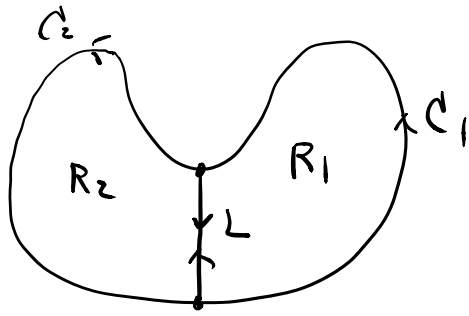


Proof of Green's Theorem for

$R =$ finite union of simple regions with intersections only along some boundary line segments, and those line segments touch only at the end points at most.

eg.



$R_1, R_2 = \text{simple}$
 but $R = R_1 \cup R_2 \neq \text{simple}$
 $\partial R_1 = C_1 + L$ ($L: \downarrow$)
 $\partial R_2 = C_2 - L$
 with anti-clockwise orientation &
 $\partial R = C_1 + C_2.$

By assumption $R = \cup R_i$ finite union s.t.

R_i are simple and

$R_i \cap R_j =$ line segment of a common boundary portion,

denote by L_{ij} ($i \neq j$) (may be empty)

$$\text{Then } \iint_R \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA = \sum_i \iint_{R_i} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA$$

$$= \sum_i \oint_{\partial R_i} M dx + N dy \quad \left(\begin{array}{l} \text{by the Green's Theorem} \\ \text{of simple region} \end{array} \right)$$

Denote $C_i =$ the part of ∂R_i with no intersection with any other R_j (except at the end points)

Then $\partial R_i = C_i + \sum_{j \neq i} L_{ij}$,

where L_{ij} is oriented according to the anti-clockwise orientation of ∂R_i

$$\begin{aligned} \text{Hence } \iint_R \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA &= \sum_i \oint_{C_i + \sum_{j \neq i} L_{ij}} M dx + N dy \\ &= \sum_i \int_{C_i} M dx + N dy + \sum_i \int_{\sum_{j \neq i} L_{ij}} M dx + N dy \end{aligned}$$

Note that, as C_i is not a common boundary of any other R_j ,

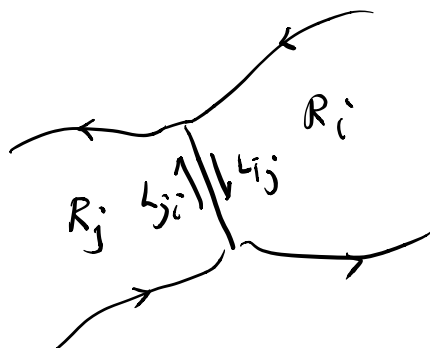
$$\sum_i C_i = \partial R.$$

$$\therefore \sum_i \int_{C_i} M dx + N dy = \oint_{\partial R} M dx + N dy$$

Finally, we have

$$L_{ji} = -L_{ij}$$

as R_i & R_j are located on the two different sides of the common boundary.



$$\begin{aligned}
\sum_i \int_{\sum_j (j \neq i) L_{ij}} M dx + N dy &= \sum_i \sum_{\substack{j \\ (j \neq i)}} \int_{L_{ij}} M dx + N dy \\
&= \sum_{\substack{i < j \\ i \neq j}} \int_{L_{ij}} M dx + N dy \\
&= \sum_{i < j} \int_{L_{ij}} M dx + N dy + \sum_{\substack{i > j \\ j < i}} \int_{L_{ji}} M dx + N dy \quad \left(\begin{array}{l} \text{interchanging} \\ \text{the notation for} \\ \text{the dummy indexes} \end{array} \right) \\
&= \sum_{i < j} \int_{L_{ij}} M dx + N dy + \sum_{i < j} \int_{L_{ji}} M dx + N dy \\
&= \sum_{i < j} \left(\int_{L_{ij}} M dx + N dy + \int_{L_{ji}} M dx + N dy \right) \\
&= \sum_{i < j} \left(\int_{L_{ij}} M dx + N dy + \int_{-L_{ij}} M dx + N dy \right) \\
&= 0 \quad \left(\text{since } L_{ji} = -L_{ij} \right)
\end{aligned}$$

This 2nd case basically include almost all situations in the level of Advanced Calculus.

The proof of general case needs "analysis" and will be omitted here. ✘

Def 12: The divergence of $\vec{F} = M\hat{i} + N\hat{j}$ is defined to be

$$\text{div } \vec{F} = \frac{\partial M}{\partial x} + \frac{\partial N}{\partial y}$$

Note: $\text{div } \vec{F} = \lim_{\epsilon \rightarrow 0} \frac{1}{\text{Area}(\bar{D}_\epsilon(x,y))} \iint_{\bar{D}_\epsilon(x,y)} \left(\frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} \right) dA$

$$= \lim_{\epsilon \rightarrow 0} \frac{1}{\text{Area}(\bar{D}_\epsilon(x,y))} \oint_{\partial \bar{D}_\epsilon(x,y)} \vec{F} \cdot \hat{n} \, ds$$

(called) "flux density"

Notation: For $f(x,y)$, $\vec{\nabla} f = \frac{\partial f}{\partial x} \hat{i} + \frac{\partial f}{\partial y} \hat{j}$ (gradient)

$$= \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} \right) f$$

It is convenient to denote

$$\left| \vec{\nabla} = \hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} \right|$$

(∇ "nabla")

Then

$$\vec{\nabla} \cdot \vec{F} = \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} \right) \cdot (M\hat{i} + N\hat{j})$$

$$= \frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} = \text{div } \vec{F}$$

Hence we also write

$$\left| \text{div } \vec{F} = \vec{\nabla} \cdot \vec{F} \right|$$

Def 13: Define $\text{rot } \vec{F}$ to be

$$\text{rot } \vec{F} = \frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \quad \text{fa } \vec{F} = M\hat{i} + N\hat{j}$$

Note: $\text{rot } \vec{F} = \lim_{\epsilon \rightarrow 0} \frac{1}{\text{Area}(\overline{D}_\epsilon(x,y))} \iint_{\overline{D}_\epsilon(x,y)} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA$

$$= \lim_{\epsilon \rightarrow 0} \frac{1}{\text{Area}(\overline{D}_\epsilon(x,y))} \oint_{\partial \overline{D}_\epsilon(x,y)} \vec{F} \cdot \hat{T} ds$$

(called)
= circulation density

Using $\vec{\nabla} = \hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y}$, we can write

$$\boxed{\text{rot } \vec{F} = (\vec{\nabla} \times \vec{F}) \cdot \hat{k}}$$

Since $\vec{F} = M\hat{i} + N\hat{j} + 0\hat{k}$ (in \mathbb{R}^3)

$$\vec{\nabla} = \hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \quad (\text{in } \mathbb{R}^3)$$

$$\Rightarrow \vec{\nabla} \times \vec{F} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ M & N & 0 \end{vmatrix} = \begin{vmatrix} \frac{\partial}{\partial x} & \frac{\partial}{\partial y} \\ M & N \end{vmatrix} \hat{k}$$

$$= \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) \hat{k}$$

$$\Rightarrow \text{rot } \vec{F} = (\vec{\nabla} \times \vec{F}) \cdot \hat{k} \quad \text{i.e. } \hat{k}\text{-component of "curl } \vec{F}\text{"}$$

where $\boxed{\text{curl } \vec{F} = \vec{\nabla} \times \vec{F}}$

Using the notations of $\left. \begin{array}{l} \text{div } \vec{F} = \vec{\nabla} \cdot \vec{F} \\ \text{curl } \vec{F} = \vec{\nabla} \times \vec{F} \end{array} \right\}$ and

the Green's Thm can be written as

Vector forms of Green's Thm

normal form

$$\oint_C \vec{F} \cdot \hat{n} ds = \iint_D \text{div } \vec{F} dA$$

or

$$\oint_C \vec{F} \cdot \hat{n} ds = \iint_D \vec{\nabla} \cdot \vec{F} dA$$

tangential form

$$\oint_C \vec{F} \cdot \hat{T} ds = \iint_D \text{curl } \vec{F} \cdot \hat{k} dA$$

or

$$\oint_C \vec{F} \cdot \hat{T} ds = \iint_D (\vec{\nabla} \times \vec{F}) \cdot \hat{k} dA$$

And Theorem 10 can be written as

Thm 10' $\Omega = \text{simply-connected}$, $\vec{F} \in C^1$. Then

$$\vec{F} = \text{conservative} \iff \text{curl } \vec{F} = \vec{\nabla} \times \vec{F} = 0$$

(check: $n=3$ case)

Note: (i) $\text{curl } \vec{F} = \vec{\nabla} \times \vec{F}$ defined only in \mathbb{R}^3 ($\supset \mathbb{R}^2$)

(ii) but $\text{div } \vec{F} = \vec{\nabla} \cdot \vec{F}$ can be defined on \mathbb{R}^n for any n .

In particular, in \mathbb{R}^3

Def 12' The divergence of $\vec{F} = M\hat{i} + N\hat{j} + L\hat{k}$ is defined to be

$$\begin{aligned}\operatorname{div} \vec{F} &= \vec{\nabla} \cdot \vec{F} = \left(\hat{i} \frac{\partial}{\partial x} + \hat{j} \frac{\partial}{\partial y} + \hat{k} \frac{\partial}{\partial z} \right) \cdot (M\hat{i} + N\hat{j} + L\hat{k}) \\ &= \frac{\partial M}{\partial x} + \frac{\partial N}{\partial y} + \frac{\partial L}{\partial z}\end{aligned}$$

Then one can easily check the following facts: (Ex!)

For C^2 function f and C^2 vector field \vec{F} :

- (i) $\vec{\nabla} \times (\vec{\nabla} f) = 0$ (i.e. $\operatorname{curl} \vec{\nabla} f = 0$)
- (ii) \vec{F} conservative $\Rightarrow \operatorname{curl} \vec{F} = \vec{\nabla} \times \vec{F} = 0$
- (iii) $\vec{\nabla} \cdot (\vec{\nabla} \times \vec{F}) = 0$ (i.e. $\operatorname{div}(\operatorname{curl} \vec{F}) = 0$)

Remark: $\vec{\nabla} \cdot (\vec{\nabla} f) \neq 0$ in general, and it is called the

Laplacian of f , and is denoted by

$$\vec{\nabla}^2 f = \vec{\nabla} \cdot (\vec{\nabla} f) = \operatorname{div}(\vec{\nabla} f) = \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2}$$

[In graduate level, it will be denoted by $\Delta = \vec{\nabla}^2$ or $\Delta = -\vec{\nabla}^2$]

The "operator" $\vec{\nabla}^2$ is called the Laplace operator and the equation $\vec{\nabla}^2 f = 0$ is called the Laplace equation.

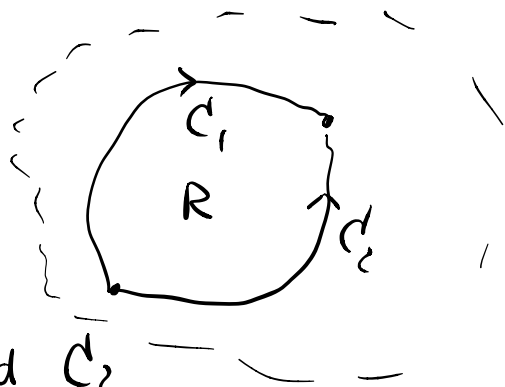
Solutions to Laplace equation are called harmonic functions.

Pf of Thm 10 ($n=2$)

We only need to show Ω simply-connected & $\vec{\nabla} \times \vec{F} = 0$ ($\frac{\partial M}{\partial y} = \frac{\partial N}{\partial x}$)
then \vec{F} is conservative.

Suppose $C_1, C_2 \subseteq \Omega$ have the same starting point and end point.

Case 1: C_1, C_2 have no intersection



Then " Ω is simply-connected"

\Rightarrow the region R enclosed by C_1 and C_2

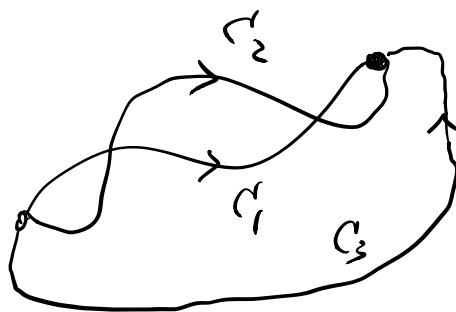
lies completely inside Ω . Then by Green's Thm,

$$0 = \iint_R \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA = \pm \left(\int_{C_1} - \int_{C_2} \right) (Mdx + Ndy)$$

$$\Rightarrow \int_{C_1} Mdx + Ndy = \int_{C_2} Mdx + Ndy$$

Case 2: C_1, C_2 intersect

Pick another curve C_3 with the same starting point and end point, and do not intersect C_1 or C_2 .



$$\begin{aligned} \text{Then by Case 1, } \int_{C_1} Mdx + Ndy &= \int_{C_3} Mdx + Ndy \\ &= \int_{C_2} Mdx + Ndy \end{aligned}$$

$\therefore \int_C \vec{F} \cdot d\vec{r}$ is independent of the path and hence conservative $\#$

In order to apply Green's Thm to more general situations, we need a more general form of Green's Thm:

Suppose we have a simple closed curve C in \mathbb{R}^2



Suppose C_1, C_2, \dots, C_n be pairwise disjoint, piecewise smooth, simple closed curves, such that C_1, \dots, C_n are enclosed by C .

(All C, C_1, \dots, C_n are anti-clockwise oriented)

Let R be the region between C and C_1, \dots, C_n .

Suppose $\vec{F} = M\hat{i} + N\hat{j}$ is defined on some open set containing R and is C^1 . Then we have

$$\iint_R \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA = \oint_C M dx + N dy - \sum_{i=1}^n \oint_{C_i} M dx + N dy$$

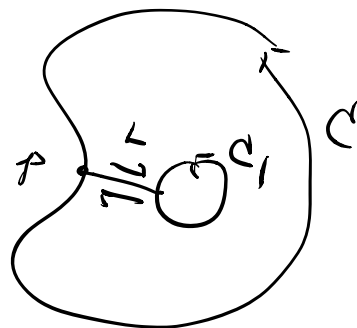
(This is the tangential form. The normal form is similar.)

Sketch of Proof:

For simplicity, only one C_1 inside C .

We connect C & C_1 by an "arc" L

and consider the "simple" closed



curve (starting from p):

$$C^* = C + L - C_1 - L$$

then the region R enclosed between C & C_1 ,
is the region enclosed by C^* except the arc L

$$\text{Hence } \iint_R \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA = \iint_{R \setminus L} \left(\frac{\partial N}{\partial x} - \frac{\partial M}{\partial y} \right) dA$$

$$\stackrel{\text{Green's}}{=} \oint_{C^*} M dx + N dy$$

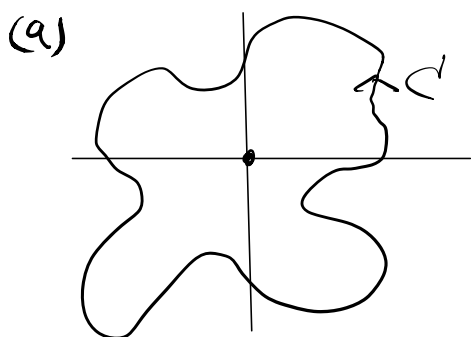
$$= \left(\oint_C + \int_L - \oint_{C_1} - \int_L \right) (M dx + N dy)$$

$$= \oint_C M dx + N dy - \oint_{C_1} M dx + N dy \quad \cancel{\int_L}$$

$$\text{eg 4.8: } \vec{F} = \frac{-y}{x^2+y^2} \hat{i} + \frac{x}{x^2+y^2} \hat{j} \text{ on } \mathbb{R}^2 \setminus \{(0,0)\} = \Omega$$

we've calculated $\oint_{C_1} \vec{F} \cdot d\vec{r} = 2\pi$ for $C_1 = x^2 + y^2 = 1$
(anti-clockwise)

How about



(b)

