

## Section 6. Asymptotic Behavior of Weak Solutions for System of Conservation Laws

$$\begin{cases} \partial_t u + \partial_x f(u) = 0, & u \in \mathbb{R}^n \\ u(x, t = 0) = u_0(x) \end{cases} \quad (6.1)$$

$$\lim_{x \rightarrow +\infty} u_0(x) = u_r, \quad \lim_{x \rightarrow -\infty} u_0(x) = u_l$$

$u_l, u_r \in \mathbb{R}^n$  are two constant states.

A1 (6.1) is strictly hyperbolic.

$$\begin{array}{l} \text{right eigenvector} \\ \text{left eigenvector} \end{array} \quad \begin{array}{l} \lambda_1(u) < \lambda_2(u) < \cdots < \lambda_n(u), \quad u \in \mathbb{R}^n \\ \gamma_1(u), \gamma_2(u) \cdots \gamma_n(u) \\ l_1(u), l_2(u) \cdots l_n(u) \end{array}$$

A2 Each characteristic field is either genuinely nonlinear or linearly degenerate.

$$\begin{cases} \partial_t u + \partial_x f(u) = 0 \\ u(x, t = 0) = \begin{cases} u_l & x < 0 \\ u_r & x > 0 \end{cases} \end{cases}$$

$u_R(x, t)$ ,  $\exists u_0 = u_l, u_1, u_2, \dots, u_n = u_r$  such that  $u_i = T_i(u_{i-1})$ , here  $T_i(u)$  is the  $i$ -th wave curve through the base point  $u$ .

Let  $U(x, t)$  be the unique viscosity solution to the Cauchy problem (6.1) - (6.2). Then one can regard  $U(x, t)$  as a limit of approximate solutions constructed by Glimm's method.

Goal: What will be large time asymptotic behavior of  $U(x, t)$ , as  $t \rightarrow +\infty$ ?

If  $n = 1$ ,  $f$  is convex, then this problem is well understood.

Case 1:  $u_l = u_r$ ,  $u_R(x, t) = u_l = u_r$ , then  $u(x, t) - u_r \rightarrow 0$  as  $t \rightarrow \infty$ , with the decay rate  $\left(\frac{1}{\sqrt{t}}\right)$ .

Case 2:  $u_l > u_r$ ,  $U(x, t) \rightarrow u_R$  with a phase shift (shock).

Case 3:  $u_l < u_r$ ,  $U(x, t) \rightarrow u_R$  (rarefaction wave).

Conclusion ( $n = 1$ ). The large time behavior of  $U(x, t)$  is determined completely by the far fields of the initial data  $(u_l, u_r)$ , in other words, the Riemann solution is “stable”.

Question: What happens for  $n > 1$ ?

Let  $U(x, t)$  be a Glimm solution. By the strict hyperbolicity,

$$\lambda_1(U(x, t)) < \lambda_2(U(x, t)) < \cdots < \lambda_n(U(x, t))$$

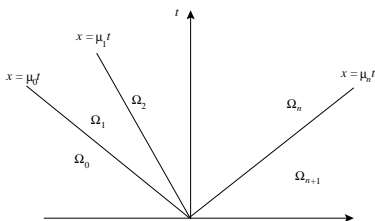
$\exists \delta > 0, \mu_i, i = 0, 1, \dots, n$  such that

$$\mu_0 + \delta \leq \min_{(x,t)} \lambda_1(U(x, t))$$

$$\max_{(x,t)} \lambda_i(U(x, t)) + \delta \leq \mu_i \leq \min_{(x,t)} \lambda_{i+1}(U(x, t)) - \delta, \quad i = 1, \dots, n-1$$

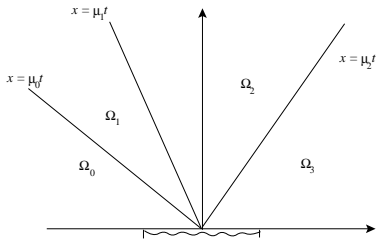
$$\max_{(x,t)} \lambda_n(U(x, t)) + \delta \leq \mu_n$$

Primary Region:  $\Omega_i$ ,  $i = 0, 1, \dots, n + 1$ , is defined as



$$\begin{aligned}\Omega_0 &= \{(x, t); x < \mu_0 t\} \\ \Omega_1 &= \{(x, t); \mu_{i-1} t < x < \mu_i t\} \quad i = 1, \dots, n \\ \Omega_{n+1} &= \{(x, t); x > \mu_n t\}\end{aligned}$$

Example:  $n = 2$ .



### Theorem 6.1 (Asymptotic behavior toward Riemann Solutions)

Let  $U(x, t)$  be the viscosity solution to (6.1) and (6.2) with small initial total variation. Let  $u_R(x, t)$  be the corresponding Riemann solution  $w, \gamma, t, (u_l, u_r)$  solved by elementary waves  $(u_{i-1}, u_i)$  as described before. Then

(1)  $U(x, t) \rightarrow u_i$  as  $t \rightarrow +\infty$  as  $\frac{x}{t} = \mu_i$ .

(2) If  $\nabla \lambda_i \cdot \gamma_i > 0, i = \alpha_1, \dots, \alpha_p$ .

$(u_{i-1}, u_i)$  is  $i$ -rarefaction wave, i.e.,  $\lambda_i(u_{i-1}) \leq \lambda_i(u_i)$ . Then the amount of  $i$ -shock wave in  $\Omega_i$  approach zero as  $t \rightarrow \infty$  and  $U(x, t)$  approaches the centered rarefaction waves  $(u_{i-1}, u_i)$ .

- (3) If  $\nabla \lambda_i \cdot \gamma_i > 0$ ,  $i = \alpha_1, \dots, \alpha_p$  and  $(u_{i-1}, u_i)$  is an  $i$ -shock, i.e.,  $\lambda_i(u_{i-1}) > \lambda_i(u_i)$ . Then in  $\Omega_i$ , the solution  $U(x, t)$  approaches  $(u_{i-1}, u_i)$  both in strength and in shock speed, furthermore, the total variation of  $U(x, t)$  in  $\Omega_i$  away from the shock approach zero.
- (4) If  $\nabla \lambda_i \cdot \gamma_i \equiv 0$ ,  $i = \beta_1, \dots, \beta_{n-p}$ . In this case,  $(u_{i-1}, u_i)$  is an  $i$ -contact discontinuity, and  $\lambda_i(u_{i-1}) = \lambda_i(u_i)$ . Then

$$\lambda_i(U(x, t)) \rightarrow \lambda_i(U_i) = \lambda_i(U_{i-1}), \quad (x, t) \in \Omega_i, \quad t \rightarrow +\infty.$$

The distance between  $\{U(x, t), (x, t) \in \Omega_i\}$  and  $T_i u_{i-1} = T_i u_i$  approach zero as  $t \rightarrow +\infty$ .



## Main Idea of the Proof:

- (1) Nonlinearity introduces dissipation: expansion waves cancels compressive waves (due to entropy condition).
- (2) Decoupling of waves  $\longleftrightarrow$  nonlinear superposition principle  
 $\iff$  when  $t \gg 1$ , only  $i$ -wave dominates on  $\Omega_i$ .

## Theorem 6.2

Let  $X_k^i (i = 1, 2; 1 \leq k \leq n)$  be two generalized  $k$ -characteristic issued from two points on  $t = t_0$ , with  $X_k^1 \leq X_k^2$ , let  $t_1 (\geq t_0)$  be any time after which  $X_k^j$  does not intersect  $X_i^{j'}$  for  $i \neq k$ . Denote by  $D_k(t)$  the distance between  $X_k^1$  and  $X_k^2$ , i.e.

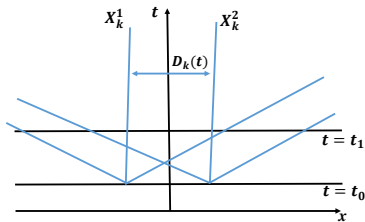
$$D_k(t) = X_k^2(t) - X_k^1(t)$$

$\mathbb{X}_k^+(t)$ : amount of  $k$ -rarefaction wave between  $X_k^1$  and  $X_k^2$ .

$\mathbb{X}_k^-(t)$ : amount of  $k$ -shock between  $X_k^1$  and  $X_k^2$  (does not include  $X_k^i$ ,  $i = 1, 2$ ). Then for  $t > t_1$ ,

$$\mathbb{X}_k^+(t) \leq \frac{D_k(t)}{t - t_1} + O(1)[Q_k(t_0, t) + h_k(t_0, t)]$$

where  $Q_k(t_0, t)$  is amount of wave interaction between  $t_0$  and  $t$  and  $X_k^1$  and  $X_k^2$ ,  $h_k(t_0, t)$  is the total amount of  $i$ -waves crossing  $X_k^1$  ( $X_k^2$ ) for all  $i > k$  ( $i < k$ ) between  $t_0$  and  $t$ .



## Proof of Theorem 6.2

### Step 1: Approximation conservation laws

Let  $\Lambda$  be a region bounded by either generalized characteristic or space like curves.

$$L_i^\pm(\Lambda) = E_i^\pm(\Lambda) \mp c_i(\Lambda) + O(1) Q(\Lambda)$$

$L_i(\Lambda)$  :  $i$ -waves leaving  $\Lambda$ .

$E_i(\Lambda)$  :  $i$ -waves entering  $\Lambda$ .

$c(\Lambda)$  :  $i$ -wave cancellation in  $\Lambda$ .

$Q(\Lambda)$  : interaction happening in  $\Lambda$ .

## Step 2: Expansion of rarefaction waves

Recall that a generalized characteristic curve is piecewise Lipschitz continuous which is either a genuine characteristic or a shock.

$\tilde{X}_k(t)$  is the total amount of  $j$ -waves ( $j \neq k$ ) between  $X_k^1(t)$  and  $X_k^2(t)$ .

Denote  $u_k^{\pm i}(t) = U(X_k^i(t) \pm, t)$ , then

$$\begin{aligned} \dot{D}_k(t) &= \frac{d}{dt} D_k(t) = \frac{d}{dt} (X_k^2(t)) - \frac{d}{dt} X_k^1(t) \\ &= \lambda_k(u_k^{+2}(t), u_k^{-2}(t)) - \lambda_k(U_k^{+1}(t), U_k^{-1}(t)) \end{aligned}$$

Fact 1:  $\exists \theta(t) \in (0, 1)$  such that

$$\dot{D}_k(t) = \theta(t)(\lambda_k^{-2}(t) - \lambda_k^{+1}(t)) + (1 - \theta(t))(\lambda_k^{+2}(t) - \lambda_k^{-1}(t))$$

where  $\lambda_k^{\pm i}(t) = \lambda_k(U(X_k^i(t) \pm, t))$  (due to entropy condition).

In fact,  $\exists \theta \in (0, 1)$  such that

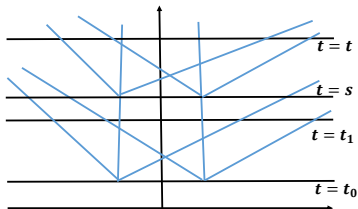
- (1)  $\dot{D}_k(t) \leq \theta(\lambda_k^{-2}(t) - \lambda_k^{+1}(t)) + (1 - \theta)(\lambda_k^{+2}(t) - \lambda_k^{-1}(t))$ .
- (2)  $\dot{D}_k(t) = (\lambda_k^{-2}(t) - \lambda_k^{+1}(t)) + (1 - \theta(t))$  (strength  $X_k^1$  + strength  $X_k^2$ ).

Indeed,

$$\dot{D}_k(t) = \theta(t) (\lambda_k^{-2}(t) - \lambda_k^{+1}(t)) + (1 - \theta(t)) (\lambda_k^{-2}(t) - \lambda_k^{-2}(t) + \lambda_k^{-2}(t) - \lambda_k^{+1}(t) + \lambda_k^{+1}(t) - \lambda_k^{-1}(t)).$$

Fact 2:

- (1)  $\exists t_1 > t_0$ , such that all  $X_k^j$  ( $k \neq k^1$ ) crosses  $X_{k'}^{j'}$  before  $t_1$ .
- (2)  $\forall t > t_1$ ,  $\exists s$  such that  $t - s = O(1)$   $D_k(t)$  such that  $X_{k-1}^2$  crosses  $X_k^1$  and  $X_{k+1}^1$  crosses  $X_k^2$  before  $t$ .



Step 2.1: Estimate  $\tilde{X}_k(t)$  (total amount of  $i$ -waves ( $i \neq k$ ) crossing  $(X_k^1(t), X_k^2(t))$ ).

If  $i < k$ , by applying approximation conservation laws.

$$\tilde{X}_k^i(t) \leq O(1) \int_s^t d(h_k(t_0, \tau) + Q_k(t_0, \tau))$$

for  $i > k$ , similar estimate also holds, so,

$$\tilde{X}_k(t) \leq O(1) \int_s^t d(h_k(t_0, \tau) + Q_k(t_0, \tau))$$

Step 2.2:

$$\begin{aligned} & \lambda_k^{-2}(t) - \lambda_k^{+1}(t) \\ = & \mathbb{X}_k^+(t) + \mathbb{X}_k^-(t) + O(1) \tilde{\mathbb{X}}_k(t) \\ = & \mathbb{X}_k^+(t) + \mathbb{X}_k^-(t) + O(1) \int_s^t d(Q_k(t_0, \tau) + h_k(t_0, \tau)) \end{aligned}$$

Then

$$\begin{aligned} \dot{D}_k(t) &= \lambda_k^{-2}(t) - \lambda_k^{-1}(t) + (1 - \theta(t)) (\text{str } X_k^2(t) + \text{str } X_k^1(t)) \\ &= \mathbb{X}_k^+(t) + \mathbb{X}_k^-(t) + O(1) \int_s^t d(Q_k(t_0, \tau) + h_k(t_0, \tau)) \\ &\quad + (1 - \theta(t)) (\text{str } X_k^2(t) + \text{str } X_k^1(t)) \end{aligned}$$



Integrate with respect to  $t$  from  $t_1$  to  $t$ ,

$$\begin{aligned}
 & D_k(t) - D_k(t_1) \\
 = & \int_{t_1}^t (\mathbb{X}_k^+(\tau) + \mathbb{X}_k^-(\tau)) d\tau + \int_{t_1}^t (1 - \theta(\tau)) (\text{str } X_k^2(\tau) + \text{str } X_k^1(\tau)) d\tau \\
 & + O(1) \int_{t_1}^t \int_s^\tau d(Q_k(t_0, \xi) + h_k(t_0, \xi)) d\tau \\
 = & \int_{t_1}^t (\mathbb{X}_k^+(\tau) + \mathbb{X}_k^-(\tau) + (1 - \theta(\tau)) (\text{str } X_k^2(\tau) + \text{str } X_k^1(\tau))) d\tau \\
 & + O(1) \int_{t_1}^t (\tau - s) d(Q_k(t_0, \tau) + h_k(t_0, \tau))
 \end{aligned}$$

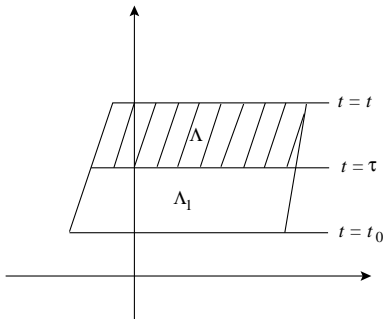
Recall the approximate conservation law

$$\begin{aligned}
 L_i^\pm(\Lambda) &= E_i^\pm(\Lambda) \mp c_i(\Lambda) + O(1) Q(\Lambda) \tag{6.3} \\
 &\left( \text{for } \alpha, \beta, c(\alpha, \beta) = \frac{1}{2}(|\alpha| + |\beta| - |\alpha + \beta|) \right)
 \end{aligned}$$

So if we apply (6.3) to  $\Lambda$ ,

$$\begin{aligned}\mathbb{X}_k^+(t) &= \mathbb{X}_k^+(\tau) - c_k(\Lambda) + O(1) Q_k(\tau, t) \\ &\leq \mathbb{X}_k^+(\tau) + O(1) Q_k(\tau, t)\end{aligned}$$

i.e.  $\mathbb{X}_k^+(\tau) \geq \mathbb{X}_k^+(t) + O(1) Q_k(\tau, t)$



Similarly, applying (6.3) to  $\Lambda_1$ ,

$$\begin{aligned}\mathbb{X}_k^-(\tau) &= \mathbb{X}_k^-(t_0) + c_k(\Lambda_1) + O(1) Q_k(t_0, \tau) \\ &\geq \mathbb{X}_k^-(t_0) + O(1) Q_k(t_0, \tau)\end{aligned}$$

Therefore,

$$\begin{aligned}D_k(t) &\geq D_k(t_1) + \mathbb{X}_k^+(t)(t - t_1) \\ &\quad + O(1) \int_{t_1}^t Q_k(\tau, t) d\tau + \mathbb{X}_k^-(t_0)(t - t_1) \\ &\quad + O(1) \int_{t_1}^t Q_k(t_0, \tau) d\tau \\ &\quad + O(1) \int_{t_1}^t (1 - \theta(\tau)) (\text{str } X_2(\tau) + \text{str } X_1(\tau)) d\tau \\ &\quad + O(1) \int_{t_1}^t D_k(\tau) d(Q_k(t_0, \tau) + h_k(t_0, \tau))\end{aligned}$$

so,

$$\begin{aligned} & \mathbb{X}_k^+(t) \\ \leq & \frac{D_k(t)}{t-t_1} + \left( -\mathbb{X}_k^-(t_0) + \frac{(-1)}{t-t_1} \int_{t_1}^t (1-\theta(\tau)) \right) (\text{str } X_2(t) + \text{str } X_1(\tau)) d\tau \\ & + O(1) Q_k(t_0, t) + O(1) \frac{1}{t-t_1} \int_{t_1}^t D_k(\tau) d(Q_k(t_0, \tau) + h_k(t_0, \tau)) \end{aligned}$$

Immediately, we obtain that

$$\mathbb{X}_k^+(t) \leq \frac{D_k(t)}{t-t_1} + O(1) [-\mathbb{X}_k^-(t_0) - \max \text{str } X_k(\tau) + Q_k(t, t_0) + h_k(t, t_0)]$$

This is true for any two characteristic, just do this procedure for the increasing variation part.

Next, we turn to the Proof of Theorem 6.1.

### Lemma 6.1

$\exists \delta_0 > 0$ , such that if  $T.V. u_0 \leq \delta_0$ , then

$$(1) \quad T.V. u(\cdot, t) \leq c_0 \delta_0 \quad \forall t > 0.$$

$$(2) \quad Q(0, t) \leq c_1 \quad \forall t > 0.$$

here  $Q(t_1, t_2)$  is the total amount of wave interaction taken place between  $t_1$  and  $t_2$ .

**Proof of Lemma 6.1:** Since  $u(x, t)$  is a solution generated by Glimm's scheme, so (1) is true for the Glimm approximate solution, thus it is true for its limit.

To see (2), we consider any  $J$ -curve  $J$  and its immediate successor,

$$\begin{aligned} Q(J') - Q(J) &\leq -D(\Delta) + O(1) D(\Delta) L(J) \\ &\leq -\frac{1}{2} D(\Delta) \quad (\text{by (1), } L(J) \text{ is sufficiently small.}) \end{aligned}$$

Now, we consider a region  $\Lambda$  whose domain of dependence contains a mesh curve  $J$ , and  $\Lambda$  consists of all diamonds. Then summing the above inequality up, we can get

$$Q(\Lambda) = \sum_{\Delta \in \Lambda} D(\Delta) \leq 2Q(J) \leq c_1.$$

Take limit to Glimm's approximate solution

$$Q(\Lambda) \leq c_1 \implies Q(0, t) \leq c_1.$$

### Lemma 6.2

$\forall \varepsilon > 0, \exists t_0 = t_0(\varepsilon)$  and  $M = M(\varepsilon)$  such that

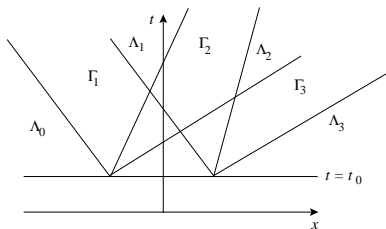
- (1)  $Q(t_0, t) < \varepsilon \quad \forall t > t_0.$
- (2)  $T.V._{\{|x| \geq M\}} u(\cdot, t_0) < \varepsilon.$

**Proof of Lemma 6.2:** These follow from Lemma 6.1.

### Lemma 6.3

Let  $\Gamma_i$  be the region between  $X_i^1$  and  $X_i^2$ ,  $\Lambda_0$  be the region left of  $X_1^1$ ,  $\Lambda_i$  be the region between  $X_1^2$  and  $X_{i+1}^1$ ,  $i = 1, \dots, n-1$ , and  $\Lambda_n$  be the right of  $X_n^2$ . Then for any  $t \geq t_1$ .

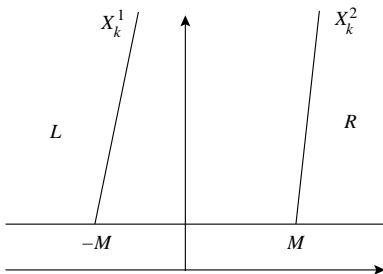
- (1)  $\overset{T.V.}{\Lambda_i} u(\cdot, t) = O(\varepsilon).$
- (2) The amount of  $j$ -waves outside of  $\Gamma_j$  at time  $t$  is  $O(\varepsilon).$
- (3)  $\overset{\text{osc}}{\Lambda_i} u(\cdot, \cdot) = O(1) \varepsilon.$
- (4)  $\mathbb{X}_j^+(t) \leq \frac{D_j(t)}{t - t_1} + O(1) \varepsilon.$





**Proof of Lemma 6.3:** We start with (2). Applying approximate conservation law to the Region  $R$  (which is the right of  $X_k^2$  for  $i \leq k$ ), one can get

$$L_i^\pm(R) \leq E_i^\pm(R) + O(1)Q(R).$$



Therefore, the total amount of  $i$ -wave ( $i < k$ ) crossing  $X_k^2 = O(1)\varepsilon$ , the total amount of  $k$ -waves on the right of  $X_k^2 = O(1)\varepsilon$ . Similarly, one can apply the approximate conservation law to  $L$  (left of  $X_k^1$ ) for  $i \geq k$ , then the total amount of  $i$ -wave ( $i > k$ ) crossing  $X_k^1 = O(1)\varepsilon$ , the total amount of  $k$ -wave on the left of  $X_k^1 = O(1)\varepsilon$ .

Thus (2) is true, and

$$h_k(t_0, t) = O(1)\varepsilon.$$

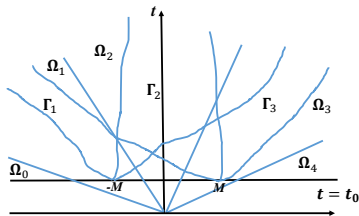
Immediately,

$$\mathbb{X}_k^+(t) \leq \frac{D_k(t)}{t - t_1} + O(1)\varepsilon,$$

(1) and (3) are consequence of (2).

## Lemma 6.4

The total amount of  $i$ -waves in the region  $\Omega_j$ ,  $j \neq i$ ,  $j = 0, 1, \dots, n, n+1$  at time  $t$  approaches 0 as  $t \rightarrow +\infty$ .



**Proof of Lemma 6.4:** By the strict hyperbolicity, for large  $t$ , one has

$$\Gamma_i \subset \Omega_i,$$

so the conclusions follows from Lemma 6.3.

### Lemma 6.5 (Emergency of contact waves)

Assume that  $\nabla \lambda_i \cdot \gamma_i \equiv 0$ ,  $\forall i = \beta_1, \dots, \beta_{n-p}$ . Then

$\forall (x_k, t_k) \in \Lambda_k$ ,  $k = i - 1, i$ ,  $i = \beta_i, \dots, \beta_{n-p}$ .

$$(1) \lambda_i(u(x_i, t_i)) = \lambda_i(u(x_{i-1}, t_{i-1})) + O(\varepsilon).$$

$$(2) U(x_i, t_i) \in T_i(u(x_{i-1}, t_{i-1})) + O(\varepsilon).$$

**Proof of Lemma 6.5:** By (3) of Lemma 6.3, without loss of generality,  $t_{i-1} = t_i$ . Since  $\lambda_i(u(\cdot, t_i))$  changes only when it crosses  $j$ -waves for  $j \neq i$  which is of order  $O(\varepsilon)$ .

### Lemma 6.6 (Emergency of shock wave)

Suppose

$$\nabla \lambda_i(u) \cdot \gamma_i(u) > 0, \quad i = \alpha_1, \dots, \alpha_p .$$

Then  $\exists k_0$  such that if

$$\lambda_i(u(x_i, t_i)) \leq \lambda_i(u(x_{i-1}, t_{i-1})) - k_0 \varepsilon, \quad (x_i, t_i) \in \Lambda_i .$$

Then for sufficiently large  $t$

- (1)  $\mathbb{X}_i^+(t) = O(\varepsilon)$ .
- (2)  $X_i^1$  and  $X_i^2$  collide to form an  $i$ -shock with strength

$$\lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1})) + O(\varepsilon).$$

**Proof of Lemma 6.6:** Recall that

$$\begin{aligned}\lambda_i^{-2}(t) - \lambda_i^{+1}(t) &= \mathbb{X}_i^+(t) + \mathbb{X}_i^-(t) + O(1) \tilde{\mathbb{X}}_i(t) \\ &= \mathbb{X}_i^+(t) + \mathbb{X}_i^-(t) + O(1)\varepsilon\end{aligned}$$

$$\lambda_i^{+2}(t) - \lambda_i^{-1}(t) = \lambda_i(U(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1})) + O(\varepsilon)$$

$$\begin{aligned}\dot{D}(t) &\leq \theta(\lambda_i^{-2}(t) - \lambda_i^{+1}(t)) + (1 - \theta)(\lambda_i^{+2} - \lambda_i^{-1}(t)) \quad (0 < \theta < 1) \\ &= \theta(\mathbb{X}_i^+(t) + \mathbb{X}_i^-(t)) + (1 - \theta)(\lambda(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1}))) + O(\varepsilon) \\ &\leq \theta \mathbb{X}_i^+(t) + (1 - \theta)(\lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1}))) + O(\varepsilon) \\ &\leq \theta \frac{D_i(t)}{t - t_1} + (1 - \theta)(\lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1}))) + O(\varepsilon)\end{aligned}$$

Set

$$H_i(t) = D_i(t) - [\lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1}))](t - t_1).$$

Then

$$\dot{H}_i(t) \leq \theta \frac{H(t)}{t - t_1} + O(\varepsilon).$$

Solving this differential inequality

$$H_i(t) \leq (t - t_1)^\theta H_i(t_1 + 1) + O(\varepsilon)(t - t_1).$$

Thus,

$$\begin{aligned} D_i(t) &\leq (\lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1}))) (t - t_1) \\ &\quad + O(1) (t - t_1)^\theta + O(\varepsilon) (t - t_1) \quad (\star) \\ &= [(\lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1}))) \\ &\quad + O(1) \varepsilon] (t - t_1) + O(1) (t - t_1)^\theta \end{aligned}$$

Choose  $k_0$  sufficiently large, then

$$D_i(t) < 0 \quad \text{for } t \gg 1,$$

so the conclusions follows.

**Lemma 6.7** (Emergency of rarefaction waves)

Assume that  $\lambda_i(U(x_i, t_i)) - \lambda_i(U(x_{i-1}, t_{i-1})) \geq -O(1)\varepsilon$  for some uniform constant  $O(1) \geq 0$ . Then

- (1)  $|\mathbb{X}_i^-(t)| + |\text{str } X_i^j(t)| = O(1)\varepsilon$ .
- (2)  $U(x_i, t_i) \in R_i^+(u(x_{i-1}, t_{i-1})) + O(1)\varepsilon$ .



**Proof of Lemma 6.7:** By Lemma 6.3,

$$\begin{aligned}
 \mathbb{X}_i^+(t) &\leq \frac{D_i(t)}{t - t_1} + O(1)\varepsilon \\
 \text{(By } (\star)\text{)} &\leq \lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1})) \\
 &\quad + O(1)(t - t_1)^{\theta-1} + O(1)\varepsilon \\
 &\leq [\lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1}))] \\
 &\quad + O(1)\varepsilon \quad \text{for large } t.
 \end{aligned}$$

On the other hand,

$$\begin{aligned}
 &\lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1})) \\
 = &\lambda_i^{+2}(t) - \lambda_i^{-1}(t) + O(1)\varepsilon \\
 = &\mathbb{X}_i^+(t) + \mathbb{X}_i^-(t) + O(1)\tilde{\mathbb{X}}_i(t) + \text{str } X_i + O(1)\varepsilon
 \end{aligned}$$

so,

$$\begin{aligned} & \mathbb{X}_i^+(t) - [\lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1}))] \\ = & |\mathbb{X}_i^-(t)| + |\text{str } X_i| - O(1)\varepsilon \\ \geq & -O(1)\varepsilon \end{aligned}$$

Thus,

$$\mathbb{X}_i^+(t) = \lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1})) + O(1)\varepsilon,$$

so,

$$|\mathbb{X}_i^-(t)| + |\text{str } X_i(t)| = O(1)\varepsilon.$$

We need to relate  $(u_{i-1}, u_i)$  in  $u_R(x, t)$  to  $(u(x_{i-1}, t_{i-1}), u(x_i, t_i))$ .

**Lemma 6.8** (Comparison with the Riemann solution)

Let  $(u_{i-1}, u_i)$  be the  $i$ -th wave in the Riemann solution  
 $u_R(x, t) = u\left(\frac{x}{t}\right),$

$$\begin{cases} \partial_t u + \partial_x f(u) = 0 \\ u(x, 0) = \begin{cases} u_l & x < 0 \\ u_R & x > 0 \end{cases} \end{cases}$$

Then

$$|u(x_i, t_i) - u_i| = O(1)\varepsilon, \quad \forall (x_i, t_i) \in \Lambda_i .$$

**Proof of Lemma 6.8** It follows from Lemmas 6.5, 6.6, 6.7. We can find  $\tilde{u}_i$  such that

(1)  $|\tilde{u}_i - u(x_i, t_i)| = O(1)\varepsilon.$

(2)  $\tilde{u}_i \in T_i(\tilde{u}_{i-1}).$

(3)  $(\tilde{u}_{i-1}, \tilde{u}_i)$  is an  $i$ -th elementary wave.

i.e., the superposition of  $(\tilde{u}_{i-1}, \tilde{u}_i)$ ,  $i = 1, \dots, n$  solves

$$\begin{cases} \partial_t u + \partial_x f(u) = 0 \\ u(x, t = 0) = \begin{cases} \tilde{u}_0 & x < 0 \\ \tilde{u}_n & x > 0 \end{cases} \end{cases}$$

On the other hand, by definition,  $(u_{i-1}, u_i)$  is the  $i$ -th elementary of  $u_R(x, t) = U\left(\frac{x}{t}\right)$ ,

$$\begin{cases} \partial_t u + \partial_x f(u) = 0 \\ u(x, t = 0) = \begin{cases} u_- = u_0 & x < 0 \\ u_+ = u_n & x > 0 \end{cases} \end{cases}$$

so,

$$|\tilde{u}_0 - u_-| = |\tilde{u}_0 - u_0| = O(1)\varepsilon.$$

Similarly,

$$|\tilde{u}_n - u_+| = O(1)\varepsilon.$$

By continuous dependence of Riemann solution,

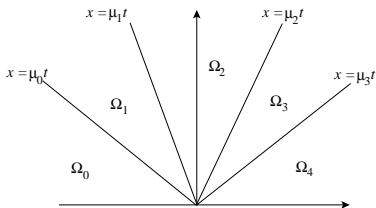
$$\begin{aligned} |\tilde{u}_i - u_i| &= O(1)\varepsilon, \quad i = 1, \dots, n, \\ \text{so } |u(x_i, t_i) - u_i| &= O(1)\varepsilon, \quad i = 1, \dots, n. \end{aligned}$$

## Final proof of Theorem 6.1

Proof of (1):  $\forall \varepsilon > 0$ , let  $X_k^j$ ,  $\Gamma_k$ ,  $\Lambda_k$  defined as before, clearly for large enough  $t$ ,  $\Gamma_i \subset \Omega_i$ . Furthermore,

$$x = \mu_i t \subset \Lambda_i,$$

so for  $x = \mu_i t$ ,  $u(x, t) = u_i + O(1)\varepsilon$ . Since  $\varepsilon$  is arbitrary, so  $u(x, t) \rightarrow u_i$  on  $\frac{x}{t} = \mu_i$ , as  $t \rightarrow \infty$ .



Proof of (2):

Step 1:  $i = \alpha_1, \dots, \alpha_p$ ,  $\nabla \lambda_i \cdot \gamma_i > 0$ ,  $(u_{i-1}, u_i)$  is a centered rarefaction wave, i.e.  $\lambda_i(u_i) \geq \lambda_i(u_{i-1})$ . Then by Lemma 6.8,

$$\lambda_i(u(x_i, t_i)) - \lambda_i(u(x_{i-1}, t_{i-1})) \geq -O(1) \varepsilon$$

$$(x_i, t_i) \in \Lambda_i, (x_{i-1}, t_{i-1}) \in \Lambda_{i-1} .$$

Then Lemma 6.7 implies

$$|\mathbb{X}_i^-(t)| + |\text{str } X_i(t) = O(1) \varepsilon.$$

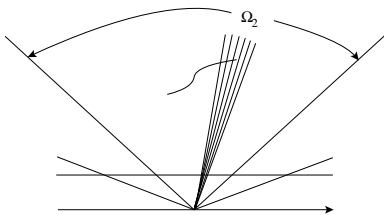
On the other hand, for large  $t$ ,  $\Gamma_i \subset \Omega_i$ , so by Lemma 6.3 that the total amount of  $i$ -shock wave in  $\Omega_i$  is order of  $O(\varepsilon)$ , which tends to zero as  $t \rightarrow +\infty$  since  $\varepsilon$  is arbitrary  $\implies$  only  $i$ -rarefaction waves left.

Step 2: We need to show in fact  $u(x, t) \rightarrow (u_{i-1}, u_i)$  in  $\Omega_i$  as  $t \rightarrow \infty$ . By Step 1,  $\exists t_2 > t_1$  such that

$|\mathbb{X}_i^-(t)| + |\text{str } X_i(t)| \leq O(1)\varepsilon$  and also the speed  $X_i^1$  and  $X_i^2$  are given by  $\lambda_i(u_{i-1}) + O(\varepsilon)$  and  $\lambda_i(u_i) + O(\varepsilon)$  respectively.

Let  $l_i^j$  ( $j = 1, 2$ ) be the edges of the centered rarefaction wave  $(u_{i-1}, u_i)$ ,  $l_i^1 = \{(x, t) \mid \frac{x}{t} = \lambda_i(u_{i-1})\}$ ,  $l_i^2 = \{(x, t) \mid \frac{x}{t} = \lambda_i(u_i)\}$ . Then for  $t \geq t_2 + O(1)D_i(t_2)$ , one has

$$|X_i^1(t) - l_i^1(t)| + |X_i^2(t) - l_i^2(t)| = O(1)\varepsilon(t - t_2) + O(1).$$





Let  $u^*(x, t)$  be the centered rarefaction wave.

Claim: (1)

$$|u^*(x, t) - u(x, t)| = O(1)\varepsilon, \quad \forall (x, t) \in (\Lambda_i \cup \Lambda_{i-1}) \cap \Omega_i.$$

$$(2) \quad |u^*(x, t) - u(x, t)| = O(1)\varepsilon, \quad \forall (x, t) \in \Gamma_i \subset \Omega_i.$$

(1)  $(x, t) \in \Lambda_{i-1} \cap \Omega_i$ , then

$$\begin{aligned} |u^*(x, t) - u(x, t)| &\leq |u^*(x, t) - u_{i-1}| + |u_{i-1} - u(x, t)| \\ &= |u^*(x, t) - u^*(l_i^1(t), t)| + O(1)\varepsilon \\ &\leq O(1) \frac{|l_i^1(t) - X_i^1(t)|}{t} + O(1)\varepsilon \\ &\leq O(1)\varepsilon \quad \text{for } t \text{ large enough.} \end{aligned}$$

Now we fix  $(x, t) \in \Gamma_i$ ,  $t \geq t_2 + O(1)D_i(t_2)$ . By Step 1, no  $i$ -th shocks and other  $j$ -waves ( $j \neq i$ ) (mod  $O(\varepsilon)$ ), and since (1) holds true.  $\exists x^* \in (l_i^1(t), l_i^2(t))$  such that

$$|u^*(x^*, t) - u(x, t)| = O(\varepsilon).$$

Through  $(x^*, t)$  we draw a generalized backward characteristics curve  $X$ , its speed changes only when it crosses other waves, since  $X$  stays in  $\Gamma_i$ , thus the total amount of other family waves are of the order  $O(\varepsilon)$ , and when it crosses the  $i$ -shocks, then its strength is  $O(\varepsilon)$ , so the speed of  $X$  is  $\lambda_i(u(x, t)) + O(\varepsilon)$

$$|x^* - x| = O(1) \varepsilon |t - t_2| + O(1)$$

then

$$\begin{aligned} |u^*(x, t) - u(x, t)| &\leq |u^*(x, t) - u^*(x^*, t)| + |u^*(x^*, t) - u(x, t)| \\ &\leq O(1) |\lambda_i(u^*(x, t)) - \lambda_i(u^*(x^*, t))| + O(\varepsilon) \\ &= O(1) \left| \frac{x - x^*}{t} \right| + O(\varepsilon) \\ &= O(\varepsilon) \quad \text{for } t \text{ sufficiently large.} \end{aligned}$$