

IV. The Reciprocal Lattice

- For every lattice in real space, there is a corresponding reciprocal lattice. Here, we introduce the reciprocal lattice as a mathematical construction. You will find that it is a very useful idea, for handling periodic structures.
- Recall: The primitive vectors  $\vec{a}_1, \vec{a}_2, \vec{a}_3$  define a lattice (one of the Bravais lattices)

A. Constructing the Reciprocal Lattice

Given  $\vec{a}_1, \vec{a}_2, \vec{a}_3$  defining a lattice, we can construct the following 3 vectors:

$$\begin{aligned} \vec{b}_1 &= \frac{2\pi}{\vec{a}_2 \times \vec{a}_3} \\ \vec{b}_2 &= \frac{2\pi}{\vec{a}_3 \times \vec{a}_1} \\ \vec{b}_3 &= \frac{2\pi}{\vec{a}_1 \times \vec{a}_2} \end{aligned}$$

Note:  
The denominator  $\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3) = \Omega_0$   
= volume of primitive unit cell

+ Refs: Ch. 2 of Kittel. Sec. 4.2 of Christman.  
In Chemistry books, the factor  $2\pi$  may be missing.

- Taking  $\vec{b}_1, \vec{b}_2, \vec{b}_3$  as the primitive vectors to generate a lattice, the resulting lattice is the reciprocal lattice of the lattice generated by  $\vec{a}_1, \vec{a}_2, \vec{a}_3$ .
- $\vec{a}_1, \vec{a}_2, \vec{a}_3$  have the dimension of length (e.g. Å), thus  $\vec{b}_1, \vec{b}_2, \vec{b}_3$  have the dimension of 1/length (e.g. Å<sup>-1</sup>).
- ∴ For a 3D lattice in real space, the reciprocal lattice is a 3D lattice in a space of linear dimension length. The space is called the reciprocal space or sometimes k space.
- The reciprocal lattice is the "dual lattice" of the lattice generated by  $\vec{a}_1, \vec{a}_2, \vec{a}_3$  (sometimes called direct lattice).
- For 2D lattices, we have  $\vec{a}_1, \vec{a}_2$  that define the plane of the system (say x-y plane). Taking  $\vec{a}_3 = \hat{z}$ , normal to the 2D plane, then  $\vec{b}_1, \vec{b}_2$  generate a 2D reciprocal lattice.

+ In waves,  $\vec{k}$  is the wavevector.  $|\vec{k}| \sim 1/\lambda$  has the dimension 1/length, and thus the name k space.

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some properties:

$$\vec{a}_1 \cdot \vec{b}_1 = \vec{a}_2 \cdot \vec{b}_2 = \vec{a}_3 \cdot \vec{b}_3 = 2\pi$$

check:  $\vec{a}_1 \cdot \vec{b}_1 = \vec{a}_1 \cdot \frac{2\pi (\vec{a}_2 \times \vec{a}_3)}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = 2\pi,$

similarly for the other cases.

$$\vec{a}_i \cdot \vec{b}_j = 0 \text{ if } i \neq j$$

⊥ plane defined by  $\vec{a}_i, \vec{a}_j$ 

check:  $\vec{a}_1 \cdot \vec{b}_2 = 2\pi \frac{\vec{a}_1 \cdot (\vec{a}_3 \times \vec{a}_1)}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = 0,$

similarly for the other cases.

Thus, we have

$$\vec{a}_i \cdot \vec{b}_j = 2\pi \delta_{ij} \leftarrow \text{a very useful property}$$

Recall:  $\Omega_c = |\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)| = \text{Volume of primitive unit cell in real space.}$

$\therefore \mathcal{V}_G \equiv \text{Volume of the reciprocal lattice primitive cell}$   
 $= |\vec{b}_1 \cdot (\vec{b}_2 \times \vec{b}_3)|$

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 $\Omega_c$  and  $\mathcal{V}_G$  are related by:

$$\mathcal{V}_G = \frac{(2\pi)^3}{\Omega_c}$$

↖ another useful property

check:  $\vec{b}_2 \times \vec{b}_3 = \frac{(2\pi)^2}{(\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3))^2} (\vec{a}_3 \times \vec{a}_2) \times (\vec{a}_1 \times \vec{a}_2)$

$$= \frac{(2\pi)^2}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} \vec{a}_1$$

$$[\vec{A} \times (\vec{B} \times \vec{C}) = \vec{B}(\vec{A} \cdot \vec{C}) - \vec{C}(\vec{A} \cdot \vec{B})]$$

$$\therefore \mathcal{V}_G = \vec{b}_1 \cdot (\vec{b}_2 \times \vec{b}_3)$$

$$= \frac{(2\pi)^3}{\vec{a}_1 \cdot (\vec{a}_2 \times \vec{a}_3)} = \frac{(2\pi)^3}{\Omega_c}$$

- The reciprocal lattice of the reciprocal lattice is the direct lattice.

Why? One of the primitive vectors of the reciprocal lattice of a reciprocal lattice can be constructed by

$$\frac{2\pi \vec{b}_2 \times \vec{b}_3}{\vec{b}_1 \cdot (\vec{b}_2 \times \vec{b}_3)} = \frac{2\pi}{\vec{b}_1 \cdot (\vec{b}_2 \times \vec{b}_3)} \cdot \frac{(2\pi)^2 \vec{a}_1}{(2\pi)^3 \vec{a}_1} = \vec{a}_1 \quad (\text{primitive vectors of direct lattice})$$

Similarly,

$$\frac{2\pi \vec{b}_3 \times \vec{b}_1}{\vec{b}_2 \cdot (\vec{b}_2 \times \vec{b}_3)} = \vec{a}_2$$

$$\frac{2\pi \vec{b}_1 \times \vec{b}_2}{\vec{b}_3 \cdot (\vec{b}_2 \times \vec{b}_3)} = \vec{a}_3$$

### B. Reciprocal Lattice Vectors $\vec{G}$

Reciprocal lattice vectors are defined as:

$$\vec{G} = n_1 \vec{b}_1 + n_2 \vec{b}_2 + n_3 \vec{b}_3$$

where  $n_1, n_2, n_3 = 0, \pm 1, \pm 2, \pm 3, \dots$   
integers

These are the "position vectors" of the points forming the reciprocal lattice.† The set  $\{\vec{G}\}$  generates all the points in a reciprocal lattice.

#### Key Property 1

- Let  $\vec{R}$  be a lattice vector in the direct lattice, and  $\vec{G}$  be a reciprocal lattice vector in the reciprocal lattice.

We have 
$$e^{i\vec{G} \cdot \vec{R}} = 1$$

[See appendix for an immediate application]

†  $\vec{G}$  is the analog of  $\vec{R}$  in direct lattice, only that  $\vec{G}$  lives in the reciprocal lattice

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It is because:

$$\begin{aligned} \vec{G} \cdot \vec{R} &= (n_1 \vec{b}_1 + n_2 \vec{b}_2 + n_3 \vec{b}_3) \cdot (u_1 \vec{a}_1 + u_2 \vec{a}_2 + u_3 \vec{a}_3) \\ &= (n_1 u_1 + n_2 u_2 + n_3 u_3) \cdot 2\pi \\ &= \text{integer} \cdot 2\pi \quad (\text{see notes}) \\ \therefore e^{i\vec{G} \cdot \vec{R}} &= 1 \end{aligned}$$

Key Property 2

- Distances between adjacent parallel crystal planes in direct lattice

$$\Gamma \text{ Recall: } d = \frac{\Omega_c}{|\vec{q}|}$$

$$\text{where } \vec{q} = h(\vec{a}_2 \times \vec{a}_3) + k(\vec{a}_3 \times \vec{a}_1) + l(\vec{a}_1 \times \vec{a}_2)$$

$h, k, l$  have no common divisors

$(hkl)$  specifies a set of parallel planes

$\vec{q}$  is normal to the planes  $(hkl)$

□

+ It follows that for some vector  $\vec{G}$ , if  $\vec{G} \cdot \vec{R} = 2\pi$  (integer) for ALL lattice vectors  $\vec{R}$ ,  $\vec{G}$  is a reciprocal lattice vector.

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- From the definition of  $\vec{b}_1, \vec{b}_2, \vec{b}_3$ , the vector  $\vec{G}(hkl) = h\vec{b}_1 + k\vec{b}_2 + l\vec{b}_3$  is normal to the lattice planes with Miller indices  $(hkl)$ .
- $\vec{q}$  and  $\vec{G}$  only differ by a constant.

$$|\vec{G}| = \frac{2\pi}{\Omega_c} |\vec{q}| \quad (\text{see definition of } \vec{b}_1, \vec{b}_2, \vec{b}_3)$$

It follows that

$$d = \frac{\Omega_c}{|\vec{q}|} = \frac{\Omega_c}{2\pi |\vec{G}|} = \frac{2\pi}{|\vec{G}|}$$

$\therefore$

$$d(hkl) = \frac{2\pi}{|\vec{G}(hkl)|}$$

where  $d(hkl)$  = separation between adjacent crystal

planes with Miller indices  $(hkl)$

$$\text{and } \vec{G}(hkl) = h\vec{b}_1 + k\vec{b}_2 + l\vec{b}_3$$

C. The First Brillouin Zone (1st B.Z. or B.Z.) IV-9

- (1) Take a point in the reciprocal lattice as origin.
- (2) The 1st B.Z. is the volume in the reciprocal lattice that is closer to the origin than other points in the reciprocal lattice, i.e. the 1st B.Z. is the Wigner-Seitz cell centered on the origin of the reciprocal lattice. +

Examples:

(a) Direct lattice: Simple Cubic lattice

$$\vec{a}_1 = a \hat{x}, \quad \vec{a}_2 = a \hat{y}, \quad \vec{a}_3 = a \hat{z}, \quad \Omega_0 = a^3$$

$$\vec{b}_1 = \frac{2\pi}{a^3} \vec{a}_2 \times \vec{a}_3 = \frac{2\pi}{a^3} \cdot a^2 \hat{x} = \frac{2\pi}{a} \hat{x}$$

$$\vec{b}_2 = \frac{2\pi}{a} \hat{y}, \quad \vec{b}_3 = \frac{2\pi}{a} \hat{z}$$

$$V_{\text{BZ}} = \frac{(2\pi)^3}{a^3} = \text{Volume of reciprocal lattice primitive cell} \\ = \text{Volume of BZ.}$$

+ Recall the construction of the Wigner-Seitz cell.

In direct lattices, we call it the Wigner-Seitz cell. In reciprocal lattice, we call it the construction the 1st B.Z.. The 1st B.Z. is a primitive cell in the reciprocal lattice.

• With  $\vec{b}_1, \vec{b}_2, \vec{b}_3$ , we see that:

The reciprocal lattice to simple cubic lattice (in real space) is also a simple cubic lattice. +

The lattice constant of the reciprocal lattice is  $\frac{2\pi}{a}$ .

• The reciprocal lattice vectors are given by:

$$\vec{G} = \frac{2\pi}{a} (v_1 \hat{x} + v_2 \hat{y} + v_3 \hat{z})$$

• Separation between adjacent planes (hkl) in direct lattice:

$$d = \frac{2\pi}{|\vec{G}(hkl)|} \quad \text{where } \vec{G} = h\vec{b}_1 + k\vec{b}_2 + l\vec{b}_3 \\ = \frac{2\pi}{a} (h\hat{x} + k\hat{y} + l\hat{z})$$

$$= \frac{2\pi}{a \sqrt{(h^2 + k^2 + l^2)}} \\ = \frac{2\pi}{\sqrt{h^2 + k^2 + l^2}}$$

same as before. (see Ch. II)

• The 1st B.Z. is a cube of edge length  $\frac{2\pi}{a}$  with the origin at the centre of the cube.

+ This is sometimes referred to as "self-dual" of simple cubic lattices, i.e., the "dual lattice" of SC is SC.

(b) Direct Lattice: Body Centered Cubic

For bcc,  $\vec{a}_1 = \frac{a}{2}(\hat{x} + \hat{y} - \hat{z})$ ;  $\vec{a}_2 = \frac{a}{2}(-\hat{x} + \hat{y} + \hat{z})$ ;

$\vec{a}_3 = \frac{a}{2}(\hat{x} - \hat{y} + \hat{z})$  (edge of cube = a)

This gives:  $\vec{a}_1 = \frac{a}{2}(\hat{x} + \hat{y})$   
 $\vec{a}_2 = \frac{a}{2}(\hat{y} + \hat{z})$   
 $\vec{a}_3 = \frac{a}{2}(\hat{z} + \hat{x})$

Note:  $[\vec{a}_1, \vec{a}_2, \vec{a}_3]$  is the edge of conventional unit cell. These are the primitive vectors of a fcc lattice of cube edge of  $(\frac{a}{2})$ .

The reciprocal lattice to bcc (in real space) is a fcc.

There are 12 nearest neighbors to a point. The 12 shortest reciprocal lattice vectors  $\vec{G}$ 's are:

$\frac{2\pi}{a}(\pm\hat{x} \pm \hat{y})$ ;  $\frac{2\pi}{a}(\pm\hat{y} \pm \hat{z})$ ;  $\frac{2\pi}{a}(\pm\hat{z} \pm \hat{x})$

The 1st BZ is a rhombic dodecahedron† (12-faces).

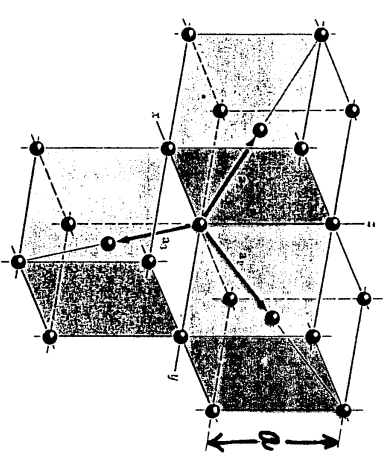
(See Figure)

Reciprocal lattice vectors:

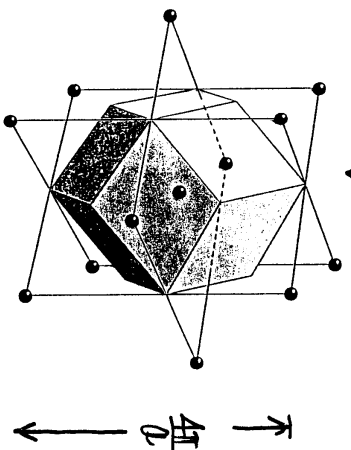
$\vec{G} = v_1 \vec{b}_1 + v_2 \vec{b}_2 + v_3 \vec{b}_3$   
 $= \frac{2\pi}{a} [(v_1 + v_2)\hat{x} + (v_1 + v_3)\hat{y} + (v_2 + v_3)\hat{z}]$

+ This is also the Wigner-Seitz cell of a fcc.

bcc in direct lattice (e.g. simple metals)



Primitive basis vectors of the body-centered cubic lattice.



First Brillouin zone of the body-centered cubic lattice. The figure is a regular rhombic dodecahedron.

Reciprocal lattice to bcc is a fcc.

The shaded volume is the 1st B.Z.

$\Omega_c = \frac{a^3}{2}$  (Volume of primitive cell)

$\Omega_{G1} = \left(\frac{2\pi}{a}\right)^3$   
 $= 2 \left(\frac{2\pi}{a}\right)^3 =$  Volume of primitive cell of reciprocal lattice  
 $=$  Volume of 1st B.Z.

See's Reciprocal Lattices: FCC with edge  $(\frac{a}{2})$

4 reciprocal lattice points in this cube  
 $\Omega_{G1} = \frac{1}{4} \left(\frac{4\pi}{a}\right)^3 = 2 \left(\frac{2\pi}{a}\right)^3$

IV-13 (c) Direct lattice: Face centered cubic (e.g. semiconductors)

For fcc,  $\vec{a}_1 = \frac{a}{2}(\hat{x} + \hat{y})$ ;  $\vec{a}_2 = \frac{a}{2}(\hat{y} + \hat{z})$ ;  $\vec{a}_3 = \frac{a}{2}(\hat{z} + \hat{x})$   
 (edge of cube = a)

This gives:  $\vec{b}_1 = \frac{\sqrt{2}}{a}(\hat{x} + \hat{y} - \hat{z})$   
 $\vec{b}_2 = \frac{\sqrt{2}}{a}(-\hat{x} + \hat{y} + \hat{z})$   
 $\vec{b}_3 = \frac{\sqrt{2}}{a}(\hat{x} - \hat{y} + \hat{z})$

These are the primitive vectors of a bcc lattice of cube edge of  $(\frac{\sqrt{2}}{a})$ .

The reciprocal lattice to fcc (in real space) is a bcc.

There are 8 nearest neighbors to a point.

The 8 shortest reciprocal lattice vectors  $\vec{G}_i$ 's are:

$$\frac{\sqrt{2}}{a}(\pm\hat{x} \pm \hat{y} \pm \hat{z}).$$

The 1st B.Z. is a truncated octahedron. †

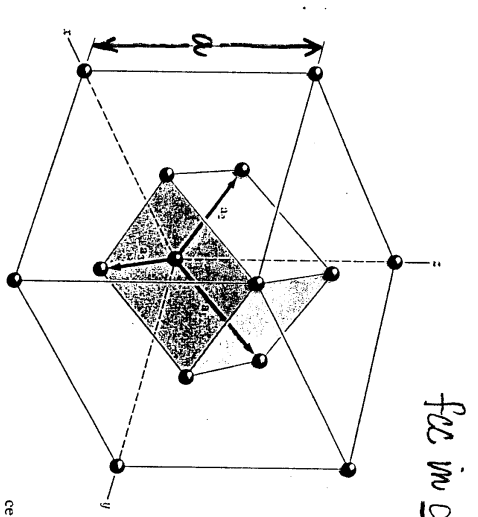
(See Figure).

Reciprocal lattice vectors:

$$\vec{G} = n_1 \vec{b}_1 + n_2 \vec{b}_2 + n_3 \vec{b}_3$$

$$= \frac{\sqrt{2}}{a} [(n_1 - n_2 + n_3)\hat{x} + (n_1 + n_2 - n_3)\hat{y} + (-n_1 + n_2 + n_3)\hat{z}]$$

+ This is also the Wigner-Seitz of a bcc.



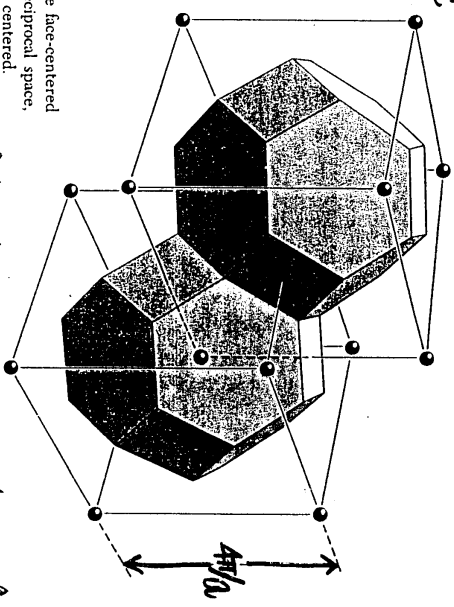
fcc in direct lattice

$$\Omega_c = \frac{a^3}{4}$$

Primitive basis vectors of the face-centered cubic lattice.

Reciprocal lattice to fcc is a bcc

Brillouin zones of the face-centered cubic lattice. The cells are in reciprocal space, and the reciprocal lattice is body centered.



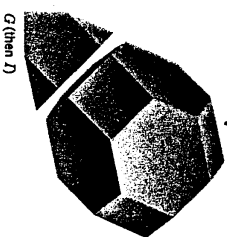
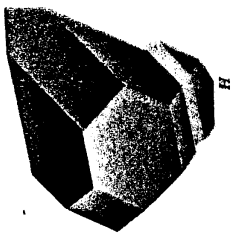
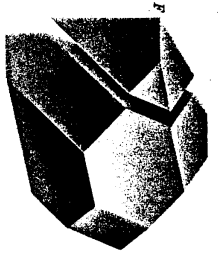
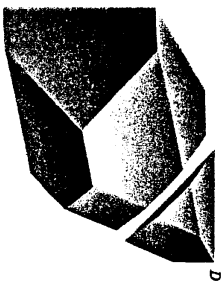
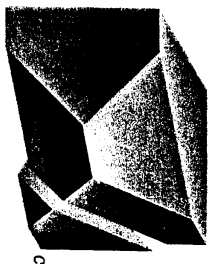
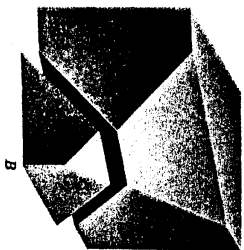
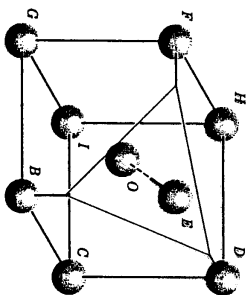
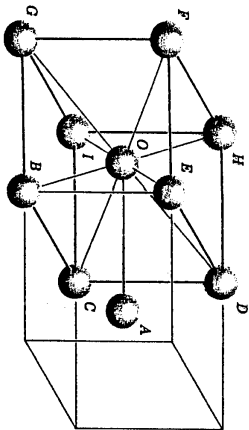
The shaded volume is the 1st B.Z. †

$$\Omega_{G_1} = \frac{(4\pi)^3}{\Omega_c} = 4 \left(\frac{4\pi}{a}\right)^3$$

+ This is probably the most useful 1st BZ in solid state physics. The reason is that many semiconductors take on the fcc structure, when we consider the band structure E(k),  $k_0$  lies in 1st BZ.

Construction of the Proximity cell of a bcc lattice

+ This could be the Wigner-Seitz cell of a bcc direct lattice, or the first Brillouin zone of a fee direct lattice for which the reciprocal lattice is a bcc.



Final object

Summary

1. Given a direct lattice (hence  $\vec{a}_1, \vec{a}_2, \vec{a}_3$ ), we can construct  $\vec{b}_1, \vec{b}_2, \vec{b}_3$ .
2.  $\vec{b}_1, \vec{b}_2, \vec{b}_3$  are the primitive vectors of the reciprocal lattice.
3.  $\vec{a}_i \cdot \vec{b}_j = 2\pi \delta_{ij}$
4.  $\Omega_G = \frac{(2\pi)^3}{\Omega_G}$
5.  $\vec{G} = n_1 \vec{b}_1 + n_2 \vec{b}_2 + n_3 \vec{b}_3$ , where  $n_1, n_2, n_3 = 0, \pm 1, \pm 2, \dots$  are the reciprocal lattice vectors.
6.  $e^{i\vec{G} \cdot \vec{R}} = 1$  since  $\vec{G} \cdot \vec{R} = 2\pi \cdot \text{integer}$
7.  $d(hkl) = \frac{2\pi}{|\vec{G}(hkl)|}$   
 where  $\vec{G}(hkl) = h\vec{b}_1 + k\vec{b}_2 + l\vec{b}_3$   
 and  $(hkl)$  are the Miller indices
8. The first Brillouin zone (1<sup>st</sup> BZ) is the Wigner-Seitz cell in the reciprocal lattice.



## 9. Examples

<u>Direct</u>	<u>Reciprocal</u>
sc ( $a$ )	sc ( $\frac{2\pi}{a}$ )
bcc ( $a$ )	fcc ( $\frac{4\pi}{a}$ )
fcc ( $a$ )	bcc ( $\frac{4\pi}{a}$ )

Here  $a$  is the cube edge.

## 10. Why BZ?

- The BZ shows the symmetry properties of the reciprocal lattice, which in turn reflects the symmetry properties of the direct lattice.
- In many problems (due to the discrete translational symmetry of crystals) in crystals, the solutions labelled by wavevectors  $\vec{k}$  and  $\vec{k} + \vec{G}$  (any  $\vec{G}$ ) are redundant. Thus, restricting  $\vec{k}$  to within a primitive unit cell in reciprocal space is sufficient and the 1<sup>st</sup> BZ is usually chosen.