

SQ11

For  $l=0$ , the radial part of the Schrödinger equation is:

$$\hat{H}_r R(r) = \left[ -\frac{\hbar^2}{2\mu} \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) - \frac{e^2}{4\pi\epsilon_0 r} \right] R(r) = ER(r) \quad \dots \quad (1)$$

Applying variation method with trial function:

$$\phi(r) = A e^{-\lambda r^2} \quad \dots \quad (2)$$

where  $A = \sqrt{\frac{2\lambda}{\pi}}$  is a normalisation constant.

and  $\lambda$  is a parameter to be determined.

By variation method, we would like to find the best  $\lambda$  where  $\langle \hat{H}_r \rangle_\phi$

$\langle \hat{H}_r \rangle_\phi$  attains its minimum.

Compute  $\langle \hat{H}_r \rangle_\phi$  in terms of  $\lambda$ :

$$\begin{aligned} \langle \hat{H}_r \rangle_\phi &= \int_0^\infty \phi^*(r) \hat{H}_r \phi(r) r^2 dr \\ &= \int_0^\infty A^2 e^{-\lambda r^2} \left[ -\frac{\hbar^2}{2\mu} \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) - \frac{e^2}{4\pi\epsilon_0 r} \right] e^{-\lambda r^2} r^2 dr \\ &= A^2 \int_0^\infty e^{-\lambda r^2} \left[ -\frac{\hbar^2}{2\mu} \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 (-2\lambda r e^{-\lambda r^2}) \right) - \frac{e^2}{4\pi\epsilon_0 r} e^{-\lambda r^2} \right] r^2 dr \\ &= A^2 \int_0^\infty e^{-\lambda r^2} \left[ -\frac{\hbar^2}{2\mu} (-2\lambda) \frac{\partial}{\partial r} (r^3 e^{-\lambda r^2}) - \frac{e^2}{4\pi\epsilon_0 r} e^{-\lambda r^2} \right] r^2 dr \\ &= A^2 \int_0^\infty e^{-\lambda r^2} \left[ \frac{\hbar^2 \lambda}{\mu} \frac{\partial}{\partial r} (3r^2 e^{-\lambda r^2} - 2\lambda r^4 e^{-\lambda r^2}) - \frac{e^2}{4\pi\epsilon_0 r} e^{-\lambda r^2} \right] r^2 dr \\ &= A^2 \int_0^\infty \left[ -\frac{e^2}{4\pi\epsilon_0} r e^{-2\lambda r^2} + \frac{3\hbar^2 \lambda}{\mu} r^2 e^{-2\lambda r^2} - \frac{2\hbar^2 \lambda^2}{\mu} r^4 e^{-2\lambda r^2} \right] r^2 dr \end{aligned}$$

With some useful Gaussian integral formula:

$$\int_0^\infty r^4 e^{-\lambda r^2} dr = \frac{3}{8\lambda^2} \sqrt{\frac{\pi}{\lambda}}, \quad \int_0^\infty r^2 e^{-\lambda r^2} dr = \frac{1}{4\lambda} \sqrt{\frac{\pi}{\lambda}}, \quad \int_0^\infty r e^{-\lambda r^2} dr = \frac{1}{2\lambda}$$

(You can check them out from table or simply do it as an exercise)

$$\begin{aligned} \therefore \langle \hat{H}_r \rangle_\phi &= A^2 \left( -\frac{e^2}{4\pi\epsilon_0} \frac{1}{2(2\lambda)} + \frac{3\hbar^2\lambda}{\mu} \frac{1}{4(2\lambda)} \sqrt{\frac{1}{2\lambda}} - \frac{2\lambda^2\hbar^2}{\mu} \frac{3}{8(2\lambda)^2} \sqrt{\frac{1}{2\lambda}} \right) \\ &= \frac{3\lambda\hbar^2}{2\mu} - \frac{2e^2}{4\pi\epsilon_0} \sqrt{\frac{2\lambda}{\pi}} \quad \dots (3) \quad (\text{Remember } A^2 = \frac{8}{\sqrt{\frac{2\lambda}{\pi}}}) \end{aligned}$$

Now, we obtained  $\langle \hat{H}_r \rangle_\phi$  as a function of  $\lambda$ . (The function is sketched below)

The next step is to find the minimum  $\langle \hat{H}_r \rangle_\phi$  with respect to  $\lambda$ :

Diff. (3) with respect to  $\lambda$ , and set:

$$\left. \frac{\partial \langle \hat{H}_r \rangle_\phi}{\partial \lambda} \right|_{\lambda=\lambda_{\text{best}}} = 0$$

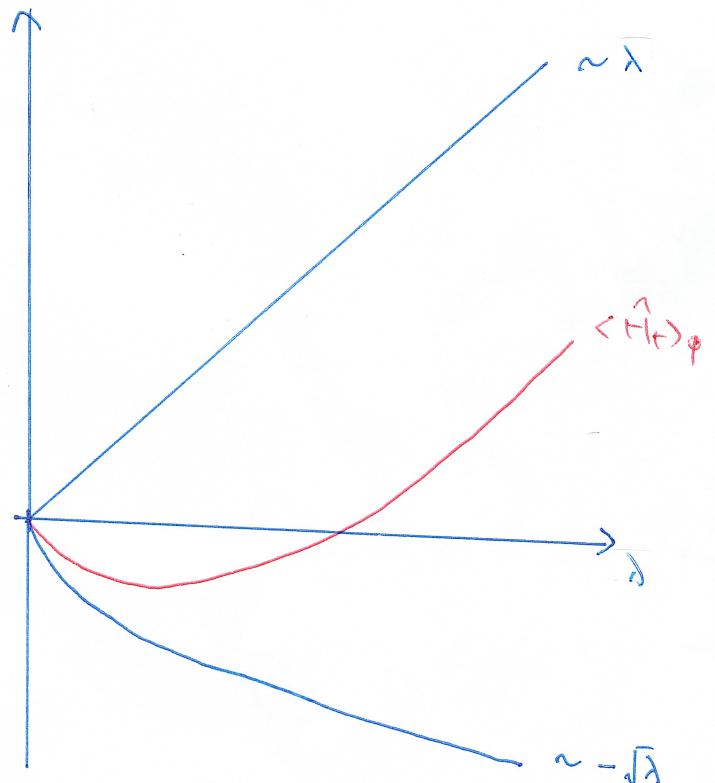
$$\therefore \frac{3\hbar^2}{2\mu} - \frac{e^2}{4\pi\epsilon_0} \sqrt{\frac{2}{\pi\lambda_{\text{best}}}} = 0$$

$$\sqrt{\lambda_{\text{best}}} = \frac{e^2}{4\pi\epsilon_0} \frac{2\mu}{3\hbar^2} \sqrt{\frac{2}{\pi}}$$

$$\text{or } \lambda_{\text{best}} = \frac{e^4}{18\pi^2\epsilon_0^2} \frac{\mu^2}{\hbar^4} //$$

The corresponding min  $\langle \hat{H}_r \rangle_\phi$  is:

$$\begin{aligned} \text{min } \langle \hat{H}_r \rangle_\phi &= -\frac{4}{3\pi} \frac{\mu e^4}{\hbar^2 (4\pi\epsilon_0)^2} \\ &\approx -0.424 \frac{\mu e^4}{\hbar^2 (4\pi\epsilon_0)^2} // \end{aligned}$$



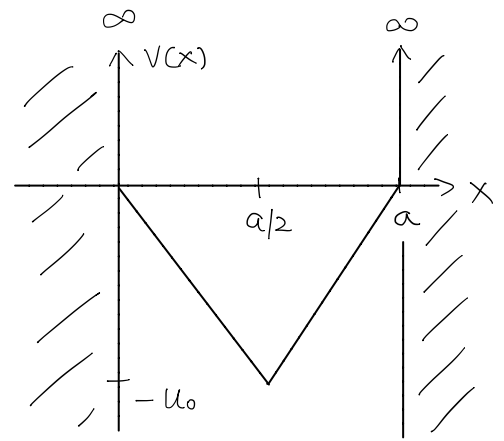
Compared to the true value  $E_{GS} = -0.5 \frac{\mu e^4}{\hbar^2 (4\pi\epsilon_0)^2}$ ,

1. The result is quite close, the variational method is a good estimate to the true hydrogen system.
2. The estimated value  $\text{min } \langle \hat{H}_r \rangle_\phi$  is slightly larger than  $E_{GS}$ , since the true ground state energy should have the lowest energy level.

SQ12

$$\hat{H} = -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x)$$

$$V(x) = \begin{cases} -\frac{2U_0}{a}x & \text{for } 0 < x < \frac{a}{2} \\ \frac{2U_0}{a}x - 2U_0 & \text{for } \frac{a}{2} < x < a \\ \infty & \text{otherwise} \end{cases}$$



$$\phi_{\text{trial}}(x) = c_1 \psi_1(x) + c_3 \psi_3(x)$$

$$= \begin{cases} c_1 \sqrt{\frac{2}{a}} \sin\left(\frac{\pi x}{a}\right) + c_3 \sqrt{\frac{2}{a}} \sin\left(\frac{3\pi x}{a}\right) & \text{for } 0 < x < a \\ 0 & \text{otherwise} \end{cases}$$

Variational method: find the best  $c_1$  and  $c_3$  values which give the best estimate of  $E_{\text{gs}}$  (i.e. minimize  $\langle \hat{H} \rangle_{\phi} = \langle \phi | \hat{H} | \phi \rangle / \langle \phi | \phi \rangle$ )

Using the same notation as lecture notes:

$$E(c_1, c_3) \equiv \langle \hat{H} \rangle_{\phi}$$

$$H_{ij} = \int \psi_i^* \hat{H} \psi_j dx$$

$$H_{ji}^* = H_{ij} \quad (\hat{H} \text{ is Hermitian})$$

$$S_{ij} = \int \psi_i^* \psi_j dx = \delta_{ij}$$

( $\psi_i$  are normalized eigenfunctions of 1D infinite well so they are orthonormal)

By Eq (B7) in notes AM-B18,

$$\begin{pmatrix} H_{11} - ES_{11} & H_{13} - ES_{13} \\ H_{31} - ES_{31} & H_{33} - ES_{33} \end{pmatrix} \begin{pmatrix} c_1 \\ c_3 \end{pmatrix} = 0 \quad (*)$$

SQ12

$$\begin{aligned}
H_{11} &= \frac{2}{a} \int_0^a \hat{\sin}\left(\frac{\pi x}{a}\right) \left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2}\right) \hat{\sin}\left(\frac{\pi x}{a}\right) dx + \frac{2}{a} \int_0^{a/2} \hat{\sin}\left(\frac{\pi x}{a}\right) \left(\frac{-2U_0}{a} x\right) \hat{\sin}\left(\frac{\pi x}{a}\right) dx \\
&\quad + \frac{2}{a} \int_{a/2}^a \hat{\sin}\left(\frac{\pi x}{a}\right) \left(\frac{2U_0}{a} x - 2U_0\right) \hat{\sin}\left(\frac{\pi x}{a}\right) dx \\
&= \frac{\hbar^2}{2m} \frac{2}{a} \frac{\pi^2}{a^2} \int_0^a \hat{\sin}^2\left(\frac{\pi x}{a}\right) dx - \frac{4U_0}{a^2} \int_0^{a/2} \hat{\sin}^2\left(\frac{\pi x}{a}\right) x dx \\
&\quad + \frac{4U_0}{a^2} \int_{a/2}^a \hat{\sin}^2\left(\frac{\pi x}{a}\right) x dx - \frac{4U_0}{a} \int_{a/2}^a \hat{\sin}^2\left(\frac{\pi x}{a}\right) dx \\
&= \frac{\hbar^2}{2m} \frac{\pi^2}{a^2} - \frac{4U_0}{a^2} \frac{a^2(\pi^2+4)}{16\pi^2} + \frac{4U_0}{a^2} \frac{a^2(3\pi^2-4)}{16\pi^2} - \frac{4U_0}{a} \frac{a}{4} \\
&= \frac{\hbar^2}{2m} \frac{\pi^2}{a^2} - \frac{U_0(\pi^2+4)}{2\pi^2}
\end{aligned}$$

$$\begin{aligned}
H_{33} &= \frac{2}{a} \int_0^a \hat{\sin}\left(\frac{3\pi x}{a}\right) \left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2}\right) \hat{\sin}\left(\frac{3\pi x}{a}\right) dx + \frac{2}{a} \int_0^{a/2} \hat{\sin}\left(\frac{3\pi x}{a}\right) \left(\frac{-2U_0}{a} x\right) \hat{\sin}\left(\frac{3\pi x}{a}\right) dx \\
&\quad + \frac{2}{a} \int_{a/2}^a \hat{\sin}\left(\frac{3\pi x}{a}\right) \left(\frac{2U_0}{a} x - 2U_0\right) \hat{\sin}\left(\frac{3\pi x}{a}\right) dx \\
&= \frac{\hbar^2}{2m} \frac{2}{a} \frac{9\pi^2}{a^2} \int_0^a \hat{\sin}^2\left(\frac{3\pi x}{a}\right) dx - \frac{4U_0}{a^2} \int_0^{a/2} \hat{\sin}^2\left(\frac{3\pi x}{a}\right) x dx \\
&\quad + \frac{4U_0}{a^2} \int_{a/2}^a \hat{\sin}^2\left(\frac{3\pi x}{a}\right) x dx - \frac{4U_0}{a} \int_{a/2}^a \hat{\sin}^2\left(\frac{3\pi x}{a}\right) dx \\
&= \frac{\hbar^2}{2m} \frac{9\pi^2}{a^2} - \frac{4U_0}{a^2} \frac{a^2(9\pi^2+4)}{144\pi^2} + \frac{4U_0}{a^2} \frac{a^2(27\pi^2-4)}{144\pi^2} - \frac{4U_0}{a} \frac{a}{4} \\
&= \frac{\hbar^2}{2m} \frac{9\pi^2}{a^2} - \frac{U_0(9\pi^2+4)}{18\pi^2}
\end{aligned}$$

SQ12

$$\begin{aligned}
 H_{13} &= \frac{2}{a} \int_0^a \sin\left(\frac{\pi x}{a}\right) \left(-\frac{\hbar^2}{2m} \frac{d^2}{dx^2}\right) \sin\left(\frac{3\pi x}{a}\right) dx + \frac{2}{a} \int_0^{a/2} \sin\left(\frac{\pi x}{a}\right) \left(-\frac{2U_0}{a} x\right) \sin\left(\frac{3\pi x}{a}\right) dx \\
 &+ \frac{2}{a} \int_{a/2}^a \sin\left(\frac{\pi x}{a}\right) \left(\frac{2U_0}{a} x - 2U_0\right) \sin\left(\frac{3\pi x}{a}\right) dx \\
 &= \frac{\hbar^2}{2m} \frac{2}{a} \frac{9\pi^2}{a^2} \int_0^a \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{3\pi x}{a}\right) dx - \frac{4U_0}{a^2} \int_0^{a/2} \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{3\pi x}{a}\right) x dx \\
 &+ \frac{4U_0}{a^2} \int_{a/2}^a \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{3\pi x}{a}\right) x dx - \frac{4U_0}{a} \int_{a/2}^a \sin\left(\frac{\pi x}{a}\right) \sin\left(\frac{3\pi x}{a}\right) dx \\
 &= -\frac{4U_0}{a^2} \left(-\frac{a^2}{4\pi^2}\right) + \frac{4U_0}{a^2} \left(\frac{a^2}{4\pi^2}\right) \\
 &= \frac{2U_0}{\pi^2}
 \end{aligned}$$

$$H_{31} = H_{13}^* = \frac{2U_0}{\pi^2}$$

$\therefore$  (\*) becomes

$$\begin{pmatrix} H_{11} - ES_{11} & H_{13} - ES_{13} \\ H_{31} - ES_{31} & H_{33} - ES_{33} \end{pmatrix} \begin{pmatrix} C_1 \\ C_3 \end{pmatrix} = 0$$

$$\begin{pmatrix} \frac{\hbar^2 \pi^2}{2m a^2} - \frac{U_0(\pi^2 + 4)}{2\pi^2} - E & \frac{2U_0}{\pi^2} \\ \frac{2U_0}{\pi^2} & \frac{\hbar^2 9\pi^2}{2m a^2} - \frac{U_0(9\pi^2 + 4)}{18\pi^2} - E \end{pmatrix} \begin{pmatrix} C_1 \\ C_3 \end{pmatrix} = 0$$

SQ12

To have non-trivial solution (i.e.  $C_1 \neq 0$  and  $C_2 \neq 0$ ), we have

$$\begin{vmatrix} \frac{\hbar^2 \pi^2}{2m a^2} - \frac{U_0(\pi^2+4)}{2\pi^2} - E & \frac{2U_0}{\pi^2} \\ \frac{2U_0}{\pi^2} & \frac{\hbar^2 9\pi^2}{2m a^2} - \frac{U_0(9\pi^2+4)}{18\pi^2} - E \end{vmatrix} = 0$$

$$\left[ \frac{\hbar^2 \pi^2}{2m a^2} - \frac{U_0(\pi^2+4)}{2\pi^2} - E \right] \left[ \frac{\hbar^2 9\pi^2}{2m a^2} - \frac{U_0(9\pi^2+4)}{18\pi^2} - E \right] - \frac{4U_0^2}{\pi^4} = 0$$

$$E = \frac{5\hbar^2 \pi^2}{2ma^2} - \frac{U_0(9\pi^2+20)}{18\pi^2} \pm \sqrt{\frac{4\hbar^4 \pi^4}{m^2 a^4} + \frac{32\hbar^2 U_0}{9ma^2} + \frac{388U_0^2}{81\pi^4}}$$

The lower one is an estimate to  $E_{GS}$

$$\begin{aligned} \langle \hat{H} \rangle_{\phi, \min} &= \frac{5\hbar^2 \pi^2}{2ma^2} - \frac{U_0(9\pi^2+20)}{18\pi^2} - \sqrt{\frac{4\hbar^4 \pi^4}{m^2 a^4} + \frac{32\hbar^2 U_0}{9ma^2} + \frac{388U_0^2}{81\pi^4}} \\ &= \frac{5\hbar^2 \pi^2}{2ma^2} - \frac{U_0(9\pi^2+20)}{18\pi^2} - \frac{2\hbar^2 \pi^2}{ma^2} \sqrt{1 + \frac{8U_0 ma^2}{9\hbar^2 \pi^4} + \frac{97U_0^2 m^2 a^4}{81\hbar^4 \pi^8}} \end{aligned}$$

Remarks =

1) when  $U_0 = 0$ ,  $\langle \hat{H} \rangle_{\phi, \min} = \frac{\pi^2 \hbar^2}{2ma^2}$  which is the GS

Energy of 1D infinite well. It should be expected!