

Unified treatment of Green's, Stokes', and Divergence Theorems

Stokes' Thm in notations of differential forms (in \mathbb{R}^3)

Working definition of differential forms

(1) A differential 1-form (or simply 1-form)

is a linear combination of the symbols dx, dy & dz :

$$\omega = \omega_1 dx + \omega_2 dy + \omega_3 dz$$

with coefficients $\omega_1, \omega_2, \omega_3$ are functions on \mathbb{R}^3 .

e.g.: The total differential of a smooth function f
is a differential 1-form:

$$df = \frac{\partial f}{\partial x} dx + \frac{\partial f}{\partial y} dy + \frac{\partial f}{\partial z} dz$$

(2) Wedge product: Let " \wedge " be an operation such that

$$\left\{ \begin{array}{l} dx \wedge dx = dy \wedge dy = dz \wedge dz = 0 \\ dx \wedge dy = -dy \wedge dx \\ dy \wedge dz = -dz \wedge dy \\ dz \wedge dx = -dx \wedge dz \end{array} \right.$$

and satisfies other usual rules in arithmetic.

i.e. If $\omega = \omega_1 dx + \omega_2 dy + \omega_3 dz$

$$\eta = \eta_1 dx + \eta_2 dy + \eta_3 dz$$

then we have

$$\begin{aligned}\omega \wedge \eta &= (\omega_1 dx + \omega_2 dy + \omega_3 dz) \wedge (\eta_1 dx + \eta_2 dy + \eta_3 dz) \\&= \cancel{\omega_1 dx} \cancel{\wedge \eta_1 dx} + \omega_2 dy \wedge \eta_1 dx + \omega_3 dz \wedge \eta_1 dx \\&\quad + \omega_1 dx \wedge \cancel{\eta_2 dy} + \cancel{\omega_2 dy} \wedge \eta_2 dy + \omega_3 dz \wedge \eta_2 dy \\&\quad + \omega_1 dx \wedge \eta_3 dz + \omega_2 dy \wedge \eta_3 dz + \cancel{\omega_3 dz} \cancel{\wedge \eta_3 dz} \\&= (\omega_1 \eta_2 - \omega_2 \eta_1) dx \wedge dy \\&\quad + (\omega_2 \eta_3 - \omega_3 \eta_2) dy \wedge dz \\&\quad + (\omega_3 \eta_1 - \omega_1 \eta_3) dz \wedge dx\end{aligned}$$

∴

$$\begin{aligned}\omega \wedge \eta &= (\omega_2 \eta_3 - \omega_3 \eta_2) dy \wedge dz \\&\quad + (\omega_3 \eta_1 - \omega_1 \eta_3) dz \wedge dx \\&\quad + (\omega_1 \eta_2 - \omega_2 \eta_1) dx \wedge dy\end{aligned}$$

- Linear combinations of $dy \wedge dz$, $dz \wedge dx$ & $dx \wedge dy$ are called differential 2-forms (on \mathbb{R}^3)

$$\xi = \xi_1 dy \wedge dz + \xi_2 dz \wedge dx + \xi_3 dx \wedge dy$$

Similarly, if ω is a 1-form and
 ξ is a 2-form

then we can define $\omega \wedge \xi$

Eg: If $\omega = dx$, $\zeta = dy \wedge dz$

then $\omega \wedge \zeta = dx \wedge dy \wedge dz$

Note that we insist on the anti-commutativity of wedge product, we have

$$\begin{aligned} dx \wedge dy \wedge dz &= -dy \wedge dx \wedge dz \\ &= dy \wedge dz \wedge dx \\ &= -dz \wedge dy \wedge dx \\ &= dz \wedge dx \wedge dy \\ &= -dx \wedge dz \wedge dy \end{aligned}$$

And $dx \wedge dx \wedge dy = \dots = 0$ whenever one of the dx, dy, dz repeated.

Hence, as $\dim \mathbb{R}^3 = 3$, all "linear combinations" of "3-fams"

are just $f dx \wedge dy \wedge dz$

which is called a differential 3-form (also called a volume form if $f > 0$)

Note: It is convenient to call smooth functions f the differential 0-form.

Summary (on \mathbb{R}^3)

$$0\text{-form} = f$$

$$1\text{-form} = \omega_1 dx + \omega_2 dy + \omega_3 dz$$

$$2\text{-form} = \varsigma_1 dy \wedge dz + \varsigma_2 dz \wedge dx + \varsigma_3 dx \wedge dy$$

$$3\text{-form} = g dx \wedge dy \wedge dz.$$

where, $f, g, \omega_i, \varsigma_i$ are (smooth) functions

Note = One can certainly define k -form for any $k \geq 0$. But in

\mathbb{R}^3 , k -forms are zero for $k > 3$:

$$dx^i \wedge dx \wedge dy \wedge dz = 0, \text{ where } dx^i = dx, dy, \text{ or } dz.$$

Change of Variables Formula: (\mathbb{R}^2)

$$\begin{cases} x = x(u, v) \\ y = y(u, v) \end{cases}$$

$$\Rightarrow \begin{cases} dx = x_u du + x_v dv \\ dy = y_u du + y_v dv \end{cases}$$

$$\Rightarrow dx \wedge dy = (x_u du + x_v dv) \wedge (y_u du + y_v dv)$$

$$= (x_u y_v - x_v y_u) du \wedge dv$$

$$= \begin{vmatrix} x_u & x_v \\ y_u & y_v \end{vmatrix} du \wedge dv$$

$$dx \wedge dy = \frac{\partial(x, y)}{\partial(u, v)} du \wedge dv$$

↑ Jacobian determinant.

Hence naturally

$$\iint f(x, y) dx \wedge dy = \iint f(x(u, v), y(u, v)) \frac{\partial(x, y)}{\partial(u, v)} du \wedge dv$$

Compare with

$$\iint f(x, y) dx dy = \iint f(x(u, v), y(u, v)) \left| \frac{\partial(x, y)}{\partial(u, v)} \right| du dv$$

Similarly for

$$\begin{cases} x = x(u, v, w) \\ y = y(u, v, w) \\ z = z(u, v, w) \end{cases}$$

$$dx \wedge dy \wedge dz = \frac{\partial(x, y, z)}{\partial(u, v, w)} du \wedge dv \wedge dw \quad (\text{Ex!})$$

(using $dx = x_u du + x_v dv + x_w dw, \dots$)

- “Oriented” change of variables formula

- “ $dx \wedge dy$ ” oriented area element

- “ $dx \wedge dy \wedge dz$ ” oriented volume element.

(see later remark)

Exterior differentiation "d" on a form " ω ".

0-form f	df	(1-form)
1-form $\omega = \omega_1 dx + \omega_2 dy + \omega_3 dz$	$d\omega = d\omega_1 \wedge dx + d\omega_2 \wedge dy + d\omega_3 \wedge dz$	(2-form)
2-form $\zeta = \zeta_1 dy \wedge dz + \zeta_2 dz \wedge dx + \zeta_3 dx \wedge dy$	$d\zeta = d\zeta_1 \wedge dy \wedge dz + d\zeta_2 \wedge dz \wedge dx + d\zeta_3 \wedge dx \wedge dy$	(3-form)
3-form $f dx \wedge dy \wedge dz$	$df \wedge dx \wedge dy \wedge dz = 0$	(4-form) in \mathbb{R}^3

eg. $d(dx) = d(dy) = d(dz) = 0$. ($d^2x = d^2y = d^2z = 0$)

eg 1 (in \mathbb{R}^2) $\omega = M dx + N dy$ ($M = M(x, y), N = N(x, y)$)

then $d\omega = dM \wedge dx + dN \wedge dy$
 $= (M_x dx + M_y dy) \wedge dx + (N_x dx + N_y dy) \wedge dy$
 $= (N_x - M_y) dx \wedge dy$ (+ve) oriented area

In this notation, Green's Thm $\oint_{C=\partial R} M dx + N dy = \iint_R (N_x - M_y) dx dy$,

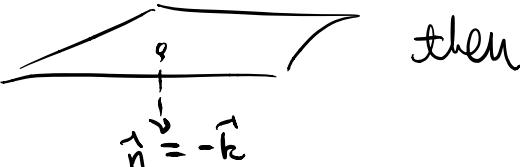
can be written as
$$\left\{ \begin{array}{l} \oint_{C=\partial R} \omega = \iint_R d\omega \end{array} \right\}$$

Remark: If we let $\vec{F} = M\hat{i} + N\hat{j} \leftrightarrow \omega = Mdx + Ndy$

then $(\vec{\nabla} \times \vec{F}) \cdot \hat{n} dA = (N_x - M_y) \underbrace{\hat{k} \cdot \hat{n}}_{dx \wedge dy} dA = d\omega$

$(\hat{n} = \hat{k})$

Hence, if we use



then

$$\hat{k} \cdot \hat{n} dA = \begin{cases} dx \wedge dy & \text{if } \hat{n} = \hat{k} \\ dy \wedge dx & \text{if } \hat{n} = -\hat{k} \end{cases}$$

orientation of the "surface"

e.g.: $\Sigma = \Sigma_1 dy \wedge dz + \Sigma_2 dz \wedge dx + \Sigma_3 dx \wedge dy$

Then $d\Sigma = d\Sigma_1 dy \wedge dz + d\Sigma_2 dz \wedge dx + d\Sigma_3 dx \wedge dy$

$$= (\frac{\partial \Sigma_1}{\partial x} dx + \dots) \wedge dy \wedge dz$$

$$+ (\dots + \frac{\partial \Sigma_2}{\partial y} dy + \dots) \wedge dz \wedge dx$$

$$+ (\dots + \frac{\partial \Sigma_3}{\partial z} dz) \wedge dx \wedge dy$$

$$= \left(\frac{\partial \Sigma_1}{\partial x} + \frac{\partial \Sigma_2}{\partial y} + \frac{\partial \Sigma_3}{\partial z} \right) dx \wedge dy \wedge dz$$

$$= \operatorname{div} \vec{F} dx \wedge dy \wedge dz$$

where $\vec{F} = \Sigma_1 \hat{i} + \Sigma_2 \hat{j} + \Sigma_3 \hat{k}$

Hence the divergence theorem can be written as:

$$\iiint_D dS = \iiint_D \left(\frac{\partial \xi_1}{\partial x} + \frac{\partial \xi_2}{\partial y} + \frac{\partial \xi_3}{\partial z} \right) dx dy dz \quad \begin{matrix} (+ve) \text{ oriented} \\ \text{volume} \end{matrix}$$

$$= \iiint_D \operatorname{div} \vec{F} dV = \iint_{S=\partial D} \vec{F} \cdot \hat{n} d\sigma \quad \begin{matrix} \vec{F} \cdot \hat{n} d\sigma \\ \text{outward} \end{matrix}$$

To see the relation between $\vec{F} \cdot \hat{n} d\sigma$ and S ,

we parametrize S :

$$\vec{r}(u, v) = x(u, v) \hat{i} + y(u, v) \hat{j} + z(u, v) \hat{k}$$

$$\Rightarrow \begin{cases} \vec{r}_u = x_u \hat{i} + y_u \hat{j} + z_u \hat{k} \\ \vec{r}_v = x_v \hat{i} + y_v \hat{j} + z_v \hat{k} \end{cases}$$

$$\Rightarrow \vec{r}_u \times \vec{r}_v = \begin{vmatrix} y_u & y_v \\ z_u & z_v \end{vmatrix} \hat{i} + \begin{vmatrix} z_u & z_v \\ x_u & x_v \end{vmatrix} \hat{j} + \begin{vmatrix} x_u & x_v \\ y_u & y_v \end{vmatrix} \hat{k}$$

If $\vec{r}_u \times \vec{r}_v$ is outward, then

$$\hat{n} = \frac{\vec{r}_u \times \vec{r}_v}{|\vec{r}_u \times \vec{r}_v|} \quad \text{and} \quad d\sigma = |\vec{r}_u \times \vec{r}_v| du dv \quad \begin{matrix} (\text{correct} \\ \text{orientation}) \end{matrix}$$

$$= |\vec{r}_u \times \vec{r}_v| du dv$$

then $\vec{F} \cdot \hat{n} d\sigma = \vec{F} \cdot \frac{\vec{r}_u \times \vec{r}_v}{|\vec{r}_u \times \vec{r}_v|} |\vec{r}_u \times \vec{r}_v| du dv$

$$= \left(\xi_1 \frac{\partial(y, z)}{\partial(u, v)} + \xi_2 \frac{\partial(z, x)}{\partial(u, v)} + \xi_3 \frac{\partial(x, y)}{\partial(u, v)} \right) du dv$$

$$= \xi_1 dy dz + \xi_2 dz dx + \xi_3 dx dy$$

$$= S$$

$$\begin{aligned} \therefore \iint_S \vec{F} \cdot \hat{n} d\sigma &= \iint_{(u,v)} \xi_1 dy dz + \xi_2 dz dx + \xi_3 dx dy \\ &= \iint_S \xi \\ S &= \partial D \end{aligned}$$

Hence divergence thm is

$$\boxed{\iint_D dS = \iint_{S=\partial D} \xi} \quad \xi = z - \text{fam}$$

eg3 Stokes' Thm

$$\vec{F} = M \vec{i} + N \vec{j} + L \vec{k} \Leftrightarrow \omega = M dx + N dy + L dz$$

$$\text{Then } d\omega = (L_y - N_z) dy dz + (M_z - L_x) dz dx \quad (\text{Ex!}) \\ + (N_x - M_y) dx dy$$

$$= (\vec{\nabla} \times \vec{F}) \cdot \hat{n} d\sigma \quad (\text{Ex!})$$

Stokes Thm becomes
 $\oint_C \vec{F} \cdot d\vec{r}$
 $C = \partial S$

$$\boxed{\oint_C \omega = \iint_S d\omega} \quad \iint_S (\vec{\nabla} \times \vec{F}) \cdot \hat{n} d\sigma$$

Generalization to manifold of n-dimension with boundary (Skipped)

- $M = n$ diiml Manifold (oriented)
- $\partial M = \text{boundary}$ (oriented with induced orientation)
- $\omega = (n-1)\text{-form on } M$ (smooth)

Then

$$\boxed{\int_M d\omega = \int_{\partial M} \omega}$$

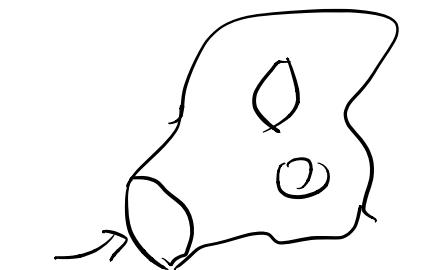
↑
 n -diiml
integral

↑
 $(n-1)$ -diiml
integral

Note: ∂M is always closed, i.e. no boundary.

$$\therefore \boxed{\partial(\partial M) = \partial^2 M = 0}$$

boundary has no boundary



∂S is a closed curve

Hence if $\omega = d\eta$, for some $(n-2)$ -form η , then

then

$$\int_M d(d\eta) = \int_M d\omega = \int_{\partial M} \omega$$

$$= \int_{\partial M} d\eta = \int_{\partial(\partial M)} \eta = 0 \quad (\text{fr any } \eta.)$$

This suggests $\boxed{d^2\eta = 0}$, & differential form

Ex: Verify this for 0-form and 1-form in \mathbb{R}^3
and observes that these are just

$$\left\{ \begin{array}{l} \vec{\nabla} \times \vec{\nabla} f = 0 \quad (d^2f = 0) \\ \vec{\nabla} \cdot (\vec{\nabla} \times \vec{F}) = 0 \quad (d^2\omega = 0) \end{array} \right.$$

e.g.: Let $\omega = \frac{-y}{x^2+y^2} dx + \frac{x}{x^2+y^2} dy$

check: $d\omega = 0$

But $\omega \neq df$ for any smooth function on $\mathbb{R}^2 \setminus \{(0,0)\}$

(Since $\omega = d\theta$ and θ is not defined on $\mathbb{R}^2 \setminus \{(0,0)\}$)

Hence $d\omega = 0 \not\Rightarrow \omega = dy$ in general

(\Leftarrow)
↑
yes

Note: Then Ω can be written as :

$\Omega \subset \mathbb{R}^2$ simply-connected, then smooth.

$d\omega = 0 \Leftrightarrow \omega = df$ for some f on Ω .