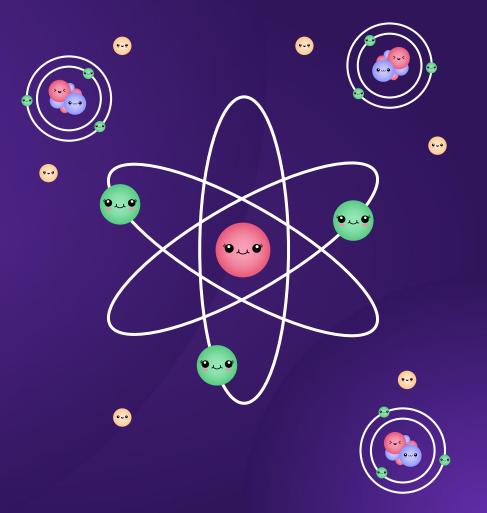
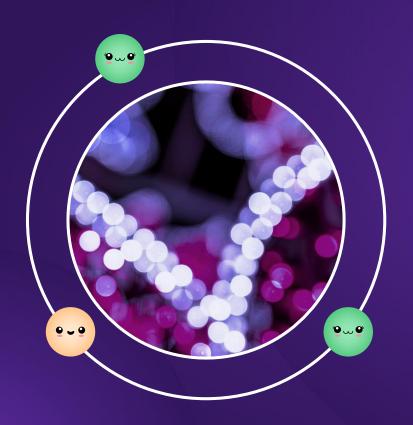
# Fusion

Chow Ka Chun Cheng Wai Yuen





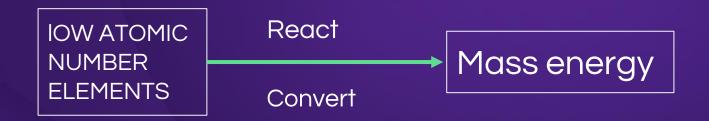
# Introduction of Fusion

- 1.1 Basics of Fusion
- 1.2 Feasibility of Fusion



#### 1.1 Basics of Fusion

- An inexhaustible source of energy for the future
- Mechanism: Under proper conditions,



E.g.

$$D+T \rightarrow {}^{4}He + n \quad (Q = 17.6 \text{ MeV})$$

#### 1.1 Basics of Fusion

Prerequisite:





> Repulsive Coulomb force



Short-range attractive nuclear force →

High temperature

The fusion reaction rate per unit volume can be written

$$R(fusions/m^3) = n_1 n_2 \langle \sigma \nu \rangle_{12}$$

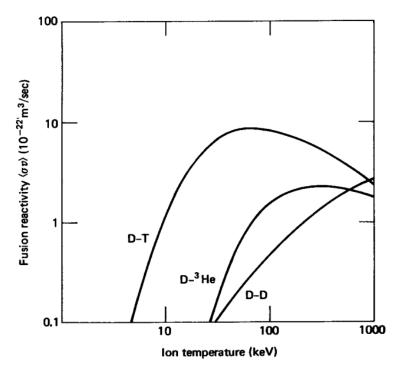


Figure 1.1 Fusion reactivity.

 D-T fusion reactivity is much greater than other fusion reactants



necessary conditions is the principal goal

# 1.2 Feasibility of Fusion

Scientific Feasibility **Engineering Feasibility** Feasibility **Practical Feasibility** Economic Feasibility & Fuel Resources

### 1.2 Scientific Feasibility

∵thermonuclear temperatures → sufficient positive energy.

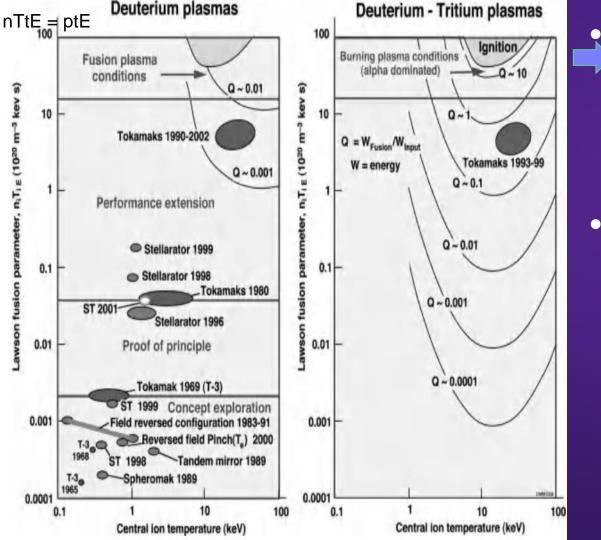
$$Q_p = \frac{\text{fusion power}}{\text{external heating power}}$$

- Using plasma power amplification factor as the measurement
  - →Qp > 1 = breakeven
  - →Qp =∞: ignition
- ∴practical definition= Qp large enough(probably >10) that net electrical power can be economically produced

The plasma power balance can be written as

fusion heating + external heating  $\ge$  radiation loss + transport loss

$$rac{1}{4}n^2\langle\sigma v
angle_{fis}U_{lpha}\left(1+rac{5}{Q_p}
ight)\geq f_z n^2 L_z+rac{3nT}{ au_E}$$



- The triple product nTtE = ptE

  High plasma pressure & long
  plasma energy confinement
  time power balance
  at high Qp (all required
  achieving)
- Magnitude of the confining  $\beta \equiv \frac{plasma\ pressure}{magnetic\ pressure} = \frac{nkT}{B^2/2\mu_0}$



Sufficiently large values of β (practical scientific feasibility)

### 1.2 Engineering Feasibility

Plasma support technologies (heat & confine the plasma)

Required technologies

- Systems included: vacuum, magnet, plasma heating, fueling etc.
- Survive in the high particle & heat fluxes
   →ITER & future fusion reactors
- R&D programme to carry out ITER

Nuclear technologies (power reactor)

- To breed tritium required for D-T plasma operation
- Heat removal & power conversion

#### 1.2 Practical Feasibility

- Sustainable modes of operation is the key
  - →E.g. For continuous/ quasi-continuous operation

Preferable

Short-pulse operation with long down-times between pulses

Compared to

- Non-inductive current-drive techniques
- →conventional tokamak confinement concept (plused

Stellarator

Practical feasibility issue

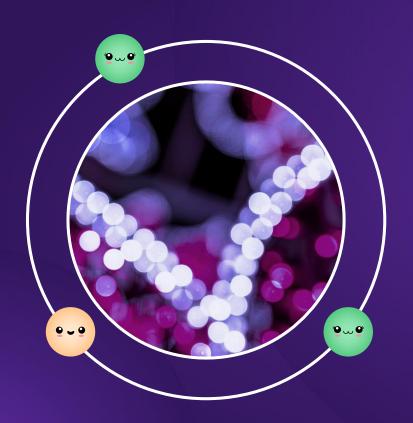
#### 1.2 Economic Feasibility & Fuel Resources

#### Today's situation

- Cost estimate of fusion electricity is expensive
   →neglected: cost of preventing carbon emission, energy potential of uranium
  - Mainly Fossil fuels

#### For future consideration

- World's estimated electricity usage in 2050 → fossil fuels gone
- Lithium (fusion)
   would last more
   than 6000 years



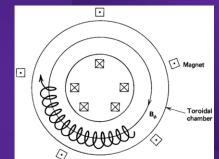
# **Confinement Concepts**

- 2.1 Magnetic Confinement
- 2.2 Tokamaks
- 2.3 Stellarator
- 2.4 Spherical Torus
- 2.5 Reversed-Field Pinch



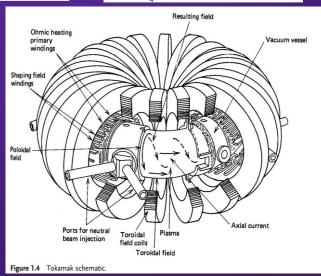
#### 2.1 Magnetic Confinement

Closed Toroidal Confinement Systems



Magnetic Confinement

- Within a confinement chamber by the proper choice of position zand currents in a set of magnetic coils
- Poloidal field produces by a toroidal current flowing in the plasma or by external coils



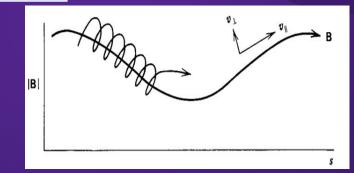
Open Toroidal Confinement Systems

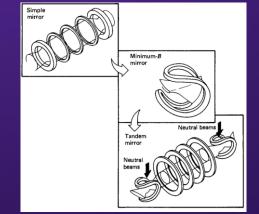
#### 2.1 Magnetic Confinement

#### Open Toroidal Confinement Systems

- Specified of confine the confinement of confinement chamber, by trapping the charged particles in a magnetic well
- When particles > values of µ
   →RHS vanishes → reflects &
   travels back along the field
   line→↑ magnetic field strength
   → >RHS vanishes again (simple mirror)

Closed Toroidal Confinement Systems





Magnetic Confinement

#### 2.2 Tokamaks

- closed toroidal confinement concepts
- The toroidal field can be represented as

$$B_{\phi}(r,\theta) = \frac{B_{\phi}^{0}}{(1+r/R_{0})\cos\theta} \equiv \frac{B_{\phi}^{0}}{1+\epsilon\cos\theta}$$

#### 2.3 Stellarator

- toroidal confinement devices
- pairs of helical conductors
- net toroidal field nulled (require toroidal field coils)
- potential advantages
  - 1. no inherent limitation to the length of operation
  - 2. Disruptive termination of the discharge



# 2.4 Spherical Torus

 minimum size central region tokamak

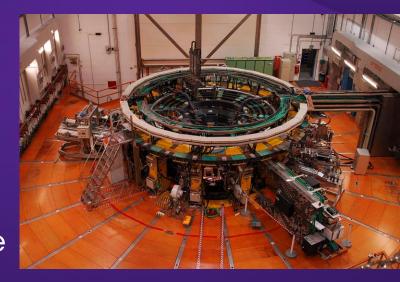
 lower magnetic field strength to confined plasma pressure

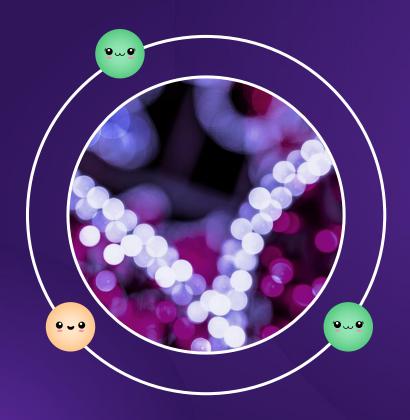


#### 2.5 Reversed-Field Pinch

 toroidal field produced by external coils

- direction of the toroidal field is reverse
- ohmic heating may be possible





# **Fusion Reactors**



# 3 Fusion Reactors (Future application)

- Plasma physics and fusion technology limits
- 2. Physical characterics of future fusion reactors and neutron sources
- 3. Discussion on the considerations of tokamak

Confinement

$$\tau_E = H_H \tau_E^{IPB98(\gamma,2)}$$

$$\tau_E = H_H \tau_E^{IPB98(\gamma,2)} \quad \tau_E^{IPB98(\gamma,2)} = 0.0562 I^{0.93} B^{0.15} P^{-0.69} \bar{n}_{e20}^{0.41} \times M^{0.19} R^{1.97} A^{-0.58} \kappa^{0.78},$$

→constraint:

- →magnetic field and plasma size → threshold power
- **Density Limit** 
  - →require achieving high particle density →high power density

$$\bar{n}_{e20} \leq \frac{I_p(MA)}{\pi a^2}$$

- →plasma current
- →plasma size
- →fusion power

- Beta Limit
  - →The MHD stability limit on the plasma pressure

$$eta_c \equiv rac{\langle n_e T_e + n_i T_i + p_lpha 
angle}{rac{B^2}{2\mu_0}} \leq eta_N rac{I(MA)}{aB} = const imes l_i rac{I}{aB}$$

 $\rightarrow$ To achieve high values of  $\beta$ N, /and  $B \rightarrow$ high plasma power density plasma size for a given fusion power level

- Kink Stability Limit
  - $\rightarrow$ related to the requirement of where  $\delta$  is the plasma triangularity
  - $\rightarrow$ limitation on the allowable combination of B,I,R and a

$$q_{95} = rac{5a^2B}{RI} \left(rac{1+\kappa^2(1+2\delta^2-1.2\delta^3)}{2}
ight) rac{\left(1.17-rac{0.65}{A}
ight)}{\left(1-rac{1}{A^2}
ight)^2} \geq 3.$$

- Startup Inductive Volt-seconds
  - →requirement consists of a component
  - 1. induce the plasma current in the absence of resistive losses
  - 2. Overcome resistive losses during startup

$$(\Delta\Phi)_{ind}=\mathit{IL}_{p}$$
  $(\Delta\Phi)_{res}=\mathit{C}_{Ejima}\mu_{0}\mathit{R}\mathit{I}$ 

- →plasma discharge maintenance after startup determines by a trade-off
- 1. Length of the burn pulse
- 2. Achievable bootstrap current
- 3. Amount of power required for non-inductive current drive
- 4. Overall design of the coil systems

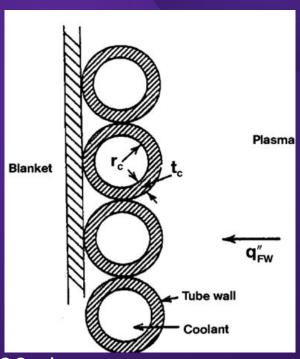
# 3 Plasma Facing Component Heat Fluxes

Stress Limits pc= coolant pressure rc= coolant tube radius tc= coolant tube thickness pdis 'B2 q=2m0= disruption pressure

$$\sigma_{th} = \frac{\alpha E \left( f_{pfc} q'' t_c + \frac{1}{2} q''' t_c^2 \right)}{2\kappa (1-\nu)} \qquad \sigma_{th} + \sigma_p \le 3S_m$$

$$\sigma_{th} + \sigma_p \leq 3S_m$$

A stress limit on the maximum allowable Heat flux to the plasma facing component q00 stress



### 3 Plasma Facing Component Heat Fluxes

Temperature Limit
 →surface temperature constrained to below Tmax

$$T_s = T_c + f_{pfc} q''_T \left(\frac{t_c}{\kappa} + \frac{1}{h}\right) + \frac{q''' t_c^2}{2\kappa} \le T_s^{\max}$$

Tc= coolant
temperature
h=surfac heat transfer
'film' coefficient
qT=Maximum
allowable heat flux

### 3 Plasma Facing Component Heat Fluxes

- Fatigue Limit
  - $\rightarrow$ subjected to cyclic loading  $\rightarrow$  fail after a mean number of cycles
  - →NS(e) which depends on the material & cyclic strain (e)

$$q_{\mathit{fat}}'' = rac{2\kappa(1-
u)}{f_{\mathit{pfc}}lpha}rac{arepsilon_{\max}(N_D)}{t_c} - rac{1}{2}q^{'''}t_c$$

- Heat Flux Limit
  - ightarrowlimit on three maximum allowable heat flu $q''_{
    m max} = \max \left\{ q''_{th}, q''_{T}, q''_{fat} 
    ight\}$
  - →high power density tokamaks may be thwarted by heat flux

limits

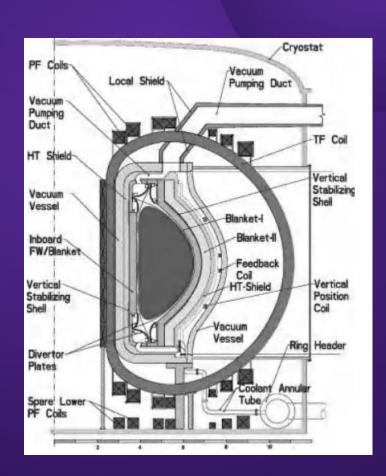
#### 3 Future Tokamak Reactors

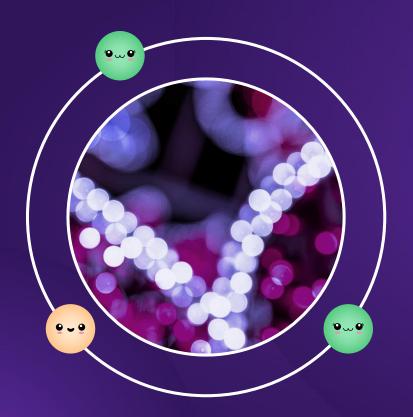
- ARIES-I
   →designed on relatively modest
   extension of ITER database
- High Field ARIES-I
   →Improvements with extended
   the toroidal field strength at the
   coil from 16 to 19T
- Reverse Shear
   →take advantage of anticipated
   physics avances → operation at
   higher βN with improved plasma
   confinement

	First stability		Reverse-shear	
	ARIES-I	High-field: ARIES-I	ARIES-RS	ARIES-AT
Major radius (m)	8.0	6.75	5.5	5.2
Plasma aspect ratio	4.5	4.5	4.0	4.0
Plasma elongation, κ	1.8	1.8	1.9	2.1
β (%) (β <sub>n</sub> )	2 (2.9)	2 (3.0)	5 (4.8)	9.2 (5.4)
Peak field at the coil (T)	16	19	16	11.5
On-axis field (T)	9.0	11	8	5.8
Neutron wail load (MW/m²)	1.5	2.5	4	3.3
ITER89-P multiplier	1.7	1.9	2.3	2.0
Plasma current (MA)	12.6	10	11.3	13
Bootstrap current fraction	0.57	0.57	0.88	0.91
Current-driver power (MW)	237	202	80	35
Recirculating power fraction	0.29	0.28	0.17	0.14
Thermal efficiency	0.46	0.49	0.46	0.59
Cost of electricity (¢/kWh)	10	8.2	7.5	5

#### 3 Future Tokamak Reactors

- ◆ ARIES-AT
   →design for an advance tokamak
   power reactor
- The volume inside of the fusion reactor pressure vessel
   1.considerably larger than Pressurized water Reactor (PWR)
  - 2. Comparable to the volume within the pressure vessel enclosing the reactor system for the gas-cooled 'next generation nuclear plant' (NGNP)





# Inertial confinement fusion

- 4.1 Introduction
- 4.2 Direct and indirect drive
- 4.3 Laser interact with matter



#### 4.1 Introduction

What laser light do

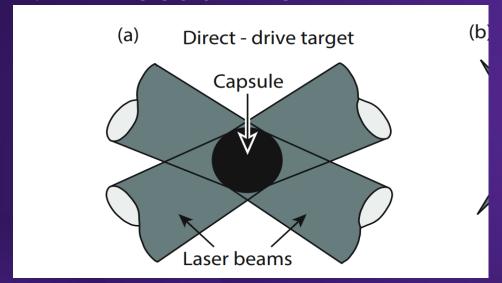
- 1.compress the D-T gas
- 2. heat up the mixture

# 4.1 Compression

$$Y = Y_0 f_B m_f^{compressed}$$

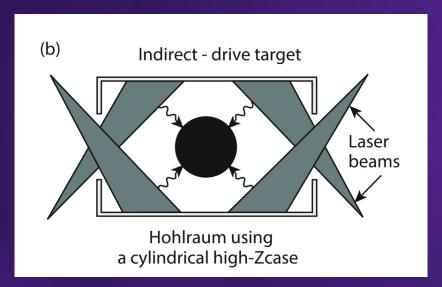
- Have practical need
  - Less energy is needed

#### 4.2 Direct drive

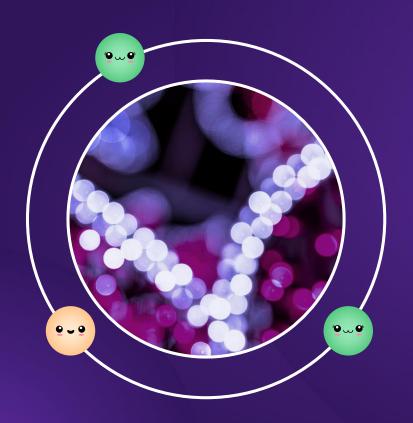


ablation of the material on the outside of the target capsule causes compression of the capsule

#### 4.2 Indrect drive



laser energy impinges on the inside of the 'hohlraum' Material of 'hohlraum' typically is high-Z material such as gold or uranium and X-rays are emitted.



# interaction of laser and matter

5.1 inverse bremsstrahlung



# 4.3.1 inverse bremsstrahlung

process by which light is absorbed in collisional plasma

$$-\ln(\alpha_T) = 2\int_0^L \kappa_{ib} \, dx = 2\frac{\nu_{ei}(n_{crit})}{c} L \int_0^1 \frac{u^2 du}{\sqrt{1-u}} = \frac{32}{15} \frac{\nu_{ei}(n_{crit})}{c} L.$$
(10.7)

Thus the total reflected energy fraction is  $\alpha_{abs} = 1 - \alpha_T$  and thus the absorbed fraction is

$$\alpha_{abs} = 1 - \exp\left(-\frac{32}{15} \frac{\nu_{ei}(n_{crit})}{c}L\right). \tag{10.8}$$

#### 5.2 resonance absorption

process by which light is absorbed at the critical density

$$\Phi(\tau) = \frac{4\left(\tau K_{\frac{1}{3}}\left(\frac{2\tau^3}{3}\right)\right)^{3/2}}{\sqrt{3\pi}\sqrt{K_{\frac{2}{3}}\left(\frac{2\tau^3}{3}\right)}}.$$
(10.12)

Here K is the modified Bessel function (these functions of order 1/3 and 2/3 are also known as Airy functions). The absorbed energy fraction is

$$f_{abs} = \Phi(\tau)^2 / 2,$$
 (10.13)

with

$$\tau \equiv (L\omega_L/c)^{1/3}\sin\theta. \tag{10.14}$$

#### **5.3 Nonlinear Effects**

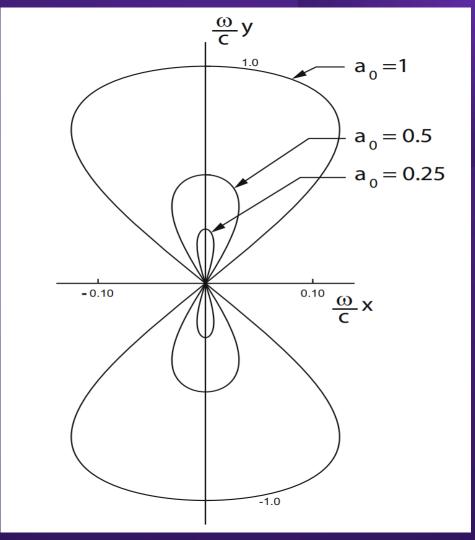
$$\mathbf{E} = E_0 \hat{\mathbf{y}} e^{i(kx - \omega t)}, \qquad \mathbf{B} = \frac{1}{c} \hat{\mathbf{x}} \times \mathbf{E}.$$

$$y(t) = -\frac{q_e E_0}{m\omega^2} \cos \omega t = -a_0 \frac{c}{\omega} \cos \omega t,$$

$$x(t) = -\frac{q_e^2 E_0^2}{8m_e^2 c} \sin 2\omega t = -\frac{a_0^2}{8} \frac{c}{\omega} \sin 2\omega t.$$

#### **5.3 Nonlinear Effects**

The motion of Electrons in High-Amplitude Electromagnetic Field would be nonlinear



# End