

Similarly, if ω is a 1-form and
 ζ is a 2-form,

then we can define $\omega \wedge \zeta$.

e.g.: If $\omega = dx$, $\zeta = dy \wedge dz$,

$$\begin{aligned}\text{then } \omega \wedge \zeta &= dx \wedge (dy \wedge dz) \\ &= dx \wedge dy \wedge dz\end{aligned}$$

Note that we insist on the anti-commutativity of wedge product, we have

$$\begin{aligned}dx \wedge dy \wedge dz &= - dy \wedge dx \wedge dz \\ &= dy \wedge dz \wedge dx \\ &= - dz \wedge dy \wedge dx \\ &= dz \wedge dx \wedge dy \\ &= - dx \wedge dz \wedge dy\end{aligned}$$

And one can see that, as $\dim \mathbb{R}^3 = 3$, all

"linear combinations" are just

$$f dx \wedge dy \wedge dz$$

which is called a 3-form (also called a volume form)
if $f > 0$

Note: It is convenient to call smooth functions
of the differential 0-form.

Summary ($\bar{\in} \mathbb{R}^3$)

0-form : f

1-form : $w_1 dx + w_2 dy + w_3 dz$

2-form : $\xi_1 dy \wedge dz + \xi_2 dz \wedge dx + \xi_3 dx \wedge dy$

3-form : $g dx \wedge dy \wedge dz$

(where f, g, w_i, ξ_i are functions)

Note: One can certainly define k -form for any $k \geq 0$.

But in \mathbb{R}^3 , k -forms are all zero for $k > 3$:

$dx^i \wedge dx \wedge dy \wedge dz = 0$, where $dx^i = dx, dy, dz$.

Change of Variables Formula : (\mathbb{R}^2)

$$\begin{cases} x = x(u, v) \\ y = y(u, v) \end{cases}$$

$$\Rightarrow \begin{cases} dx = x_u du + x_v dv \\ dy = y_u du + y_v dv \end{cases}$$

$$\begin{aligned} \Rightarrow dx \wedge dy &= (x_u du + x_v dv) \wedge (y_u du + y_v dv) \\ &= (x_u y_v - x_v y_u) du \wedge dv \\ &= \begin{vmatrix} x_u & x_v \\ y_u & y_v \end{vmatrix} du \wedge dv \end{aligned}$$

$$\boxed{dx \wedge dy = \frac{\partial(x, y)}{\partial(u, v)} du \wedge dv}$$

Hence, naturally

$$\boxed{\iint f(x, y) dx \wedge dy = \iint f(x(u, v), y(u, v)) \frac{\partial(x, y)}{\partial(u, v)} du \wedge dv}$$

Compare with

$$\boxed{\iint f(x, y) dx dy = \iint f(x(u, v), y(u, v)) \left| \frac{\partial(x, y)}{\partial(u, v)} \right| du dv}$$

we can interpret " $dx \wedge dy$ " as the "oriented" area element.

Similarly for $\begin{cases} x = x(u, v, w) \\ y = y(u, v, w) \\ z = z(u, v, w) \end{cases}$

$$(Ex!) \quad dx \wedge dy \wedge dz = \frac{\partial(x, y, z)}{\partial(u, v, w)} du \wedge dv \wedge dw$$

(using $dx = x_u du + x_v dv + x_w dw, \dots$)

Oriented Change of Variables Formula

Exterior differentiation "d" on a form " ω "

0-form f	df (1-form)
1-form $\omega = \omega_1 dx + \omega_2 dy + \omega_3 dz$	$d\omega = d\omega_1 \wedge dx + d\omega_2 \wedge dy + d\omega_3 \wedge dz$ (2-form)
2-form $\xi = \xi_1 dy \wedge dz + \xi_2 dz \wedge dx + \xi_3 dx \wedge dy$	$d\xi = d\xi_1 \wedge dy \wedge dz + d\xi_2 \wedge dz \wedge dx + d\xi_3 \wedge dx \wedge dy$ (3-form)

3-fam
 $\int dx \wedge dy \wedge dz$

$df \wedge dx \wedge dy \wedge dz = 0$ (4-fam)

e.g. $(\bar{m} \subset \mathbb{R}^2)$ $\omega = M dx + N dy$ $(M = M(x, y), N = N(x, y))$

then $d\omega = dM \wedge dx + dN \wedge dy$
 $= (M_x dx + M_y dy) \wedge dx + (N_x dx + N_y dy) \wedge dy$
 $= (N_x - M_y) dx \wedge dy$

In this notation, Green's Thm can be written as

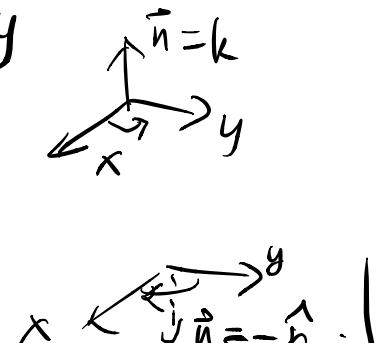
$$\oint_C \omega = \iint_R d\omega .$$

Note that $d\omega = (N_x - M_y) dx \wedge dy$ \curvearrowright is oriented.

If we let $\vec{F} = M\hat{i} + N\hat{j} \leftrightarrow \omega = Mdx + Ndy$

then $(\nabla \times \vec{F}) \cdot \vec{n} dA = (N_x - M_y) \underbrace{\hat{k} \cdot \vec{n}}_{dx \wedge dy} dA = d\omega$

Hence $\hat{k} \cdot \vec{n} dA = \begin{cases} dx \wedge dy, & \text{if } \vec{n} = \hat{k} \\ dy \wedge dx, & \text{if } \vec{n} = -\hat{k} \end{cases}$



$$\text{eg} \quad \varsigma = \varsigma_1 dy \wedge dz + \varsigma_2 dz \wedge dx + \varsigma_3 dx \wedge dy$$

$$\text{then } d\varsigma = d\varsigma_1 \wedge dy \wedge dz + d\varsigma_2 \wedge dz \wedge dx + d\varsigma_3 \wedge dx \wedge dy$$

$$= \left(\frac{\partial \varsigma_1}{\partial x} dx + \dots \right) \wedge dy \wedge dz$$

$$+ \left(\dots + \frac{\partial \varsigma_2}{\partial y} dy + \dots \right) \wedge dz \wedge dx$$

$$+ \left(\dots + \frac{\partial \varsigma_3}{\partial z} dz \right) \wedge dx \wedge dy$$

$$= \left(\frac{\partial \varsigma_1}{\partial x} + \frac{\partial \varsigma_2}{\partial y} + \frac{\partial \varsigma_3}{\partial z} \right) dx \wedge dy \wedge dz$$

$$= (\operatorname{div} \vec{F}) dx \wedge dy \wedge dz$$

$$\text{where } \vec{F} = \overset{\wedge}{\varsigma_1} \hat{i} + \overset{\wedge}{\varsigma_2} \hat{j} + \overset{\wedge}{\varsigma_3} \hat{k}$$

Hence the divergence theorem, we can be written

$$\text{as } \iiint_D d\varsigma = \iiint_D \left(\frac{\partial \varsigma_1}{\partial x} + \frac{\partial \varsigma_2}{\partial y} + \frac{\partial \varsigma_3}{\partial z} \right) dx \wedge dy \wedge dz$$

$$= \iiint_D \operatorname{div} \vec{F} dV$$

$$= \iint_{S=\partial D} \vec{F} \cdot \vec{n} d\sigma$$

To see the relation between $\vec{F} \cdot \vec{n} d\sigma$ and Σ ,
we parametrize S :

$$\vec{r}(u, v) = x(u, v)\hat{i} + y(u, v)\hat{j} + z(u, v)\hat{k}$$

$$\Rightarrow \begin{cases} \vec{r}_u = x_u \hat{i} + y_u \hat{j} + z_u \hat{k} \\ \vec{r}_v = x_v \hat{i} + y_v \hat{j} + z_v \hat{k} \end{cases}$$

$$\Rightarrow \vec{r}_u \times \vec{r}_v = \begin{vmatrix} y_u & y_v \\ z_u & z_v \end{vmatrix} \hat{i} + \begin{vmatrix} z_u & z_v \\ x_u & x_v \end{vmatrix} \hat{j} + \begin{vmatrix} x_u & x_v \\ y_u & y_v \end{vmatrix} \hat{k}$$

If $\vec{r}_u \times \vec{r}_v$ is outward, then

$$\vec{n} = \frac{\vec{r}_u \times \vec{r}_v}{|\vec{r}_u \times \vec{r}_v|} \quad \text{and} \quad d\sigma = |\vec{r}_u \times \vec{r}_v| du dv$$

$$= (\vec{r}_u \times \vec{r}_v) du \wedge dv$$

(correct orientation)

$$\begin{aligned} \vec{F} \cdot \vec{n} d\sigma &= \vec{F} \cdot \frac{\vec{r}_u \times \vec{r}_v}{|\vec{r}_u \times \vec{r}_v|} |\vec{r}_u \times \vec{r}_v| du \wedge dv \\ &= \left(\Sigma_1 \frac{\partial(y, z)}{\partial(u, v)} + \Sigma_2 \frac{\partial(z, x)}{\partial(u, v)} + \Sigma_3 \frac{\partial(x, y)}{\partial(u, v)} \right) du \wedge dv \\ &= \Sigma_1 dy \wedge dz + \Sigma_2 dz \wedge dx + \Sigma_3 dx \wedge dy = \Sigma \end{aligned}$$

$$\therefore \boxed{\iint_D d\zeta = \iint_S \zeta \quad |_{S=\partial D}}$$

eg3 Stokes' Thm

$$\text{If } \vec{F} = M\hat{i} + N\hat{j} + L\hat{k} \Leftrightarrow \omega = Mdx + Ndy + Ldz$$

$$\text{Then } d\omega = (N_y - L_z) dy \wedge dz + (M_z - L_x) dz \wedge dx \quad (\text{Ex!}) \\ + (N_x - M_y) dx \wedge dy$$

Compare to

$$\vec{\nabla} \times \vec{F} = (N_y - L_z)\hat{i} + (M_z - L_x)\hat{j} + (N_x - M_y)\hat{k}$$

$$\therefore d\omega = (\vec{\nabla} \times \vec{F}) \cdot \vec{n} d\sigma$$

Stokes' Thm \Leftrightarrow

$$\oint_C \vec{F} \cdot d\vec{r} = \iint_S (\vec{\nabla} \times \vec{F}) \cdot \vec{n} d\sigma$$

$C = \partial S$ ↓ ↓

$$\boxed{\oint_C \omega = \iint_S d\omega \quad |_{C=\partial S}}$$

Generalization to manifold of n-dimension with boundary

- $M = n$ diiml Manifold (oriented)
- $\partial M = \text{boundary}$ (oriented with induced orientation)
- $\omega = (n-1)\text{-form on } M$ (smooth)

Then

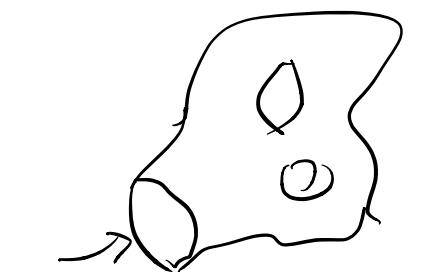
$$\boxed{\int_M d\omega = \int_{\partial M} \omega}$$

\uparrow \uparrow
 n -diiml $(n-1)$ -diiml
integral integral

Note: ∂M is always closed, i.e. no boundary.

$$\therefore \boxed{\partial(\partial M) = \partial^2 M = 0}$$

boundary has no boundary



∂S is a closed curve

Hence if $\omega = d\eta$, for some $(n-2)$ -form η , then

then

$$\int_M d(d\eta) = \int_M d\omega = \int_{\partial M} \omega$$

$$= \int_{\partial M} d\eta = \int_{\partial(\partial M)} \eta = 0 \quad (\text{fr any } \eta.)$$

This suggests

$$\boxed{d^2\eta = 0}, \text{ & differential form}$$

Ex: Verify this for 0-form and 1-form in \mathbb{R}^3
and observes that these are just

$$\left\{ \begin{array}{l} \vec{\nabla} \times \vec{\nabla} f = 0 \quad (d^2f = 0) \\ \vec{\nabla} \cdot (\vec{\nabla} \times \vec{F}) = 0 \quad (d^2\omega = 0) \end{array} \right.$$

e.g.: Let $\omega = \frac{-y}{x^2+y^2} dx + \frac{x}{x^2+y^2} dy$

check: $d\omega = 0$

But $\omega \neq df$ for any smooth function on $\mathbb{R}^2 \setminus \{(0,0)\}$

(Since $\omega = d\theta$ and θ is not defined on $\mathbb{R}^2 \setminus \{(0,0)\}$)

Hence $d\omega = 0 \not\Rightarrow \omega = dy$ in general

$$(\Leftarrow)$$

↑
yes

Note: Then Ω can be written as :

$\Omega \subset \mathbb{R}^2$ simply-connected, then smooth.

$d\omega = 0 \Leftrightarrow \omega = df$ for some \checkmark function
f on Ω .