

# Tutorial 13 28/11/2014

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① Ex:  $M =$  regular surface

$\alpha =$  unit speed curve on  $M$  with length  $L$

$$\beta_\varepsilon(s) = \alpha(s) + \varepsilon \delta(s), \quad \varepsilon \in (\text{small interval containing } 0)$$

$$\langle \delta, \alpha' \rangle = 0 \text{ and } \langle \delta, \mathbb{U} \rangle = 0$$

$\mathbb{U} =$  Gauss map (i.e. unit normal)

$$\Rightarrow \left. \frac{d}{d\varepsilon} \right|_{\varepsilon=0} L(\beta_\varepsilon) = - \int_0^L k_g \langle \mathbb{U} \times T, \delta \rangle ds$$

Prmk: If  $\alpha =$  a geodesic, then  $\alpha$  is a critical pt of such length variation.

Pf:  $\beta'_\varepsilon = \alpha' + \varepsilon \delta', \quad |\beta'_\varepsilon| = \sqrt{1 + \varepsilon^2 |\delta'|^2 + 2\varepsilon \langle \alpha', \delta' \rangle}$

$$L(\beta_\varepsilon) = \int_0^L \sqrt{1 + \varepsilon^2 |\delta'|^2 + 2\varepsilon \langle \alpha', \delta' \rangle} ds$$

$$\left. \frac{d}{d\varepsilon} \right|_{\varepsilon=0} L(\beta_\varepsilon) = \int_0^L \left. \frac{d}{d\varepsilon} \right|_{\varepsilon=0} \sqrt{1 + \varepsilon^2 |\delta'|^2 + 2\varepsilon \langle \alpha', \delta' \rangle} ds$$

$$= \int_0^L \frac{1}{2} \cdot \left( 2\varepsilon |\delta'|^2 + 2 \langle \alpha', \delta' \rangle \right) \Big|_{\varepsilon=0} ds$$

$$= \int_0^L \langle \alpha', \delta' \rangle ds$$

$$\langle \delta, \alpha' \rangle = 0 \Rightarrow \langle \delta', \alpha' \rangle + \langle \delta, \alpha'' \rangle = 0.$$

$$\begin{aligned} \frac{d}{ds} \Big|_{s=0} L(\beta_2) &= - \int_0^L \langle \delta, \alpha'' \rangle ds \\ &= - \int_0^L \langle \delta, (\alpha'')_{\text{tan}} \rangle ds \quad \text{Since } \langle \delta, U \rangle = 0 \\ &= - \int_0^L k_g \langle \delta, U \times T \rangle ds. \end{aligned}$$

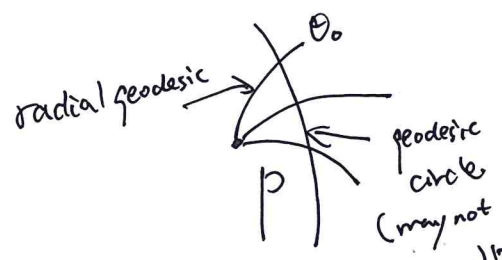
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### (2) Geodesic Polar Coordinates

( At any pt in a regular surface, we can find a local coordinate around that pt  $\Sigma_t$

$$g = \begin{pmatrix} 1 & 0 \\ 0 & G \end{pmatrix}.$$

Let  $p \in M$ ,  $(e_1, e_2)$  be an orthonormal basis of  $T_p M$ . For any unit vect.  $w = \cos \nu e_1 + \sin \nu e_2 \in T_p M$ , we can find  $\alpha_{p,w}(u)$ , a unit speed geodesic with  $\begin{cases} \alpha_{p,w}(0) = p \\ \alpha'_{p,w}(0) = w \end{cases}$  in an interval  $(-\epsilon_{p,w}, \epsilon_{p,w})$ .



Let  $\nu$  runs over  $[0, 2\pi)$ , then we get

a ~~set~~ coordinate  $X(u, \nu) = \alpha_{p, \cos \nu e_1 + \sin \nu e_2}(u)$ .

But we should check  $\epsilon_{p, \cos \nu e_1 + \sin \nu e_2}$  have a uniform lower bound. See §4.6 of do Carmo (of  $\nu$ )

for rigorous discussion.

$$g_{11} = \langle X_u, X_u \rangle = |\alpha'|^2 = 1.$$

$$(u, v) \triangleq (\rho, \theta) \quad (= (r, \theta))$$

$\alpha_{p, \cos \theta e_1 + \sin \theta e_2}^{(u)}$  is a geodesic  $\Rightarrow$   
 $(\theta = t, \theta = \text{Constant})$

$$\begin{cases} \frac{d^2 \rho}{dt^2} + \Gamma_{11}^1 \left(\frac{d\rho}{dt}\right)^2 + 2\Gamma_{12}^1 \frac{d\rho}{dt} \frac{d\theta}{dt} + \Gamma_{22}^1 \left(\frac{d\theta}{dt}\right)^2 = 0 \\ \frac{d^2 \theta}{dt^2} + \Gamma_{11}^2 \left(\frac{d\rho}{dt}\right)^2 + 2\Gamma_{12}^2 \frac{d\rho}{dt} \frac{d\theta}{dt} + \Gamma_{22}^2 \left(\frac{d\theta}{dt}\right)^2 = 0 \end{cases}$$

$$\Rightarrow \Gamma_{11}^1 = \Gamma_{11}^2 = 0$$

$$\Rightarrow \begin{cases} g^{11} (2\partial_1 g_{11} + 2\partial_1 g_{11} - 2\partial_1 g_{11}) = 0 \\ g^{21} (2\partial_1 g_{11} + 2\partial_1 g_{11} - 2\partial_1 g_{11}) = 0 \end{cases}$$

$$\Rightarrow \begin{cases} g^{12} (2\partial_1 g_{12}) = 0 \\ g^{22} (2\partial_1 g_{12}) = 0 \end{cases}$$

$\Rightarrow \partial_1 g_{12} = 0$  (\*) Since  $g^{12}, g^{22}$  can not vanish at the same time

~~And note that  $g_{12} = \langle X_1, X_2 \rangle = 0$  at  $p$  since  $X_2(p) = \frac{d}{d\theta} (\alpha(\theta))_{p, \cos \theta e_1 + \sin \theta e_2}$~~

~~$\Rightarrow g_{12} = 0$~~   $|g_{12}| \leq |\langle X_1, X_2 \rangle| = |X_2| \xrightarrow{**} 0$

~~We want to show more:~~  $\rightarrow 0$  as  $\rho \rightarrow 0$   
 Since  $\lim_{\rho \rightarrow 0} \frac{\partial}{\partial \theta} \alpha(\rho) = 0$  (Check it!)

~~$\lim_{\rho \rightarrow 0} g = 0$~~

By  $\otimes \otimes$ , we get  $g_{12} \equiv 0$ .

We would like to change coordinate to see

more.

$$\begin{cases} u = \rho \cos \theta \\ v = \rho \sin \theta \end{cases} \quad \left( (u, v) \triangleq \text{normal coordinate} \right)$$

$\theta_0 = \text{Constant}$  (radial geodesic)  $\Leftrightarrow (u, v) = t(u_0, v_0)$

In particular,  $X(t, 0) = \alpha_{\rho, e_1}^{\downarrow \text{polar coordinate}}(t)$

$X_u(0, 0) = \alpha_{\rho, e_1}^{\downarrow \text{polar coordinate}}(0) = e_1$  ( $X_u(0, 0) = e_1$ )  $\leftarrow \text{normal}$

$X(t, \frac{\pi}{2}) = \alpha_{\rho, e_2}^{\downarrow \text{polar coordinate}}(t)$

$X_v(0, \frac{\pi}{2}) = \alpha_{\rho, e_2}^{\downarrow \text{polar coordinate}}(0) = e_2$  ( $X_v(0, 0) = e_2$ )  $\leftarrow \text{normal}$

$\Rightarrow g_{\downarrow \text{normal}}(0, 0) = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$

$(u^1, u^2) \triangleq (u, v) = t(u_0, v_0)$

$\frac{d^2 u^k}{dt^2} + \Gamma_{ij}^k \frac{du^i}{dt} \frac{du^j}{dt} = 0, \quad k=1, 2$

$\Rightarrow \Gamma_{ij}^k u_0^i \cdot u_0^j = 0 \quad (u_0^1, u_0^2) \triangleq (u_0, v_0)$

Since  $(u_0, v_0)$  is arbitrary as long as  $u_0^2 + v_0^2 = 1$

$\Rightarrow \Gamma_{ij}^k = 0$  along the radial geodesic  $\forall (i, j) = 1, 2$

In particular,  $\Gamma_{ij}^k(p) = 0, \quad \forall i, j, k = 1, 2$ .

$$\frac{\partial}{\partial u^k} g_{u^i u^j} = \frac{\partial}{\partial u^k} \langle X_{u^i}, X_{u^j} \rangle$$

$$\boxed{X_{u^j} = b_{ij} \vec{n} + P_{ij}^k X_k} \quad \leftarrow$$

$$= \langle X_{u^i u^k}, X_{u^j} \rangle + \langle X_{u^i}, X_{u^j u^k} \rangle$$

$$= \langle P_{ik}^s X_{u^s}, X_{u^j} \rangle + \langle X_{u^i}, P_{jk}^s X_{u^s} \rangle$$

$$= 0 \text{ at } p.$$

$$\Rightarrow g_{u^i u^j} = \delta_{ij} + 2^{\text{nd}} \text{ or higher order terms.}$$

$$\sqrt{\det \begin{pmatrix} 1 & 0 \\ 0 & g_{u^i u^j}(p, \theta) \end{pmatrix}} = \sqrt{\det(g_{u^i u^j})} \frac{\partial(u, v)}{\partial(p, \theta)} \leftarrow (\text{Exercise!})$$

$$\text{Where } \frac{\partial(u, v)}{\partial(p, \theta)} = \det \begin{pmatrix} \frac{\partial u}{\partial p} & \frac{\partial u}{\partial \theta} \\ \frac{\partial v}{\partial p} & \frac{\partial v}{\partial \theta} \end{pmatrix} = \det \begin{pmatrix} \cos \theta & -\rho \sin \theta \\ \sin \theta & \rho \cos \theta \end{pmatrix} = \rho$$

$$\Rightarrow \sqrt{g(p, \theta)} = \rho \sqrt{\det(g_{u^i u^j})}$$

$$\Rightarrow \lim_{\rho \rightarrow 0} \sqrt{g} = 0$$

$$\left( \sqrt{g(p, \theta)} \right)_p = \sqrt{\det(g_{u^i u^j})} + \rho \left( \frac{\partial}{\partial \rho} \sqrt{\det(g_{u^i u^j})} \right)$$

$$\rightarrow 1 \text{ as } \rho \rightarrow 0.$$



~~$\langle X_{u^1}, X_{u^2} \rangle = \dots \rightarrow 0 \text{ as } \rho \rightarrow 0$~~

~~Together with the fact that~~

~~$\int_1^2 g_{12} = 0$  along the vertical geodesic~~

~~$\int_1^2 g_{12} = 0$~~

Claim: (i)  $\sqrt{G}(p, \theta) = e^{-\frac{e^3}{3!} K(p) + \cancel{(p^3)} R(p, \theta)}$  where  $\lim_{e \rightarrow 0} \frac{R(p, \theta)}{e^3} = 0$ .

(ii)  $K(p) = \lim_{e_0 \rightarrow 0} \frac{3}{\pi} \frac{2\pi e_0 - L_0}{e_0^3}$  where  $L_0 =$  length of the geodesic circle  $p = e$ .

To see (i), we need to show

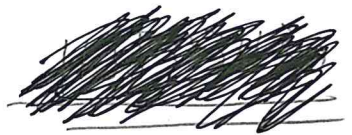
$\lim_{e \rightarrow 0} \sqrt{G} = 0$ ,  $\lim_{e \rightarrow 0} (\sqrt{G})_e = 1$ ,  $\lim_{e \rightarrow 0} (\sqrt{G})_{ee} = 0$

and  $\lim_{e \rightarrow 0} (\sqrt{G})_{eee} = -K(p)$ . (Although there is no def. at  $p$ , we may use limit to def. it.)

Recall a formula for Gauss curvature in the case of  $F=0$ :

$$K = -\frac{1}{2\sqrt{EG}} \left[ \left( \frac{E_v}{\sqrt{EG}} \right)_v + \left( \frac{G_u}{\sqrt{EG}} \right)_u \right]$$
$$= -\frac{1}{2\sqrt{G}} \left( \frac{\partial}{\partial e} \left( \frac{G_e}{\sqrt{G}} \right) \right) \quad \left( \frac{G_e}{\sqrt{G}} = 2 \frac{\partial}{\partial e} (\sqrt{G}) \right)$$
$$= -\frac{1}{2\sqrt{G}} 2 (\sqrt{G})_{ee}$$

$$(\sqrt{G})_{\rho\rho} = -\sqrt{G} k$$



$$\lim_{\rho \rightarrow 0} (\sqrt{G})_{\rho\rho} = 0$$



$$(\sqrt{G})_{\rho\rho\rho} = -(\sqrt{G})_{\rho\rho} k - k_{\rho} \sqrt{G}$$

$$\lim_{\rho \rightarrow 0} \downarrow = -k(\rho)$$

$$(ii) L_0 = \int_0^{2\pi} \sqrt{G}(\rho_0, \theta) d\theta$$

$$= \int_0^{2\pi} \left( \rho_0 - \frac{\rho_0^3}{3!} k(\rho) + \cancel{R(\rho_0, \theta)} \right) d\theta$$

$$= 2\pi \rho_0 - \frac{\pi}{3} \rho_0^3 k(\rho) + \int_0^{2\pi} R(\rho_0, \theta) d\theta$$

$$k(\rho) \Rightarrow \frac{3}{\pi} \left[ \frac{(2\pi \rho_0 - L_0)}{\rho_0^3} + \frac{3}{\pi} \int_0^{2\pi} \frac{R(\rho_0, \theta)}{\rho_0^3} d\theta \right]$$

$$= \lim_{\rho_0 \rightarrow 0} \frac{3}{\pi} \cdot \frac{2\pi \rho_0 - L_0}{\rho_0^3}$$

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