

Solutions to Final Exam

1. (a)

$$\int_{\mathbb{R}} S(t, x) dx = \int_{\mathbb{R}} \frac{1}{\sqrt{4\pi t}} e^{-x^2/4t} dx.$$

By change of variables and the Gaussian integral,

$$\int_{\mathbb{R}} S(t, x) dx = \int_{\mathbb{R}} \frac{1}{\sqrt{\pi}} e^{-x^2} dx = 1.$$

(b)

$$\max_{\delta < |x| < \infty} S(t, x) \leq \frac{1}{\sqrt{4\pi t}} e^{-\delta^2/4t}.$$

So

$$\max_{\delta < |x| < \infty} S(t, x) \rightarrow 0$$

as $t \rightarrow 0$.

2. Suppose that $f_n \rightarrow f$ uniformly, that is,

$$\sup_I |f_n - f| \rightarrow 0.$$

Since

$$\int_I |f_n - f|^2 \leq |I| \left(\sup_I |f_n - f| \right)^2,$$

$f_n \rightarrow f$ in the L^2 sense. Since

$$|f_n(x) - f(x)| \leq \sup_I |f_n - f|$$

for $x \in I$, $f_n \rightarrow f$ in the pointwise sense.

3. (a) For $(t, x) \in (0, \infty) \times (0, 1)$, since $0 \leq u \leq 1$ on the parabolic boundary $\overline{(0, 1) \times (0, t]} \setminus (0, 1) \times (0, t]$, by the maximum principle, $0 \leq u(t, x) \leq 1$.

(b) Let $u_1(t, x) = u(t, x)$ and $u_2(t, x) = u(t, 1 - x)$, then u_1 and u_2 satisfy the heat equation, $u_1(t, 0) = u_2(t, 0)$, $u_1(t, 1) = u_2(t, 1)$, and $u_1(0, x) = u_2(0, x)$. So by the uniqueness of initial boundary problems of heat equations, we have $u_1(t, x) = u_2(t, x)$ for all $(t, x) \in [0, \infty) \times [0, 1]$.

(c) Since

$$\frac{d}{dt} \int_0^1 u^2(t, x) dx = \int_0^1 2uu_t dx$$

$$\begin{aligned}
&= \int_0^1 2uu_{xx} \, dx \\
&= - \int_0^1 2u_x^2 \, dx \\
&\leq 0,
\end{aligned}$$

where the boundary terms of the integration by parts vanish since u vanishes on the boundary, $\int_0^1 u^2(t, x) \, dx$ is a decreasing function of t .

4. (a) For u_1 and u_2 satisfying the problem, consider $v = u_1 - u_2$, then v satisfies that

$$\begin{cases} v_t - v_{xx} = 0 & (t, x) \in (0, \infty) \times (0, \ell); \\ v_x(t, 0) = v_x(t, \ell) = 0 & t \in (0, \infty); \\ v(0, x) = 0 & x \in [0, 1]. \end{cases}$$

Consider

$$E(t) = \int_0^\ell v^2(t, x) \, dx.$$

Then

$$\begin{aligned}
\frac{dE(t)}{dt} &= \int_0^\ell 2vv_t \, dx \\
&= \int_0^\ell 2vv_{xx} \, dx \\
&= - \int_0^\ell 2v_x^2 \, dx \\
&\leq 0,
\end{aligned}$$

where the boundary terms of the integration by parts vanish since v_x vanishes on the boundary. Since $E(0) = 0$, $E(t) = 0$ and $v = 0$.

(b) Let $u = e^{-t}v$, then

$$e^{-t}v_t - e^{-t}v - e^{-t}v_{xx} + e^{-t}v = 0,$$

giving that

$$v_t - v_{xx} = 0.$$

Moreover,

$$v(0, x) = u(0, x) = \phi(x).$$

So

$$v(t, x) = \int_{\mathbb{R}} S(t, x - y)\phi(y) \, dy.$$

Therefore

$$u(t, x) = e^{-t}v(t, x) = e^{-t} \int_{\mathbb{R}} S(t, x - y)\phi(y) dy,$$

and it is easy to verify that u is a solution.

(c) Let $u = e^{-t^3/3}v$, then

$$e^{-t^3/3}v_t - t^2e^{-t^3/3}v - e^{-t^3/3}v_{xx} + t^2e^{-t^3/3}v = 0,$$

giving that

$$v_t - v_{xx} = 0.$$

Moreover,

$$v(0, x) = u(0, x) = \phi(x).$$

So

$$v(t, x) = \int_{\mathbb{R}} S(t, x - y)\phi(y) dy.$$

Therefore

$$u(t, x) = e^{-t^3/3}v(t, x) = e^{-t^3/3} \int_{\mathbb{R}} S(t, x - y)\phi(y) dy,$$

and it is easy to verify that u is a solution.

5. (a) By d'Alembert's formula,

$$\begin{aligned} u(t, x) &= \frac{e^{x+t} + e^{x-t}}{2} + \frac{1}{2} \int_{x-t}^{x+t} \sin s ds \\ &= \frac{e^{x+t} + e^{x-t}}{2} + \frac{\cos(x-t) - \cos(x+t)}{2}. \end{aligned}$$

(b) By d'Alembert's formula,

$$\begin{aligned} u(t, x) &= \frac{\log[1 + (x+t)^2] + \log[1 + (x-t)^2]}{2} + \frac{1}{2} \int_{x-t}^{x+t} (4+s) ds \\ &= \frac{\log[1 + (x+t)^2] + \log[1 + (x-t)^2]}{2} + \frac{(4+x+t)^2 - (4+x-t)^2}{4} \\ &= \frac{\log[1 + (x+t)^2] + \log[1 + (x-t)^2]}{2} + t(4+x). \end{aligned}$$

6. (a)

$$\begin{aligned} A_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} x \cos nx dx = 0; \\ B_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} x \sin nx dx = \frac{2(-1)^{n+1}}{n}. \end{aligned}$$

So

$$x = \sum_{n=1}^{\infty} \frac{2(-1)^{n+1}}{n} \sin nx.$$

(b) First we find separated solutions $u(t, x) = T(t)X(x)$. Then

$$T''(t)X(x) = T(t)X''(x),$$

giving that

$$-\frac{T''(t)}{T(t)} = -\frac{X''(x)}{X(x)} = \lambda.$$

$X(x)$ satisfies that $X(0) = X(\pi) = 0$. If $\lambda = 0$, then $X(x) = Ax + B$. From the boundary conditions, $A = B = 0$. If $\lambda \neq 0$, suppose that $-\lambda = \beta^2$, then $X(x) = Ae^{\beta x} + Be^{-\beta x}$. From the boundary conditions,

$$\begin{vmatrix} 1 & 1 \\ e^{\beta\pi} & e^{-\beta\pi} \end{vmatrix} = 0.$$

So $\beta = ni$ and $\lambda = -n^2$ for $n \in \mathbb{Z} \setminus \{0\}$. Then $X_n(x) = \sin nx$ for $n \in \mathbb{N}^+$. Then $T_n(t) = A_n \cos nt + B_n \sin nt$. By (a),

$$x = \sum_{n=1}^{\infty} \frac{2(-1)^{n+1}}{n} \sin nx.$$

Moreover,

$$0 = \sum_{n=1}^{\infty} 0 \sin nx.$$

Therefore,

$$u(t, x) = \sum_{n=1}^{\infty} \frac{2(-1)^{n+1}}{n} \cos nt \sin nx.$$

7. (a) We first prove that

$$\begin{aligned} v(x) &\leq \frac{1}{4\pi r^2} \int_{\partial B_r(x)} v(y) \, dy. \\ \frac{1}{4\pi r^2} \int_{\partial B_r(x)} v(y) \, dy &= \frac{1}{4\pi} \int_{\partial B_1(0)} v(x + ry) \, dy. \end{aligned}$$

So

$$\begin{aligned} \frac{d}{dr} \left(\frac{1}{4\pi r^2} \int_{\partial B_r(x)} v(y) \, dy \right) &= \frac{1}{4\pi} \int_{\partial B_1(0)} \nabla v(x + ry) \cdot y \, dy \\ &= \frac{1}{4\pi r^2} \int_{\partial B_r(x)} \nabla v(y) \cdot \mathbf{n} \, dy \\ &= \frac{1}{4\pi r^2} \int_{B_r(x)} \Delta v(y) \, dy \\ &\geq 0. \end{aligned}$$

Moreover,

$$\lim_{r \rightarrow 0} \frac{1}{4\pi r^2} \int_{\partial B_r(x)} v(y) \, dy = v(x).$$

Therefore,

$$v(x) \leq \frac{1}{4\pi r^2} \int_{\partial B_r(x)} v(y) \, dy.$$

So

$$4\pi r^2 v(x) \leq \int_{\partial B_r(x)} v(y) \, dy.$$

Integrating r , we have

$$\frac{4\pi r^3}{3} v(x) \leq \int_{B_r(x)} v(y) \, dy,$$

giving that

$$v(x) \leq \frac{3}{4\pi r^3} \int_{B_r(x)} v(y) \, dy.$$

(b) Suppose that v attains its maximum M at x_0 . If $x_0 \in \partial U$, the proof ends. If $x_0 \in U$, we then show that $v = M$ on \bar{U} . Let

$$E = \{x \in U \mid v(x) = M\}.$$

Then E is closed. Next we show that E is open. Suppose that $x \in E$ and $\overline{B_r(x)} \subset U$.

By (a),

$$M = v(x) \leq \frac{3}{4\pi r^3} \int_{B_r(x)} v(y) \, dy.$$

So $v = M$ on $B_r(x)$ and $B_r(x) \subset E$. Therefore E is open. Since U is connected and E is nonempty, $E = U$. Hence $v = M$ on \bar{U} and $\max_{\partial U} v = M$.

8. Consider $v = \pm u + \frac{|x|^2}{6} \lambda$ where $\lambda = \max_{\bar{U}} |f|$. Then

$$\Delta v = \mp f + \lambda \geq 0.$$

By 7,

$$\max_{\bar{U}} v = \max_{\partial U} v \leq \max_{\partial U} |g| + C(U)\lambda.$$

So

$$\pm u \leq \max_{\bar{U}} v \leq C(U) \left(\max_{\partial U} |g| + \max_{\bar{U}} |f| \right)$$

on \bar{U} , giving that

$$\max_{\bar{U}} |u| \leq C(U) \left(\max_{\partial U} |g| + \max_{\bar{U}} |f| \right).$$

9. (a)

$$\begin{aligned} \widehat{f}'(n) &= \frac{1}{2\pi} \int_{-\pi}^{\pi} f'(x) e^{-inx} \, dx \\ &= \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) \cdot i n e^{-inx} \, dx \\ &= i n \widehat{f}(n). \end{aligned}$$

$$\hat{f}(0) = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) dx = 0.$$

By Parseval's equality,

$$\begin{aligned} \frac{1}{2\pi} \int_{-\pi}^{\pi} |f(x)|^2 dx &= \sum_{n \neq 0} |\hat{f}(n)|^2 \\ &\leq \sum_{n \neq 0} |in\hat{f}(n)|^2 \\ &= \frac{1}{2\pi} \int_{-\pi}^{\pi} |f'(x)|^2 dx. \end{aligned}$$

(b) Consider the odd extension \tilde{f} of f across 0. Then \tilde{f} is C^1 on $[-\pi, \pi]$, $\tilde{f}(-\pi) = \tilde{f}(\pi)$, and $\int_{-\pi}^{\pi} \tilde{f}(x) dx = 0$. By (a),

$$\int_{-\pi}^{\pi} |\tilde{f}(x)|^2 dx \leq \int_{-\pi}^{\pi} |\tilde{f}'(x)|^2 dx,$$

giving that

$$\int_0^{\pi} |f(x)|^2 dx \leq \int_0^{\pi} |f'(x)|^2 dx.$$