Multiple undersampled images are often obtained by us ing multiple identical image sensors shifted from each other by subpixel displacements The reconstruction of high resolution images can be modeled as solving

$$
\mathcal{H}f = g,\tag{1}
$$

where q is the observed from $\frac{1}{\sqrt{2}}$ resolution images-formed from the low resolution images-formed from the low resolution in $\frac{1}{\sqrt{2}}$ desired high resolution image and $\mathcal H$ is the reconstruction operator. If all the low resolution images are shifted from each other with exactly half-pixel displacements, $\mathcal H$ will be a spatially invariant operator However- displacement errors may be present in practice- and the resulting operator H becomes spatially variant.

 \mathbb{R} systems are illegalled and generally not positive denite-density not positive denite-denite-denite-denite-denite-denite-denite-denite-denite-denite-denite-denite-denite-denite-denite-denite-denite-denite-denite-d using a minimization and regularization technique

$$
\min_{f} \left\{ \|\mathcal{H}f - g\|_2^2 + \alpha \mathcal{R}(f) \right\}.
$$
 (2)

Here $\mathcal{R}(f)$ is a functional which measures the regularity of f and the regularization parameter is to control the degree of the solution \mathbb{R} is paper-degree of the solution \mathbb{R} is paperregularization functionals $\|f\|_2^2$ and $\|\mathcal{L}f\|_2^2$ where $\mathcal L$ is the first order differential operator.

Because of the blurring convolution process- the boundary values of g are not completely determined by the original image f inside the scene. They are also affected by the values of f outside the scene Thus in solving f from \mathbb{R} from - we need some assumptions on the values of from \mathbb{R} outside the scene These assumptions are called boundary conditions In - Bose and Boo used the traditional choice of imposing the zero boundary conditions outside the scenes α dark background outside the scene in the image reconstruction However-Construction However-Construction Howeveris not satisfactory in a ringing extending equation of the reconstructed by \mathbf{r} images. The problem is more severe if the images are reconstructed from a large sensor array since the number of pixel values of the image affected by the sensor array increases.

in the paper- will use the this condition of the Neumann problems in the image- the interpretation of the international that the scene immediately outside is a reflection of the original scene at the boundary. The Neumann boundary condition has been studied in image restoration - - and in image compression - Properties in have shown that the Neumann in have shown that the Neumann image model is the Neu gives better reconstructed high resolution images than that under the zero or periodic boundary conditions In - we also proposed to use cosine transform preconditioners to precondition the resulting linear systems and preliminary numerical results have shown that these preconditioners are effective. The main aim of this paper is to analyze the convergence rate of these systems. We prove that when the L original is used-the spectra of the spectra of preconditioned systems are clustered around 1 for sufficiently small displacement errors.

The outline of the paper is as follows In Section - we give a mathematical formulation of the problem A brief introduction on the cosine transform preconditioners and the convergence analysis will be given in Section - $\mathbf{1}$ the effectiveness of the cosine transform preconditioners.

$\overline{2}$ The Mathematical Model

We begin with a brief introduction of the mathematical model in high resolution image recon struction. Details can be found in $[4]$.

Consider a sensor array with $L_1 \times L_2$ sensors, each sensor has $N_1 \times N_2$ sensing elements (pixels) and the size of each sensing element is $T_1 \times T_1$. Our aim is to reconstruct an image of resolution $M_1 \times M_2$, where $M_1 = L_1 \times N_1$ and $M_2 = L_2 \times N_2$. To maintain the aspect ratio of the reconstructed image-where μ and μ and assume that L is an even number in the following discussion.

In order to have enough information to resolve the high resolution image- there are subpixel displacements between the sensors In the ideal case-ideal case-ideal case-ideal case-ideal case-ideal case-idea by a value proportional to $T_1/L \times T_2/L$. However, in practice there can be small perturbations around these ideal subpixel locations due to imperfection of the mechanical imaging system Thus, for $l_1, l_2 = 0, 1, \dots, L-1$ with $(l_1, l_2) \neq (0, 0)$, the horizontal and vertical displacements $d_{l_1l_2}^{\ast}$ and $d_{l_1l_2}^{\ast}$ of the $l_1 l_2$. The limit referred the limit reference sensor array with reference sensor array are the limit l_1 given by

$$
d_{l_1l_2}^x = \frac{T_1}{L}(l_1 + \epsilon_{l_1l_2}^x) \text{ and } d_{l_1l_2}^y = \frac{T_2}{L}(l_2 + \epsilon_{l_1l_2}^y).
$$

Here $\epsilon_{l_1l_2}^*$ and $\epsilon_{l_1l_2}^*$ denote respectively the normalized horizontal and vertical displacement errors.

We remark that the parameters $\epsilon_{l_1l_2}^x$ and $\epsilon_{l_1l_2}^y$ can $\frac{l_1 l_2}{\sqrt{2}}$ camera calibration. We assume that

$$
|\epsilon_{l_1l_2}^x|<\frac{1}{2}\quad\text{and}\quad |\epsilon_{l_1l_2}^y|<\frac{1}{2}.
$$

For if not- the low resolution images observed from two dierent sensor arrays will be overlapped so much that the reconstruction of the high resolution image is rendered impossible

 \mathcal{L} , the original scene Theorem is the observed low resolution in a set of the low resoluti sensor is modeled by

$$
g_{l_1l_2}[n_1, n_2] = \int_{T_2(n_2 - \frac{1}{2}) + d_{l_1l_2}^y}^{T_2(n_2 + \frac{1}{2}) + d_{l_1l_2}^y} \int_{T_1(n_1 - \frac{1}{2}) + d_{l_1l_2}^x}^{T_1(n_1 + \frac{1}{2}) + d_{l_1l_2}^x} f(x_1, x_2) dx_1 dx_2 + \eta_{l_1l_2}[n_1, n_2], \tag{3}
$$

 $f_{\mu_1 \nu_2}$, and is the noise corresponding to the line $f_{\mu_1 \nu_2}$, and the line $f_{\mu_1 \nu_2}$ and $f_{\mu_1 \nu_2}$, and the line $f_{\mu_1 \nu_2}$ sensor. We intersperse the low resolution images to form an $M_1 \times M_2$ image by assigning

$$
g[L(n_1-1)+l_1,L(n_2-1)+l_2]=g_{l_1l_2}[n_1,n_2].
$$
\n(4)

Here g is an $M_1 \times M_2$ image and is called the *observed high resolution image*. Figure 1 shows the method of forming a 4×4 image g with a 2×2 sensor array where each g_{ij} has a 2×2 יים עים ביצור בייתה והיים עים ביום וחוד והיים והיים ביום בין היים בין היים בין היים בין ה

Using a column by column ordering for q, we obtain $q = Hf + \eta$ where H is a spatially variant operator [4]. Since $\mathcal H$ is ill-conditioned due to the averaging of the pixel values in the

Figure 1: Construction of the observed high resolution image

image model in - the classical Tikhonov regularization is used and the minimization problem is solved In this paper- we use the regularization functionals

$$
\mathcal{R}(f) = ||f||_2^2 \quad \text{and} \quad \mathcal{R}(f) = ||\mathcal{L}f||_2^2 \tag{5}
$$

where $\mathcal L$ is the first order differential operator.

-Image Boundary

The continuous image model in (3) can be discretized by the rectangular rule and approximated by a discrete image model Because of the blurring process cf - the boundary values of \mathbf{u} are also also aected by the values of f outside the scene Thus in solving f from \mathbf{u} some assumptions on the values of f outside the scene In - Bose and Boo imposed the zero \mathbf{a} image reconstruction

Let g and f be respectively the discretization of g and f using a column by column ordering. Under the blurring condition- the blurring matrix corresponding to the light μ in the local corresponding to th can be written as

$$
\mathbf{H}_{l_1l_2}(\epsilon)=\mathbf{H}^x_{l_1l_2}(\epsilon)\otimes \mathbf{H}^y_{l_1l_2}(\epsilon)
$$

where $\mathbf{H}_{l_1 l_2}^{\omega}(\epsilon)$ is an $M_1 \times M_1$ banded Toeplitz matrix with bandwidth $2L-1$

$$
\tilde{\mathbf{H}}^x_{l_1 l_2}(\epsilon) = \frac{1}{L} \left(\begin{array}{cccc} 1 & \cdots & 1 & h^{x+}_{l_1 l_2} & 0 \\ \vdots & \ddots & \ddots & \ddots & \ddots \\ 1 & & & & & h^{x+}_{l_1 l_2} \\ h^{x-}_{l_1 l_2} & \ddots & \ddots & \ddots & 1 \\ & & \ddots & \ddots & \ddots & \vdots \\ 0 & & h^{x-}_{l_1 l_2} & 1 & \cdots & 1 \end{array} \right),
$$

and

$$
h^{x\pm}_{l_1l_2}=\frac{1}{2}\pm\epsilon^x_{l_1l_2}.
$$

The $M_2 \times M_2$ banded blurring matrix ${\bf H}^s_{l_1 l_2}(\epsilon)$ is defined similarly. We note that ringing effects will occur at the boundary of the reconstructed images if f is indeed not zero close to the boundary, see for instance Figure 3 in $\S 4$. The problem is more severe if the image is reconstructed from a large sensor array since the number of pixel values of the image affected by the sensor array increases

in a proposed to use the Neumann boundary condition on the Indian \mathcal{A} . It assumes that is assumed the scene immediately outside is a reflection of the original scene at the boundary. Our numerical results have shown that the Neumann boundary condition gives better reconstructed high resolution images than that by the zero or periodic boundary conditions Under the Neumann boundary condition, the blurring matrices are still banded matrices with bandwidth $2L = 1$, but there are entries added to the upper left part and the lower right part of the matrices (see the second matrix in (6)). The resulting matrices, denoted by ${\bf H}^s_{l_1l_2}(\epsilon)$ and ${\bf H}^s_{l_1l_2}(\epsilon)$, have a Toeplitz-plus-Hankel structure:

$$
\mathbf{H}_{l_1l_2}^{x}(\epsilon) = \frac{1}{L} \begin{pmatrix} 1 & \cdots & 1 & h_{l_1l_2}^{x+} & 0 \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ 1 & \ddots & \ddots & \ddots & \vdots \\ h_{l_1l_2}^{x-} & \ddots & \ddots & \ddots & \vdots \\ \vdots & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & 1 \end{pmatrix} + \frac{1}{L} \begin{pmatrix} 1 & \cdots & 1 & h_{l_1l_2}^{x-} & 0 \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ h_{l_1l_2}^{x-} & \ddots & \ddots & \ddots & \vdots \\ h_{l_1l_2}^{x+} & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots \\ 0 & \ddots & \ddots & \ddots & \ddots & \ddots \\ 0 &
$$

and $\mathbf{H}_{l_1l_2}^{\sigma}(\epsilon)$ is defined similarly. The blurring matrix corresponding to the (l_1,l_2) -th sensor \cdot . \cdot . under the Neumann boundary condition is given by

$$
\mathbf{H}_{l_1l_2}(\epsilon)=\mathbf{H}^x_{l_1l_2}(\epsilon)\otimes \mathbf{H}^y_{l_1l_2}(\epsilon).
$$

The blurring matrix for the whole sensor array is made up of blurring matrices from each

sensor

$$
\mathbf{H}_{L}(\epsilon) = \sum_{l_1=0}^{L-1} \sum_{l_2=0}^{L-1} \mathbf{D}_{l_1 l_2} \mathbf{H}_{l_1 l_2}(\epsilon).
$$
 (7)

 $H_{\rm eff}$ are diagonal matrices with diagonal elements equal to \pm if the corresponding compo \mathbf{r}_c and \mathbf{r}_c is the left \mathbf{r}_c and \mathbf{r}_c and \mathbf{r}_c and \mathbf{r}_c are details and \mathbf{r}_c and \mathbf{r}_c Tikhonov regularization- our discretization problem becomes

$$
(\mathbf{H}_L(\epsilon)^t \mathbf{H}_L(\epsilon) + \alpha \mathbf{R}) \mathbf{f} = \mathbf{H}_L(\epsilon)^t \mathbf{g}
$$
 (8)

where **R** is the discretization matrix corresponding to the regularization functional $\mathcal{R}(f)$ in (5).

Cosine Transform Based Preconditioners

The linear system (8) will be solved by using the preconditioned conjugate gradient method. in this section, we cosmic this section present the cosine transform precondition to HLV which exploits the co banded and block structures of the matrix

Let ${\bf C}_n$ be the $n\times n$ discrete cosine transform matrix, i.e., the (\imath,\jmath) -th entry of ${\bf C}_n$ is given by

$$
\sqrt{\frac{2-\delta_{i1}}{n}}\cos\left(\frac{(i-1)(2j-1)\pi}{2n}\right), \qquad 1\leq i, j\leq n,
$$

where δ_{ij} is the Kronecker delta. Note that the matrix-vector product $\mathbf{C}_n\mathbf{z}$ can be computed in $O(n \log n)$ operations for any vector **z**, see [13, pp. 59–60]. For an $m \times m$ block matrix **B** with the size of each block equal to $n \times n$, the cosine transform preconditioner $c(\mathbf{B})$ of \mathbf{B} is defined to be the matrix $(\mathbf{C}_m \otimes \mathbf{C}_n) \Lambda(\mathbf{C}_m \otimes \mathbf{C}_n)$ that minimizes

$$
||(\mathbf{C}_m \otimes \mathbf{C}_n) \Lambda(\mathbf{C}_m \otimes \mathbf{C}_n) - \mathbf{B}||_F
$$

in the Frobenius normal any $\vert \psi \vert$, which are the computing matrix ψ costs of cost of computing η $c(B)$ -y for any vector y is $O(mn \log mn)$ operations. For banded matrices in (i) , which have $(2L-1)^2$ non-zero diagonals and are of size $M_1M_2 \times M_1M_2$, the cost of constructing $c(\mathbf{H}_L(\epsilon))$ is of $O(L/M_1M_2)$ operations only, see $|U|.$

-Spatially Invariant Case

When there are no subpixel displacement errors, i.e., when all $\epsilon_{l_1, l_2}^* = \epsilon_{l_1, l_2}^* = 0$, then l_1, l_2 , l_1, l_2 ${\bf H}_{l_1l_2}^*(0)$ and also ${\bf H}_{l_1l_2}^*(0)$ are the same for all l_1 and $l_2.$ We will denote them simply by ${\bf H}_L^x$ and H_L^y . We claim that in this case, the blurring matrix $H_L \equiv H_L(0) = H_L^z \otimes H_L^y$ can always be diagonalized by the discrete cosine transform matrix

We begin with $L = 2$. The blurring matrix $H_2 = H_2^2 \otimes H_2^2$, where H_2^2 is an $M_1 \times M_1$ tridiagonal matrix given by

$$
\mathbf{H}_{2}^{x} = \frac{1}{2} \begin{pmatrix} \frac{3}{2} & \frac{1}{2} \\ \frac{1}{2} & 1 & \frac{1}{2} \\ \vdots & \vdots & \ddots & \vdots \\ \frac{1}{2} & 1 & \frac{1}{2} \\ \frac{1}{2} & \frac{1}{2} & \frac{3}{2} \end{pmatrix} = \frac{1}{2} \begin{pmatrix} 1 & \frac{1}{2} \\ \frac{1}{2} & 1 & \frac{1}{2} \\ \vdots & \vdots & \ddots & \vdots \\ \frac{1}{2} & 1 & \frac{1}{2} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ \frac{1}{2} & 1 & \frac{1}{2} \\ \frac{1}{2} & 1 & \frac{1}{2} \end{pmatrix} + \frac{1}{2} \begin{pmatrix} \frac{1}{2} & 0 & & & \\ 0 & 0 & 0 & & \\ & \ddots & \ddots & \ddots & \vdots \\ & & 0 & 0 & 0 \\ & & & 0 & \frac{1}{2} \end{pmatrix}
$$

and \mathbf{H}_2^s is an $M_2 \times M_2$ matrix with the same structure. It is easy to see that in this case, the matrices \mathbf{H}_2^s and \mathbf{H}_2^s can be diagonalized by \mathbf{C}_{M_1} and \mathbf{C}_{M_2} respectively, see the basis given in - for the class of matrices that can be diagonalized by the cosine transform matrix Thus \mathbf{H}_2 can be diagonalized by $\mathbf{C}_{M_1}\otimes \mathbf{C}_{M_2}.$

Next we observe that the blurring matrix is ill-conditioned.

 ${\bf L}$ emma $\bf l$ Under the Neumann boundary condition, the M $_1\times$ M $_1$ matrix ${\bf H}_2^{\rm s}$ can be diagonalized by the discrete cosine transform matrix and its eigenvalues are given by

$$
\lambda_j(\mathbf{H}_2^x) = \cos^2\left(\frac{(j-1)\pi}{2M_1}\right), \quad 1 \le j \le M_1.
$$
\n(9)

In particular, the condition number $\kappa(\textbf{h}_2)$ of the matrix \textbf{h}_2 satisfies

$$
\kappa(\mathbf{H}_2^x) \ge O(M_1^2). \tag{10}
$$

Proof The formula for the eigenvalues can be derived easily using the basis given in - for the class of matrices that can be diagonalized by the cosine transform matrix. Since $\lambda_{\rm max}(\mathbf{n}_{\bar{2}})=1$ and

$$
\lambda_{\min}(\mathbf{H}_2^x)=\cos^2\left(\frac{(M_1-1)\pi}{2M_1}\right)\leq \sin^2\left(\frac{\pi}{M_1}\right)\leq \frac{\pi^2}{M_1^2},
$$

the estimate of the condition number is then given by (10) . \Box

It follows from Lemma 1 that the condition number of the matrix $\mathbf{H}_2(=\mathbf{H}_2^u\otimes \mathbf{H}_2^y)$ is of $O(M_1^2M_2^2)$. The matrix is very in-conditioned, for $L > 2$, we have the following theorem.

Theorem I Under the Neumann boundary condition, the matrix \mathbf{H}_{L} can be alagonalized by the discrete cosine transform matrix and its eigenvalues are given by

$$
\lambda_i(\mathbf{H}_L^x) = \frac{4}{L} \cos^2\left(\frac{(i-1)\pi}{2M_1}\right) p_L\left(\frac{(i-1)\pi}{M_1}\right), \quad 1 \le i \le M_1,
$$
\n(11)

where

$$
p_L\left(\frac{(i-1)\pi}{M_1}\right) = \begin{cases} \sum_{j=1}^{L/4} \cos\left(\frac{(i-1)(2j-1)\pi}{M_1}\right), & L = 4k \text{ for some positive integer } k, \\ \frac{1}{2} + \sum_{j=1}^{(L-2)/4} \cos\left(\frac{(i-1)2j\pi}{M_1}\right), & \text{otherwise.} \end{cases}
$$
(12)

Proof: We first establish a relationship between the matrices \mathbf{n}_L and \mathbf{n}_2 . From (0), for $L \geq 2$, we have

$$
\mathbf{H}_{L}^{x} = \begin{cases} \frac{2}{L} \sum_{j=1}^{L/4} \mathbf{S}_{2j-1} \mathbf{H}_{2}^{x}, & L = 4k \text{ for some positive integer } k, \\ \frac{2}{L} \sum_{j=0}^{(L-2)/4} \mathbf{S}_{2j} \mathbf{H}_{2}^{x}, & \text{otherwise,} \end{cases}
$$
(13)

where \mathbf{S}_0 is the $M_1 \times M_1$ identity matrix and

$$
\mathbf{S}_k = \text{ Toeplitz}(\mathbf{e}_{k+1}) + \text{ Hankel}(\mathbf{e}_k), \quad 1 \le k \le M_1 - 1.
$$

Here Toeplitz(\mathbf{e}_k) is the $M_1 \times M_1$ symmetric Toeplitz matrix with the k-th unit vector \mathbf{e}_k as the first column, and Hankel(${\bf e}_k$) is the $M_1\times M_1$ Hankel matrix with ${\bf e}_k$ as the first column and ${\bf e}_k$ in the reverse order as the last column

We remark that the Toeplitz part in S_k can be interpreted as the decomposition of the discrete blurring function - -- -- - into the sum of the elementary discrete blurring function - - with dierent shifts For example- for L - we have

$$
[\frac{1}{2}, 1, 1, 1, \frac{1}{2}] = [\frac{1}{2}, 1, \frac{1}{2}, 0, 0] + [0, 0, \frac{1}{2}, 1, \frac{1}{2}],
$$

where the two terms on the right together gives the Toeplitz part in S- to \pm . All we have the Toeplitz part in S-

$$
[\frac{1}{2},1,1,1,1,1,\frac{1}{2}]=[\frac{1}{2},1,\frac{1}{2},0,0,0,0]+[0,0,\frac{1}{2},1,\frac{1}{2},0,0]+[0,0,0,0,\frac{1}{2},1,\frac{1}{2}],
$$

where the first and the third terms on the right together gives the Toeplitz part in S_2 while the middle term gives the Toeplitz part of S_0 . Because we are considering the Neumann boundary condition, entries outside the blurring matrix \mathbf{H}_L are hipped into the matrix (cf. (6)). This is done by means of the Hankel part of S_k . Thus the resulting shift matrices are given by S_k and

Since $\{S_k\}_{k=0}^{M_1-1}$ is exactly a basis for the space containing all matrices that can be diagonalized by $\mathbf{C}_{M_1},$ see [5], it follows that the matrix \mathbf{h}_L can be diagonalized by the discrete cosine transform matrix. We also note that the eigenvalues of S_k are given by

$$
\lambda_i(\mathbf{S}_k) = 2 \cos \left(\frac{(i-1)k\pi}{M_1} \right), \quad 1 \leq i \leq M_1,
$$

see for instance [5]. Using (13) and (9), the eigenvalues of \mathbf{H}_{L} are given in (11). Γ

Theorem 1 states that the matrices ${\bf H}_L (= {\bf H}_L^u \otimes {\bf H}_L^y)$ are also very ill-conditioned and their condition numbers are at least of order $M_1^2 M_2^2$ (cf. (11)). We remark that some of these matrices may even be singular. For instance, when $L = 4$ and $M_1 = M_2 = 64$, $\lambda_{33}(\mathbf{H}_4) = 0$. Thus a regularization procedure such as (8) should be imposed to obtain a reasonable estimate for the original image in the high resolution reconstruction

In this paper-box this paper-box in the L and H- and Hspondingly-be-condition-to-condition-to-condition-to-condition-to-condition-to-condition-to-condition-to-condition-

$$
(\mathbf{H}_{L}^{t}\mathbf{H}_{L} + \alpha \mathbf{I})\mathbf{f} = \mathbf{H}_{L}^{t}\mathbf{g} \quad \text{or} \quad (\mathbf{H}_{L}^{t}\mathbf{H}_{L} + \alpha \mathbf{L}^{t}\mathbf{L})\mathbf{f} = \mathbf{H}_{L}^{t}\mathbf{g},\tag{14}
$$

where $\alpha > 0$, **I** is the identity matrix and **L** L is the discrete Laplacian matrix with the Neumann boundary condition. We note that $\mathbf{L}^t\mathbf{L}$ can be diagonalized by the discrete cosine transform matrix-deformation instance instance instance instance instance in the Neumann boundary condition for both the blurring matrix \mathbf{n}_L and the regularization operator $\mathbf{L} \cdot \mathbf{L}$, then the coemcient matrix in (14) can be diagonalized by the discrete cosine transform matrix and hence its inversion can be done in three 2-dimensional fast cosine transforms (one for finding the eigenvalues of the coefficient matrix- two for transforming the right hand side and the solution vector- see for instance Thus the total cost of solving the system is of OM-Line Are Complete All systems is o

We remark that for the zero boundary condition- discrete sine transform matrices can di a gonalize Toephitz matrices with at most b bands (e.g., \mathbf{H}_{2}) but not dense Toephitz matrices in general (e.g., \mathbf{H}_{4}), see [3] for instance. Therefore, in general we have to solve large block-Toephiz-Toephiz-block systems. The fastest direct Toephiz solvers require $O(M_1^T M_2^T)$ operations- see The systems can also be solved by the preconditioned conjugate gradient method with some suitable preconditioners, see IVI, its matrix is at the cost per iteration is at the cost per iteration least four 2-dimensional fast Fourier transforms. Thus we see that the cost of using the Neumann boundary condition is significantly lower than that of using the zero boundary condition.

-Spatially Variant Case

When the blurring displacement are protocometers the blurring matrix matrix has the same banded of the same ba structure as the that of HL-1) with some clients showed it is a near block to a near block the first perturbed Toeplitz-block matrix but it can no longer be diagonalized by the cosine transform matrix. Therefore we solve the corresponding linear system by the preconditioned conjugate gradient method. We will use the cosine transform preconditioner $c(\mathbf{H}_L(\epsilon))$ of $\mathbf{H}_L(\epsilon)$ as the preconditioner

Below we study the convergence rate of the preconditioned conjugate gradient method for solving the linear systems

$$
[c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon)) + \alpha \mathbf{I}]^{-1} [\mathbf{H}_L(\epsilon)^t \mathbf{H}_L(\epsilon) + \alpha \mathbf{I}] \mathbf{f} = \mathbf{H}_L(\epsilon)^t \mathbf{g}
$$
(15)

and

$$
[c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon)) + \alpha \mathbf{L}^t \mathbf{L}]^{-1} [\mathbf{H}_L(\epsilon)^t \mathbf{H}_L(\epsilon) + \alpha \mathbf{L}^t \mathbf{L}] \mathbf{f} = \mathbf{H}_L(\epsilon)^t \mathbf{g},
$$
(16)

where α is a positive constant. We prove that the spectra of the preconditioned normal systems are clustered around 1 for sufficiently small subpixel displacement errors. Hence when the conju- α . The contract method is applied to solving the preconditioned systems α , and α , and α , α , α fast convergence. Our numerical results in $\S 4$ show that the cosine transform preconditioners can indeed speed up the convergence of the method. We begin the proof with the following lemma

Lemma 2 Let
$$
\epsilon^* = \max_{0 \leq l_1, l_2 \leq L-1} \{ \epsilon_{l_1 l_2}^x, \epsilon_{l_1 l_2}^y \}
$$
. Then for all M_1 and M_2 , we have $\|\mathbf{H}_L(\epsilon) - \mathbf{H}_L\|_2 \leq 4\epsilon^*$ and $\|c(\mathbf{H}_L(\epsilon)) - \mathbf{H}_L\|_2 \leq 4\epsilon^*$. (17)

 \blacksquare rown \blacksquare . The column of \blacksquare and \blacksquare and in at most LD character and each entry is bounded by ϵ /L. It follows that

$$
\|\mathbf{H}_{L}(\epsilon)-\mathbf{H}_{L}\|_{\infty}\leq 4\epsilon^* \quad \text{and} \quad \|\mathbf{H}_{L}(\epsilon)-\mathbf{H}_{L}\|_{1}\leq 4\epsilon^*.
$$

Hence the first inequality in (17) follows by using $\|\cdot\|_2 \leq \sqrt{\|\cdot\|_1 \|\cdot\|_\infty}$. For the second inequality, we have noted that by Theorem I, σ (\equiv μ) \equiv \equiv μ) include we have

$$
\|c(\mathbf{H}_L(\epsilon)) - \mathbf{H}_L\|_2 = \|c(\mathbf{H}_L(\epsilon) - \mathbf{H}_L)\|_2 \leq \|\mathbf{H}_L(\epsilon) - \mathbf{H}_L\|_2,
$$

where the last inequality follows from $||c(\cdot)||_2 \le ||\cdot||_2$, see [2].

 ${\bf Lemma ~3}~~ Let ~ \epsilon^* = \max_{0 \leq l_1,l_2 \leq L-1} \{ \epsilon^x_{l_1l_2}, \epsilon^y_{l_1l_2} \}.~~ Then$ \cdot . \cdot . ll-

$$
\|\mathbf{H}_L(\epsilon)^t\mathbf{H}_L(\epsilon) - c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon))\|_2 < d_L(\epsilon^*)
$$

where $a_L(\cdot)$ is a function independent of M₁ and M₂ and $\lim_{\epsilon^* \to 0} a_L(\epsilon^-) = 0$.

Proof We note that

$$
\begin{aligned}\n\|\mathbf{H}_{L}(\epsilon)^{t}\mathbf{H}_{L}(\epsilon) - c(\mathbf{H}_{L}(\epsilon))^{t} c(\mathbf{H}_{L}(\epsilon))\|_{2} \\
&\leq \|\mathbf{H}_{L}(\epsilon)^{t}[\mathbf{H}_{L}(\epsilon) - c(\mathbf{H}_{L}(\epsilon))] \|_{2} + \|[\mathbf{H}_{L}(\epsilon)^{t} - c(\mathbf{H}_{L}(\epsilon))^{t}] c(\mathbf{H}_{L}(\epsilon))\|_{2}.\n\end{aligned}
$$

By Theorem 1, $||H_L||_2$ is bounded above by a constant independent of M_1 and M_2 . Hence by $(17), \|H_L(\epsilon)^t\|_2$ and $\|c(H_L(\epsilon))\|_2$ are also bounded above by some constants independent of M_1 and M_2 . Moreover, by (17) again, $\|\mathbf{H}_L(\epsilon) - c(\mathbf{H}_L(\epsilon))\|_2$ and $\|\mathbf{H}_L(\epsilon)^t - c(\mathbf{H}_L(\epsilon))^t\|_2$ are less than $\delta \epsilon$. The result therefore follows. \blacksquare

Using the above lemmas- we can analyze the convergence rate of the preconditioned systems (15) and (16) .

Theorem 2 Let $\epsilon^* = \max_{0 \le l_1, l_2 \le L-1} \{\epsilon^x_{l_1l_2}, \epsilon^y_{l_1l_2}\}.$ If ϵ^* is sufficiently small, then the spectra of lll lthe preconditioned matrices

$$
[c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon))+\alpha \mathbf{I}]^{-1} [\mathbf{H}_L(\epsilon)^t \mathbf{H}_L(\epsilon)+\alpha \mathbf{I}]
$$

are clustered around 1 and their smallest eigenvalues are bounded away from 0 by a positive constant independent independent of M-G and M-

 \blacksquare . \blacksquare . The function of \blacksquare . The that is the three terms is the three terms is the term is the term in the term is the term in the term

$$
\Vert [c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon)) + \alpha \mathbf{I}]^{-1} \Vert_2 \leq \frac{1}{\alpha}
$$

and

$$
\begin{array}{ll} &&[c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon))+\alpha\mathbf{I}]^{-1}[\mathbf{H}_L(\epsilon)^t \mathbf{H}_L(\epsilon)+\alpha\mathbf{I}] \\ = & \mathbf{I}+[c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon))+\alpha\mathbf{I}]^{-1}[\mathbf{H}_L(\epsilon)^t \mathbf{H}_L(\epsilon)-c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon))].\end{array}
$$

Hence the result follows by applying Lemma 3. \Box

Theorem 3 Let $\epsilon^* = \max_{0 \leq l_1,l_2 \leq L-1} \{ \epsilon^x_{l_1l_2}, \epsilon^y_{l_1l_2} \}.$ Ij $\begin{array}{l} \mathbb{S}^y_{\{1_1\}} \} \ \ \ \textit{If ϵ^* is sufficiently small, then the spectra of} \end{array}$ the preconditioned matrices

$$
[c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon)) + \alpha \mathbf{L}^t \mathbf{L}]^{-1} [\mathbf{H}_L(\epsilon)^t \mathbf{H}_L(\epsilon) + \alpha \mathbf{L}^t \mathbf{L}]
$$

are clustered around 1 and their smallest eigenvalues are uniformly bounded away from 0 by a positive constant independent of M- and M

Proof Since

$$
\begin{array}{ll} &&[c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon)) + \alpha \mathbf{L}^t \mathbf{L}]^{-1} [\mathbf{H}_L(\epsilon)^t \mathbf{H}_L(\epsilon) + \alpha \mathbf{L}^t \mathbf{L}] \\ = & \mathbf{I} + [c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon)) + \alpha \mathbf{L}^t \mathbf{L}]^{-1} [\mathbf{H}_L(\epsilon)^t \mathbf{H}_L(\epsilon) - c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon))], \end{array}
$$

it suffices to show that $\|[c(\mathbf{H}_L(\epsilon))^t c(\mathbf{H}_L(\epsilon))+\alpha \mathbf{L}^t \mathbf{L}]^{-1}\|_2$ is bounded above by a constant independent of M_1 and M_2 . Since $\lambda_{\min}(A) + \lambda_{\min}(B) \leq \lambda_{\min}(A + B)$ for any Hermitian matrices A and B see B se

$$
\|\left[c(\mathbf{H}_{L}(\epsilon))^{t}c(\mathbf{H}_{L}(\epsilon)) + \alpha \mathbf{L}^{t}\mathbf{L}\right]^{-1}\|_{2}
$$
\n
$$
= \frac{1}{\lambda_{\min}(c(\mathbf{H}_{L}(\epsilon))^{t}c(\mathbf{H}_{L}(\epsilon)) + \alpha \mathbf{L}^{t}\mathbf{L})}
$$
\n
$$
\leq \frac{1}{\lambda_{\min}(\mathbf{H}_{L}^{t}\mathbf{H}_{L} + \alpha \mathbf{L}^{t}\mathbf{L}) + \lambda_{\min}(c(\mathbf{H}_{L}(\epsilon))^{t}c(\mathbf{H}_{L}(\epsilon)) - \mathbf{H}_{L}^{t}\mathbf{H}_{L})}
$$
\n
$$
\leq \frac{1}{\lambda_{\min}(\mathbf{H}_{L}^{t}\mathbf{H}_{L} + \alpha \mathbf{L}^{t}\mathbf{L}) - \|c(\mathbf{H}_{L}(\epsilon))^{t}c(\mathbf{H}_{L}(\epsilon)) - \mathbf{H}_{L}^{t}\mathbf{H}_{L}\|_{2}}.
$$
\n(18)

because the matrix $\mathbf{n}_L^{\intercal}\mathbf{n}_L + \alpha \mathbf{L}^{\intercal}\mathbf{L}$ can be diagonalized by the 2-dimensional discrete cosine transform matrix-definition matrix \mathbf{M}

$$
\lambda_{(i-1)M_2+j}(\mathbf{L}^t\mathbf{L}) = 4\sin^2\left(\frac{(i-1)\pi}{2M_1}\right) + 4\sin^2\left(\frac{(j-1)\pi}{2M_2}\right),\tag{19}
$$

for $1 \leq i \leq M_1$ and $1 \leq j \leq M_2$, see [5]. By using (11) and the fact that $H_L = H_L^x \otimes H_L^y$, we obtain

$$
\lambda_{(i-1)M_2+j}(\mathbf{H}_L^t \mathbf{H}_L) = \left(\frac{4}{L}\right)^4 \cos^4\left(\frac{(i-1)\pi}{2M_1}\right) \cos^4\left(\frac{(j-1)\pi}{2M_2}\right) p_L^2 \left(\frac{(i-1)\pi}{M_1}\right) p_L^2 \left(\frac{(j-1)\pi}{M_2}\right),\tag{20}
$$

for $1 \leq i \leq M_1$ and $1 \leq j \leq M_2$, where $p_L(\cdot)$ is defined in (12).

Clearly the function $\sin^{-}(x/z)$ is zero at $x = 0$ and positive in $(0, \pi)$, whereas the function $\cos^{4}(x/2)p_{L}^{2}(x) \geq 1/4$ at $x = 0$ and is nonnegative in $[0, \pi]$. Thus we see that the function

$$
4\alpha \sin^2\left(\frac{x}{2}\right) + 4\alpha \sin^2\left(\frac{y}{2}\right) + \left(\frac{4}{L}\right)^4 \cos^4\left(\frac{x}{2}\right) \cos^4\left(\frac{y}{2}\right) p_L^2(x) p_L^2(y)
$$

is positive for all x and y in $[0, \pi]$. It follows from (19) and (20) that the matrix $\mathbf{H}_L^T \mathbf{H}_L + \alpha \mathbf{L}^T \mathbf{L}$ is positive denite and its smallest eigenvalue is bounded away from by a positive constant independent of M-In view of In view of Lemma - the right for the right hand side of \mathcal{A} is the right dependent. by a positive constant independent of M-F \sim 1 m for μ and μ and M-F μ m-F μ \blacksquare

Thus we conclude that the preconditioned conjugate gradient method applied to (15) and (16) with a will converge superlinearly for superlinearly for such a such a see for such a second to see for instance $\vert \circ \vert$. Since $\mathbf{n}_L(\epsilon)$ has only $\vert 2L-1\rangle$ hon-zero diagonals, the matrix-vector product $\mathbf{n}_L(\epsilon)$ x can be done in $O(L^2 M_1 M_2)$. Thus the cost per each iteration is $O(M_1 M_2 \log M_1 M_2 + L^2 M_1 M_2)$. p Hence the total pixely extract the total cost for no meaning the image cost formula the second in of $O(M_1M_2$ tog $M_1M_2+L_1M_1M_2$ operations.

Numerical Examples

In this section- we illustrate the eectiveness of using cosine transform preconditioners for solving high resolution image reconstruction problems. The original image is shown in Figure 2 (left). The conjugate gradient method is employed to solving the preconditioned systems (15) and (16) . The stopping criteria is $||{\bf r}^{(j)}||_2/||{\bf r}^{(0)}||_2 < 10^{-6}$, where ${\bf r}^{(j)}$ is the normal equations residual after j iterations. In the tests, the parameters $\epsilon^*_{l_1 l_2}$ and $\epsilon^*_{l_1 l_2}$ are i l_1l_2 are random values chosen between $-1/2$ and signal to the signal to the low resolution of the low resolution of the low resolution of the low resolution images

Tables 1–2 show the numbers of iterations required for convergence for $L = 2$ and 4 respectively In the tables- cos- cir or no signify that the cosine transform preconditionerthe level-2 circulant preconditioner $\lbrack 8\rbrack$ or no preconditioner is used respectively. We see from the tables that the cosine transform preconditioner converges much faster than the circulant p is the matrix for different M and m and m and α and α and α is the size of the size of the reconstruction of the reconstruction of α image and α is the regularization parameter. Also the convergence rate is independent of M for fixed α as predicted by Theorems 2 and 3.

Next we show the 256×256 reconstructed images from four 128×128 low resolution images, i.e., a 2 \times 2 sensor array is used. One of the low resolution images is shown in Figure 2 (middle).

			α 1×10^{-2} 1×10^{-3} 1×10^{-4}					
					$M \vert \cos$ cir no $\vert \cos$ cir no $\vert \cos$ cir no			
					32 8 27 48 12 58 127 20 83 325			
					64 8 27 48 11 64 130 19 125 347			
					128 8 27 48 11 68 129 17 173 345			
					256 8 27 48 10 68 129 17 181 348			

Table 1a. Number of iterations for $L = 2$ where the L_2 norm regularization is used.

α 1×10^{-2} 1×10^{-3} 1×10^{-4}						
$M \vert \cos$ cir no cos cir no cos cir no						
32 7 16 26 9 38 68 13 70 178						
64 7 16 26 9 36 69 13 88 180						
128 7 16 26 9 38 69 13 99 180						
256 6 16 26 8 38 69 13 99 180						

Table b Number of iterations for L where the H- norm regularization is used

			α 1×10^{-2} 1×10^{-3} 1×10^{-4}					
							M \cos cir no \cos cir no \cos cir no	
							$\begin{array}{c cccccc} 32 & 7 & 33 & 45 & 10 & 67 & 111 & 16 & 145 & 256 \\ 64 & 6 & 34 & 47 & 10 & 84 & 123 & 16 & 180 & 314 \\ 128 & 6 & 32 & 47 & 10 & 96 & 125 & 15 & 237 & 323 \end{array}$	
							256 6 32 47 9 92 125 15 262 323	

Table 2a: Number of iterations for $L = 4$ where the L_2 norm regularization is used.

α 1×10^{-2} 1×10^{-3} 1×10^{-4}								
							$M \vert \cos$ cir no $\vert \cos$ cir no $\vert \cos$ cir no	
							$32 \begin{array}{ c c c c c } \hline 5 & 23 & 33 \end{array}$ 8 46 72 12 86 159	
							$\begin{array}{c ccccc} 64 & 5 & 23 & 33 \\ 128 & 5 & 23 & 34 \end{array}$ $\begin{array}{c ccccc} 8 & 63 & 83 & 12 & 127 & 182 \\ 7 & 65 & 87 & 11 & 155 & 204 \end{array}$	
							256 5 22 34 7 63 86 11 178 216	

Table b Number of iterations for L where the H- norm regularization is used

The observed high resolution image ^g is shown in Figure right We tried the Neumann- zero and periodic boundary conditions to reconstruct the high resolution images. Figure 3 shows the reconstructed images. The optimal regularization parameter α is chosen such that it minimizes $\frac{1}{2}$ relative error $\frac{1}{2}$ the reconstruction image $\frac{1}{2}$ (c) is the original image $\frac{1}{2}$ i.e., it minimizes $\|\mathbf{f} - \mathbf{f}_r(\alpha)\|_2 / \|\mathbf{f}\|_2$. By comparing the figures in Figure 3, it is clear that the trees in the image are much better reconstructed under the Neumann boundary condition than that under the zero and periodic boundary conditions. We also see that the boundary artifacts under the Neumann boundary condition are less prominent than that under the other two boundary conditions

Figure The original image left- a low resolution image middle- and the observed high resolution image (right).

Figure Reconstructed image using the Neumann boundary condition left- the zero boundary condition (middle) and the periodic boundary condition (right).

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