

Lecture 7:

Recap: Analytic Spectral (Fourier) method

1. For bounded domain $[0, 2\pi]$, we introduced Fourier Series method.

Write $u(x) = \sum_{k=0}^{\infty} a_k \cos kx + \sum_{k=1}^{\infty} b_k \sin kx$

Eigenvectors of some differential operators (such as $\frac{d^2}{dx^2}$)

↓
Kill differential operator

↓
Algebraic equation (or ODE)
(by comparing coefficients)

Example: Consider $\frac{d^2y}{dx^2} + \frac{dy}{dx} = \sin x + 2\cos x$ where
 $y(0) = y(2\pi)$ (periodic). ($y: [0, 2\pi] \rightarrow \mathbb{R}$)

Write: $y(x) = a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx)$

$$\sum_{n=1}^{\infty} \left[(nb_n - n^2 a_n) \cos nx - (na_n + n^2 b_n) \sin nx \right] = \sin x + 2\cos x$$

⋮ Compare coefficients

$$b_1 - a_1 = 2$$

$$a_1 + b_1 = -1$$

Example:

Consider: $u_t = u_{xx}$, $x \in [0, 2\pi]$ Such that

$$u(0, t) = u(2\pi, t) \text{ (periodic)}$$

$$u(x, 0) = f(x) \text{ (initial condition)}$$

$$(u: [0, 2\pi] \times \mathbb{R} \rightarrow \mathbb{R})$$

$x \quad t$

Write: $u(x, t) = \sum_{n=0}^{\infty} a_n(t) \cos nx + b_n(t) \sin nx$

$$\sum_{n=0}^{\infty} a_n'(t) \cos nx + b_n'(t) \sin nx = \sum_{n=0}^{\infty} (-n^2) a_n(t) \cos nx + (-n^2) b_n(t) \sin nx$$

Comparing coefficients

$$a_n'(t) = -n^2 a_n(t), \text{ and } b_n'(t) = -n^2 b_n(t) \text{ . (ODE)}$$

2. For unbounded domain $(-\infty, \infty)$, we introduced the Fourier transform method.

Example: Consider the ODE = $\frac{dy}{dt} - 4y = H(t)e^{-4t}$ where $t \in \mathbb{R}$
and $H(t)$ is given by: $H(t) = \begin{cases} 0 & \text{if } t < 0 \\ 1 & \text{if } t \geq 0. \end{cases}$

Apply Fourier Transform on both sides:

$$\widehat{\frac{dy}{dt}}(k) - 4\hat{y}(k) = \widehat{H(t)e^{-4t}}(k) \quad \left(= \frac{1}{4+ik} \right)$$

$$(ik-4)\hat{y}(k) = \frac{1}{4+ik} \quad (\text{Algebraic eqn})$$

Example: We consider: $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ where $c \neq 0$
($u = u(x, t)$ is a function defined on $-\infty < x < \infty$ $t \geq 0$)

Apply Fourier Transform w.r.t. x on both sides:

$$\frac{d^2}{dt^2} \hat{u}(k, t) = -c^2 k^2 \hat{u}(k, t)$$

(ODE if fixing k)

Question: Extension to discrete case (Computational Math.)

Answer: Discrete Fourier Transform

Goal: ① Define discrete Fourier Transform (DFT)

② Use DFT to solve discretized differential eqt.

Definition: (Discrete Fourier Transform) Given $f_0, f_1, \dots, f_{n-1} \in \mathbb{C}$,

then the discrete Fourier Transform (DFT) is defined as:

$$\vec{c} = \begin{pmatrix} c_0 \\ c_1 \\ \vdots \\ c_{n-1} \end{pmatrix} \in \mathbb{C}^n \quad \text{where} \quad c_k = \frac{1}{n} \sum_{j=0}^{n-1} f_j e^{-i\left(\frac{2jk\pi}{n}\right)} \quad \text{for } k=0, 1, 2, \dots, n-1$$

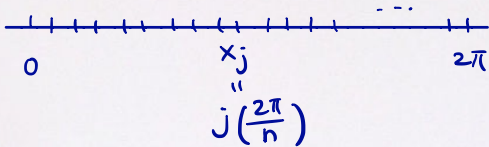
The inverse discrete Fourier Transform recovers the original signal:

$$f_j = \sum_{k=0}^{n-1} c_k e^{i\left(\frac{2jk\pi}{n}\right)} \quad \text{for } j=0, 1, 2, \dots, n-1$$

Motivation of the definition:

Let $f(x)$ be defined on $[0, 2\pi]$. Approximate $f(x)$ by:

$$F_n(x) = \sum_{k=0}^{n-1} c_k e^{ikx}, \quad x \in [0, 2\pi] \quad \text{such that:}$$

$$F_n(x_j) = f(x_j) := f_j, \quad x_j = j \left(\frac{2\pi}{n} \right)$$


Then, we have:

$$(*) \left\{ \begin{array}{l} F_n(x_0) = c_0 + c_1 + c_2 + \dots + c_{n-1} = f_0 \\ F_n(x_1) = c_0 + c_1 e^{ix_1} + c_2 e^{i2x_1} + \dots + c_{n-1} e^{(n-1)ix_1} = f_1 \\ \vdots \\ F_n(x_{n-1}) = c_0 + c_1 e^{ix_{n-1}} + \dots + c_{n-1} e^{(n-1)ix_{n-1}} = f_{n-1} \end{array} \right.$$

Let $\omega = e^{\frac{2\pi i}{n}} = e^{ix_1}$. Then: $\omega^2 = e^{\frac{4\pi i}{n}} = e^{ix_2}$, $\omega^3 = e^{ix_3}$ etc...

$\therefore (*)$ can be written as:

$$A\omega = \begin{pmatrix} 1 & \omega & \omega^2 & \dots & \omega^{n-1} \\ 1 & \omega^2 & \omega^4 & \dots & \omega^{2(n-1)} \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ 1 & \omega^{n-1} & \omega^{2(n-1)} & \dots & \omega^{(n-1)^2} \end{pmatrix} \begin{pmatrix} c_0 \\ c_1 \\ \vdots \\ c_{n-1} \end{pmatrix} = \begin{pmatrix} f_0 \\ f_1 \\ \vdots \\ f_{n-1} \end{pmatrix}$$

Note: $1 + \omega + \omega^2 + \dots + \omega^{n-1} = \frac{1 - \omega^n}{1 - \omega} = 0$ ($\because (e^{\frac{2\pi i}{n}})^n = \omega^n = 1$)

$$\begin{aligned} \therefore (A\omega \bar{A}\omega)_{j,k} &= 1 \cdot 1 + \omega^j \bar{\omega}^k + \omega^{2j} \bar{\omega}^{2k} + \dots + \omega^{(n-1)j} \bar{\omega}^{(n-1)k} \\ &= 1 + e^{\frac{2\pi i(j-k)}{n}} + \dots + e^{\frac{2\pi i(n-1)(j-k)}{n}} = \begin{cases} n & \text{if } j=k \\ \frac{1 - e^{\frac{2\pi i(j-k)}{n}}}{1 - e^{\frac{2\pi i(j-k)}{n}}} = 0 & \text{if } j \neq k \end{cases} \end{aligned}$$

$$\therefore A_w \overline{A_w} = \overline{A_w} A_w = n I .$$

$$\text{We have: } \begin{pmatrix} c_0 \\ c_1 \\ \vdots \\ c_{n-1} \end{pmatrix} = A_w^{-1} \begin{pmatrix} f_0 \\ f_1 \\ \vdots \\ f_{n-1} \end{pmatrix} = \frac{\overline{A_w}}{n} \begin{pmatrix} f_0 \\ f_1 \\ \vdots \\ f_{n-1} \end{pmatrix}$$

$$\text{Hence, } c_k = \frac{1}{n} (f_0 + e^{-\frac{2\pi i}{n} k} f_1 + e^{-\frac{4\pi i}{n} k} f_2 + \dots + e^{-\frac{2\pi i}{n} k(n-1)} f_{n-1})$$

" DFT !! $k = 0, 1, 2, \dots, n-1$

Remark: Computational cost for DFT is:

$$\begin{array}{l} n^2 \text{ multiplication} \\ + \\ n(n-1) \text{ addition} \end{array} = \mathcal{O}(n^2)$$

Example: Consider $f(t) = 5 + 2 \cos(t - \frac{\pi}{2}) + 3 \cos(2t)$.

f is 2π -periodic. Divide $[0, 2\pi]$ by 4 partitions. Find the DFT of f (discretized by 4 points).

$$f_0 = f(0) = 8; \quad f_1 = f\left(\frac{2\pi}{4}\right) = 4; \quad f_2 = f\left(\frac{4\pi}{4}\right) = 8; \quad f_3 = f\left(\frac{6\pi}{4}\right) = 0$$

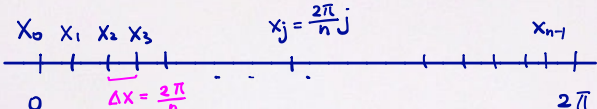
$$\therefore \text{DFT: } C_k = \frac{1}{4} \sum_{j=0}^3 f_j e^{-i\left(\frac{2jk\pi}{4}\right)} \quad \text{for } k=0, 1, 2, 3 \quad \text{or}$$

$$\begin{pmatrix} C_0 \\ C_1 \\ C_2 \\ C_3 \end{pmatrix} = \frac{1}{4} \begin{matrix} | \\ | \\ | \\ | \\ \hline \end{matrix} \begin{pmatrix} f_0 \\ f_1 \\ f_2 \\ f_3 \end{pmatrix} = \frac{1}{4} \begin{pmatrix} 1 & 1 & 1 & 1 \\ 1 & -i & -1 & i \\ 1 & -1 & 1 & -1 \\ 1 & i & -1 & -i \end{pmatrix} \begin{pmatrix} 8 \\ 4 \\ 8 \\ 0 \end{pmatrix} = \begin{pmatrix} 5 \\ -i \\ 3 \\ i \end{pmatrix}$$

Another interpretation of DFT

Recall: Fourier Transform is the extension of Fourier series to $(-\infty, \infty)$.

Fourier coefficient is given by: $C_k = \frac{1}{2\pi} \int_0^{2\pi} f(x) e^{-ikx} dx$

Divide:  2π -periodic

We can approximate the integration by:

$$\begin{aligned} C_k &\approx \frac{1}{2\pi} \sum_{j=0}^{n-1} f(x_j) e^{-ix_j k} \Delta x = \frac{1}{2\pi} \sum_{j=0}^{n-1} f(x_j) e^{-i \frac{2\pi}{n} j k} \left(\frac{2\pi}{n}\right) \\ &= \frac{1}{n} \sum_{j=0}^{n-1} f_j e^{-i \left(\frac{2j k \pi}{n}\right)} \quad \text{for } k=0, 1, 2, \dots, n-1 \end{aligned}$$

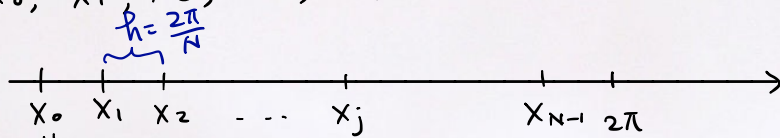
\therefore DFT = approximation of (complex) Fourier coefficient.

DFT and numerical differential equation

Consider: $\frac{d^2 u}{dx^2} = f$ for $x \in [0, 2\pi]$ with periodic boundary condition.
 $u(0) = u(2\pi)$

Suppose f is measured only at N discrete points =

$$x_0, x_1, x_2, \dots, x_{N-1}$$



Let $\vec{f} = \begin{pmatrix} f(x_0) \\ f(x_1) \\ \vdots \\ f(x_{N-1}) \end{pmatrix} \stackrel{\text{def}}{=} \begin{pmatrix} f_0 \\ f_1 \\ \vdots \\ f_{N-1} \end{pmatrix} \in \mathbb{R}^N$ and $\vec{u} = \begin{pmatrix} u(x_0) \\ u(x_1) \\ \vdots \\ u(x_{N-1}) \end{pmatrix} = \begin{pmatrix} u_0 \\ u_1 \\ \vdots \\ u_{N-1} \end{pmatrix} \in \mathbb{R}^N$ (unknown)

By Taylor's expansion,

$$u(x_j + h) \approx u(x_j) + h u'(x_j) + \frac{h^2}{2} u''(x_j) \quad \text{--- (1)}$$

$$u(x_j - h) \approx u(x_j) - h u'(x_j) + \frac{h^2}{2} u''(x_j) \quad \text{--- (2)}$$

$$\textcircled{1} + \textcircled{2} : \quad u''(x_j) \approx \frac{u(x_{j-1}) - 2u(x_j) + u(x_{j+1}))}{h^2}$$

(Central difference approximation of u'')

$$\therefore u''(x_j) \approx \frac{u_{j-1} - 2u_j + u_{j+1}}{h^2} \quad \text{where } h = \frac{2\pi}{N}$$

Thus, $\begin{pmatrix} u''(x_0) \\ u''(x_1) \\ \vdots \\ u''(x_{N-1}) \end{pmatrix} \approx \tilde{D} \vec{u}$ where $\tilde{D} = \frac{1}{h^2} \begin{pmatrix} -2 & & & & \\ 1 & -2 & & & \\ & 1 & \ddots & & \\ & & \ddots & \ddots & \\ & & & 1 & -2 \\ & & & & 1 \end{pmatrix}$

$M_{N \times N}(\mathbb{R})$

(using the fact $u_0 = u_N, u_{-1} = u_{N-1}$)

Thus, $\frac{d^2u}{dx^2} = f$ can be discretized as $\tilde{D} \vec{u} = \vec{f}$ (Linear System)

Remark: Note that \tilde{D} can be a very BIG matrix.

Goal: Design the numerical spectral method. We need to:

① Determine eigenvalues / eigenvectors of \tilde{D}

② Rank of \tilde{D} (to understand the sol. of linear system)

In continuous case, e^{ikx} is an eigenfunction of $\frac{d^2}{dx^2}$, that is periodic.

In discrete case, define:

$$\underbrace{e^{ikx}}_{\mathbb{C}^N} \stackrel{\text{def}}{=} \begin{pmatrix} e^{ikx_0} \\ e^{ikx_1} \\ \vdots \\ e^{ikx_{N-1}} \end{pmatrix}$$

(Capture the values of e^{ikx} at N discrete points)

Claim: $\overrightarrow{e^{ikx}}$ is an eigenvector of \tilde{D}

Proof: For each j , $(\tilde{D} \overrightarrow{e^{ikx}})_j = \frac{e^{ikx_{j-1}} - 2e^{ikx_j} + e^{ikx_{j+1}}}{h^2}$

$$= e^{ikx_j} \left(\frac{e^{-ikh} - 2 + e^{ikh}}{h^2} \right)$$

$$1 - 2\sin^2 \frac{kh}{2}$$

$$= e^{ikx_j} \left(\frac{2\cos^2 \frac{kh}{2} - 2}{h^2} \right)$$

$$= e^{ikx_j} \left(\frac{\cancel{\cos kh} - i\cancel{\sin kh} - 2 + \cancel{\cos kh} + i\cancel{\sin kh}}{h^2} \right)$$

$$= \left(-\frac{4\sin^2 \frac{kh}{2}}{h^2} \right) e^{ikx_j}$$

Let $-\lambda_k^2 = \left(-\frac{4\sin^2 \frac{kh}{2}}{h^2} \right)$. Then: $\tilde{D} \overrightarrow{e^{ikx}} = -\lambda_k^2 \overrightarrow{e^{ikx}}$